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## „Internal color gradients and distributions of stellar populations of early-type galaxies in z ~ 0.4 CLASH clusters"

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## The history of astronomy

is a history of receding horizons.
EDWIN POWELL HUBBLE,
THE REALM OF THE NEBULE

## Abstract english:

To better understand the processes involved in the formation and evolution of elliptical galaxies, 2D models were created in the course of this thesis, which offer the possibilities to constrain the probabilities of various scenarios. The galaxies are members of the clusters MACS J1206.2-0847 and MACS J0416.1-2403 at an average redshift of $z \sim 0.4$. Utilizing the, by the models determined structural parameters Sérsic index $n$ and effective radius $r_{e}$, as well as the total integrated magnitude, radial surface brightness profiles were derived for each of the observed galaxies in each of the twelve bands along with radial profiles for the four colors $g_{475}-I_{814}$, $r_{625}-Y_{105}, I_{814}-H_{160}$ and $Y_{105}-H_{160}$. Thereupon the logarithmic slopes of these color profiles were determined, as well as the influence on them from the ages and metallicities of the underlying stellar populations. It became apparent that the sizes of the stellar structures on average decrease with increasing wavelength, while $n$ remains constant. Already this observation indicates on average negative color gradients, which are indeed also found. As a consequence one can conclude that the stars are redder in the centers of the galaxies than in the exterior regions. Whether this is the result of radial variations in age and/or metallicity should be answered with the help of models of stellar populations. To that end for the mentioned colors a variety of values were simulated at the redshifts of the galaxy clusters, being based on differing formation redshifts and metallicities. A comparison of the simulated against the measured values yields that both a radial difference in ages and especially a metallicity gradient are the reasons for the observed color gradients. The magnitudes of these gradients favour a scenario, which predicts with progression in time a passive evolution combined with minor dry mergers with mass-poor systems, which are relaxing mainly into a more extended structure than the initial one for very distant $(z \sim 2)$, compact, but mass-rich elliptical galaxies, which means that the galaxies undergo an inside-out-growth.

## Abstract deutsch:

Um die Vorgänge bezüglich der Entstehung und Entwicklung von elliptischen Galaxien besser zu verstehen, wurden im Zuge dieser Arbeit 2D-Modelle dieser Objekte erstellt, die es ermöglichten, die Wahrscheinlichkeit diverser Szenarios einzuschränken. Die Galaxien sind Mitglieder der Haufen MACS J1206.2-0847 und MACS J0416.1-2403 bei einer mittleren Rotverschiebung von $z \sim 0.4$. Mithilfe der von den Modellen ausgegebenen strukturellen Parametern Sérsic Index $n$ und Effektiv Radius $r_{e}$, sowie der integrierten, totalen Helligkeit wurden für die Beobachtungen in zwölf verschiedenen Filtern für jede Galaxie jeweils ein radiales Oberflächenhelligkeitsprofil, als auch anschließend die radialen Profile für die vier Farben $g_{475}-I_{814}, r_{625}-Y_{105}$, $I_{814}-H_{160}$ und $Y_{105}-H_{160}$ erstellt. Daraufhin wurden die logarithmischen Steigungen dieser Farbprofile bestimmt, ebenso wie der Einfluss von Alter und Metallizität der zugrunde liegenden stellaren Populationen auf diese. Es stellte sich heraus, dass die Durchmesser der stellaren Strukturen im Mittel mit zunehmender Wellenlänge kleiner werden, während $n$ konstant bleibt. Bereits dies deutet im Durchschnitt gesehen auf negative Farbgradienten hin, die auch tatsächlich gefunden wurden. Als Konsequenz daraus lässt sich schließen, dass die Sterne in den Zentren der Galaxien röter sind als in den äußeren Bereichen. Ob dies nun ein Resultat aus radialen Variationen im Alter und/oder der Metallizität sei, sollte mithilfe der Modellierung von stellaren Populationen ermittelt werden. Zu diesem Zweck wurde eine Vielzahl an Werten für die bereits erwähnten Farben bei der Rotverschiebung der Galaxiehaufen erstellt, variierend aufgrund verschiedener „Entstehungs-rotverschiebungen ", sowie Metallizitäten. Ein Abgleich der simulierten mit den gemessenen Werten ergab, dass sowohl eine radiale Differenz des Alters, als auch insbesondere ein Metallizitätsgradient die Gründe für die beobachteten Farbgradienten sind. Die Größenordnungen dieser Gradienten bevorzugen ein Szenario, das für weit entfernte ( $z \sim 2$ ), kompakte, aber massereiche elliptische Galaxien mit Verlauf der Zeit nur noch eine passive Evolution vorhersagt, kombiniert mit gasarmen Verschmelzungen mit masseärmeren Systemen, die hauptsächlich eine ausgedehntere Struktur bilden als die ursprüngliche und die Galaxien somit von innen nach außen hin wachsen.

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## Introduction

Already in 1755 Immanuel Kant coined the term 'Welteninseln' ('island universes') for diffuse nebula-like objects observed in the night sky and assumed the universe consisting and containing numerous of those objects. But it was only in the twenties of the 20th century, that the nature of those structures became a matter of debate.

By this time several observations and measurements of those 'nebulae' took place, for example the creation of the General Catalogue of Nebulae by William and John Herschel containing survey results of the Northern and Southern hemisphere, the additions to this catalogue by John Dreyer, who consolidates these results into the New General Catalogue, which is still very much in use, or Vesto Melvin Slipher's measurements revealing the spectral lines of those objects being redshifted.
All these new insights lead to more questions than answers and on the 26th April 1920 in the Baird auditorium of the Smithsonian Museum of Natural History, Harlow Shapley and Heber Curtis were the two opponents of a debate, which was titled 'The scale of the universe'. The main question was whether these nebulae were relatively small and therefore part of our Milky Way or independent systems, making them exceedingly large and especially distant objects. Shapley favoured a scenario, in which the Milky Way is in fact the whole visible universe, while Curtis believed in those nebulae just being independent systems, much like the Milky Way itself.

Only a few years later Edwin Hubble (Hubble, 1926) was able to resolve this issue using Cepheid variables as distance indicators. As this class of stars has a pulsation period, which is precisely proportional to their luminosities, measuring their apparent magnitude can be directly used to calculate the distances, which makes them so-called standard candles. Hubble was therefore able to eventually establish those faint nebulae as distant, large and independent objects we nowadays know as galaxies, planting the seed of extragalactic research.
Our knowledge about these structures has grown enormously over the last decades, not only, but also thanks to the technological advances. Not even mentioning the emergence of powerful computers, which enable theoreticians to create complex simulations ranging from the formation of single stars to the whole universe, also obser-
vations became more precise, efficient and numerous, be it, among other things, due to the introduction of Charged Coupled Devices (CCDs), the possibilities to detect objects not only in the visible range of the spectrum, but also for example in the UV, IR or X-ray regimes thanks to satellites or last but not least the construction of monumental observatories, like the Mauna-Kea Observatories or the Paranal ESO Observatory.

Nonetheless there are still countless questions, which desire an answer. It is well known, that the nowadays estimated up to a total of $100-200$ billion galaxies in the universe, come in all shapes, sizes, colors and environments. Thus they can be found in the field as single objects with no gravitational interaction with any other galaxy right up to clusters containing hundreds, if not thousands of galaxies, or even superclusters, representing presumably the largest gravitational bound structures in the universe.

Still almost all galaxies can be described by a bimodal distribution, segregating them into blue spirals with young stellar populations in a prominent disk component and red, old elliptical galaxies. How those two populations of galaxies form and evolve over time is an ongoing topic of research and due to the finite speed of light and the launch of the Hubble Space Telescope ( $H S T$ ) and other extraordinary telescopes, we are also in the lucky position to be able to take a look into the past and therefore shed some light on the different formation and evolution scenarios of galaxies, by observing at different redshifts and also galaxies in different environments (e.g. Silk \& Mamon (2012); Conselice (2014)).

An important part plays the usage of the observed structures of the galaxies to defer properties directly related to the formation and evolution. One well established option is to examine the light profiles of the galaxies as a function of radius. Applying this method to different wavelengths and consequently exploring how the light distribution changes with different passbands (i.e. filters), opens up the possibility to derive the colors in dependence of radius. These colors, respectively their change with radius, i.e. the color-gradients, are directly correlated to the underlying spatial distribution of the stellar population of a galaxy, therefore making it possible if not to decipher, but to constrain possible formation scenarios of galaxies.

The first two chapters intend to give an overview of the theoretical framework of this thesis, as well as describing the tools applied to facilitate this study. The third chapter depicts the data, while the fourth and fifth outline the modeling of the sample galaxies and the comparison of the derived color-gradients with Simple

Stellar Population (SSP) models. In the sixth and seventh chapter results and an analysis are given, as well as a comparison to other publications and studies. Finally in the eighth and last chapter a summary of this thesis is presented.

## 1. Theory \& observations

This chapter intends to give insights on the needed background without going into too many details. The first three sections focus on observational results, the last part on the rather numerical element of simulating galaxy populations. It also gives an overview of the current status regarding the formation and evolution of galaxies to put the results presented in this thesis in up to date context.

### 1.1. Classification of galaxies

Galaxies come in a wide variety, consisting of stars formed possibly at different times and are maybe even still in the process of creating new stars. They possess different amount of gas and dust and are embedded into different environments. Galaxies therefore occur in all kinds of appearances and in order to make comparisons among diverse galaxies easier, one tries to classify them into different categories.

Edwin Hubble did not only resolve the issue of whether the observed nebulae are independent objects or not and in addition formulated his famous law correlating the recession velocity of an extragalactical object with its distance, but he also developed his notable tuning fork scheme (Hubble, 1926) for classifying galaxies solely based on their observed appearances and segregating them into elliptical early-type galaxies and late-type spiral galaxies. The former population is further distinguished by the parameter $n$ in regard to their projected ellipticity $e=1-b / a$ with $n=10 \times e$. The spiral population in turn is divided by taking their disk/bulge ratio into account and also the possible presence of a bar, marked by the extra ' B ', ranging from Sa ( SBa ), describing a spiral galaxy with tightly wound smooth spiral arms and a large, bright bulge to $\mathrm{Sc}(\mathrm{SBc}$ ), a galaxy showing very loose, fragmented spiral arms and a small, faint bulge.

With today's knowledge about galaxies this classification scheme seems far too superficial, but nevertheless aside from minor modifications, like the addition of a few more galaxy types to be able to classify for example the lenticular $\mathrm{S0}$ galaxies, in this thesis counted as early-type galaxies, or irregular galaxies, which are showing no
distinctive features, (e.g. Sandage (1961); Kormendy \& Bender (2012)), it is still in use (Fig. 1.1). Therefore we nonetheless exploit mainly these two categories, early and late-type, to classify observed galaxies.


Figure 1.1.: A modern version of the conventional Hubble tuning fork to classify galaxies. Taken from Kormendy \& Bender (2012).

The validity of this approach is supported by the fact, that numerous physical properties of galaxies correlate with their morphologies and show also a bimodal distribution. It has been well known for more than four decades that early-type galaxies are more massive and redder (Fig. 1.2), with these galaxies occupying a tight region in the color-magnitude-space (the so-called red sequence) out to $z \sim 2$ (Kodama \& Arimoto, 1997) and also showing little to no star formation in comparison to late-type galaxies, which leaves them passively evolving (Holmberg, 1958). But this bimodality does not only concern mass, color and star formation, but also for example metallicity $Z$, gas content, mass-to-light ratio $(M / L)$ and the bulge fraction (Conselice, 2006). So early-type galaxies possess a higher metallicity, low or almost no gas content, a large $M / L$ and a very high bulge fraction in the form of those galaxies being pure bulges.

Considering all these correlations one can conclude, that the morphology of a galaxy is the product of the formation and especially its evolution over time, including, amongst other factors, possible interactions with their environment or other galaxies in the form of mergers or internal perturbations (Buta, 2011).

Nevertheless it must be considered, that although distant galaxies seem to exhibit same morphologies as local galaxies, they do not possess the same properties like sizes, colors or star formation rate. Consequently the conventional Hubble sequence is not completely established at higher redshifts, is not present to the full extent at


Figure 1.2.: Left side: Color-magnitude diagram depicting the red sequence of early-type galaxies, shown as red squares. Later-types are displayed as blue crosses and intermediates as green triangles. Right side: Fractions of different morphological types at different stellar masses. The solid, red line represents early-types, the dashed green line spirals and the blue dotted line irregular galaxies. Diagrams taken from (Conselice, 2006).
$z \gtrsim 1$ (van den Bergh et al., 1996) and entirely vanishes at $z \gtrsim 2$. Hence spiral and elliptical galaxies in the classical sense following Hubble's classification only appear at $z \sim 1.5$ with the fraction of early-type galaxies increasing from $\sim 20-30$ percent at $z \sim 3$ to $\sim 70$ percent at $z \sim 0$ with elliptical galaxies being the dominating morphological type for massive galaxies since only $z \sim 1$ (Buitrago et al., 2013). However this trend seems to show a dependency on mass as the transition redshift, where the combined fraction of early- and late-type galaxies equal that of peculiars, shifts from higher redshifts $(z \sim 2.22 \pm 0.82)$ for higher mass galaxies $\left(\log \mathrm{M}_{*}>\right.$ 10.5) to lower redshifts $(z \sim 1.73 \pm 0.57)$ for lower mass galaxies $\left(\log \mathrm{M}_{*}<10.25\right)$ (Mortlock et al., 2013).

Yet in the local universe the distinction between early- and late-type galaxies is an important indication for the understanding of the physical processes involved in the formation of galaxies. As spheroidal galaxies (ellipticals and S0) contain $\sim 60$ percent of all stellar mass (Hogg et al. (2002b);Driver et al. (2006)) and the highest numbers of early-type galaxies can be found in highest density regimes (Dressler (1980);Dressler (1984)), i.e. the center of clusters, elliptical galaxies are the targets of this analysis. Furthermore at the redshifts examined in this thesis they represent the majority of galaxies and an understanding of their evolution is an important part to explain the general development of galaxies.

### 1.2. Formation and evolution of galaxies

One way to gain insights into the formation and evolution of galaxies is to examine the structure of galaxies, which links directly to galaxy properties from which the evolution for single galaxies can be derived. Hence it is naturally a complex issue to describe the formation and evolution of all galaxies in its entirety, due to their wide variety in numerous properties, e.g. star formation history, formation epoch, sizes, mass. Nevertheless the extragalactic research made great strides in the last decades and the knowledge regarding this manifold topic increased significantly. As early-type galaxies are the topic of this thesis the focus lies on the description of the formation and evolution of spheroidal galaxies, although of course the general conclusions are also applicable to late-type galaxies, all the more, as spiral galaxies can also play an important role in the progression of elliptical galaxies, due to interactions.

Historically two main frameworks for the formation of galaxies exist: The (revised) monolithic or top-down model and the hierarchical or bottom-up scenario. The former predicts a formation of spheroids at a very early epoch of the universe resulting from a global starburst followed by a passive evolution to the present day. Subsequently spheroids may form a disk component by accreting gas from the surrounding environment causing limited incidents of star formation, thus making spheroids preceding disks. The latter scheme envisions big spheroids as the result of a gas-rich, 'wet', merging event of two disk galaxies, thus making disks preceding spheroids, leading to an intense burst of star formation, forming the bulk of the newly created elliptical galaxy (Renzini, 2006). Another possibility would be a 'dry' merger scenario with passive galaxies forming a brighter more massive early-type galaxy.

Unfortunately both models are found to have shortcomings in comparison to simulations and also cannot completely explain observations, such as the down-sizing problem describing more massive galaxies being older, thus contradicting the hierarchical scheme. In the wake of these complications a new scheme arose and became widely accepted, termed 'inside-out-growth', one might say an adaption of the hierarchical scheme. Supported by observations and simulations early-type galaxies still form at high redshifts $(z \sim 4-5)$, but in contrast to the monolithic model as compact spheroids resulting from gas-rich mergers. This would also clarify the down-sizing phenomena, as the star formation starts first in the regions of highest densities, which will develop into the most massive ( $\gtrsim 10^{11} \mathrm{M}_{\odot}$ ) galaxies.

With time those compact early-type galaxies are taking part in subsequent 'dry', minor mergers increasing the galaxy in size, but hardly in mass. This picture is consistent with the findings of Driver et al. (2013), describing a turning point of the cosmic star formation history for the Universe at $z \sim 1.7$ shifting from a 'hot mode evolution' - spheroid formation and growth via mergers and/or collapse - to a 'cold mode evolution', characterizing a disk formation or growth by gas infall and minor mergers.

Indeed this scenario can also be described successfully by hydrodynamical cosmological simulations (Naab et al., 2009) and by using the Virial theorem one can demonstrate why minor mergers should increase the size of a galaxy, but not the mass to the same extent (Bezanson et al., 2009). Denoting the kinetic energy of the remnant after the merger as the sum of the kinetic energies of the progenitor galaxies

$$
\begin{equation*}
K_{R}=K_{1}+K_{2}=\frac{1}{2} M_{R} v_{R}^{2}=\frac{1}{2} M_{1} v_{1}^{2}+\frac{1}{2} M_{2} v_{2}^{2} \tag{1.1}
\end{equation*}
$$

and considering that $M_{R}=M_{1}+M_{2}$ with $M_{1} \gg M_{2}$, which leads to

$$
\begin{equation*}
\frac{M_{R}}{M_{1}} \approx \frac{v_{1}^{2}}{v_{R}^{2}} \tag{1.2}
\end{equation*}
$$

as well as $v^{2} \propto \frac{G M}{r}$ one arrives at

$$
\begin{equation*}
\frac{r_{R}}{r_{1}} \approx \frac{M_{R}}{M_{1}} \frac{v_{1}^{2}}{v_{R}^{2}} \approx\left(\frac{M_{R}}{M_{1}}\right)^{2} . \tag{1.3}
\end{equation*}
$$

Thus the radius grows by the square of the mass, rather than just linearly, which would be the case for a merger of equal mass galaxies $\left(M_{1}=M_{2}\right)$.

Numerous observations exist to confirm both parts of this theory, one being the early formation of the core of early-type galaxies, the other the growth in size. So it is meanwhile well established, that the majority of stars of spheroids were formed at high redshifts ( $z \approx 3-5$ ) and on short timescales ( $\tau \sim 1$ Gyr) (Thomas et al., 2005, 2009). This notion is supported by several correlations between global properties of early-type galaxies, such as color-magnitude, the fundamental plane (Bender et al., 1992) or the Mg- $\sigma$-relation (Ziegler \& Bender, 1997).

The case of small sized, massive galaxies at high redshifts $(z \sim 1.5-2)$ is likewise well observed, requiring a growth in size, but not in mass, which is most likely achieved by means of minor mergers (Daddi et al., 2005; Trujillo et al., 2007, 2011;
van Dokkum et al., 2008, 2010). As a consequence this would result in an evolution of the effective radius by a factor of $\sim 1.5$ from $z \sim 1.5$ to $z=0$ (Buitrago et al., 2008; Longhetti et al., 2007) and a factor of $\sim 4$ since $z \sim 2$ (Chan et al., 2016), which relates to $r \propto(1+z)^{-1.48}$ (van der Wel et al., 2014). Furthermore such small, but massive galaxies have no local counterpart, endorsing that those galaxies are the progenitors of the local massive early-type galaxies.

However recent studies show that this inside-out-growth scenario is only valid for high-mass early-type galaxies. For galaxies in the mass range like our Milky Way ( $\sim 10^{10} \mathrm{M}_{\odot}$ ) and also containing disks, the mass growth appears at all radii (van Dokkum et al., 2013).

### 1.3. Light profiles and color gradients

### 1.3.1. Sérsic profile

A simple and common way to measure magnitudes and light distributions is aperture photometry, in which one measures the intensities at certain radii and ends one's measurements at a chosen radius, where the flux does not seem to increase significantly any more. Ideally this radius corresponds to an aperture large enough to encompass all the flux of the object. However due to the noise from the background and maybe also from interfering neighbours the choice of finding the right aperture is a complex task and in most cases the way in which the radius is chosen leads to a loss of flux. Furthermore multiple images and therefore also varying signal to noise ( $\mathrm{S} / \mathrm{N}$ ) values, various light distributions and also different wavelengths can impede such efforts. Especially observing galaxies at different wavelengths implicates disparate results, as the diverse fluxes in differing wavelengths result in different apertures leading to wrong estimates.

One approach to evade such problems is to describe the intensity progression with radius and in the course of that also the structure of galaxies with integrated light profiles, which is one of the earliest, but nevertheless still commonly used methods. Simply spoken, one measures the average intensity of a galaxy at a certain radius and determines from that how the intensity changes with centric distance.

The first one to apply this method on a large sample was Gérard de Vaucouleurs, who demonstrated that practically all early-type galaxies show the same smooth light distribution, following the de Vaucouleurs' $R^{1 / 4}$ law (de Vaucouleurs, 1948). Almost a decade later José Luis Sérsic started his work at the 1.54 m -telescope at the

Astrophysical Station at Bosque Alegre in Argentina, with his groundbreaking work being summarized in his galaxy atlas 'Galaxias Australes' (Sérsic, 1968). Already a few years before publishing his famous work he deduced a generalised version of the de Vaucouleurs' law to describe not only early-type but all types of galaxies (Sérsic, 1963). Applying his Sérsic's $R^{1 / n}$ law to all sufficiently large galaxies in his atlas, he was in the position to derive total magnitudes and structural parameters of those objects, which eventually showed a high correlation with the visually observed morphologies.

The Sérsic profile is best known in the version of an intensity profile (Graham \& Driver, 2005), which outlines as

$$
\begin{equation*}
I(R)=I_{e} \exp \left\{-b_{n}\left[\left(\frac{R}{r_{e}}\right)^{1 / n}-1\right]\right\} . \tag{1.4}
\end{equation*}
$$

$I_{e}$ is the intensity at the effective radius $r_{e}$, which describes the radius at which half of the total emitted light of the galaxy is encompassed, which is ensured by the dimensionless constant $b_{n}$. The last parameter $n$, also called Sérsic index is responsible for the curvature of the profile and it hence provides information about the morphology of the scrutinized galaxy (Fig. 1.3).

So indicates an $n=4$ a light distribution related to an early-type galaxy as the Sérsic profiles turns into the known de Vaucouleurs' $R^{1 / 4}$ law for spheroidal galaxies, whereas an $n=1$ converts the Sérsic profile into an exponential profile, best fitting for disks (Fig. 1.3). Therefore an early-type galaxy can be easily described by a simple one component Sérsic profile, unlike spiral galaxies, which require usually a combination of two profiles with different $n$, due to their spheroid galaxy resembling bulge component and additional disk element, a circumstance first described by Freeman (1970). Of course the presence of further features like bars may demand even more sophisticated model combinations. As a result a particular $n$ was introduced to separate the early- and late-type populations. Most commonly used and also adapted in this thesis is $n=2.5$ (Shen et al., 2003), as this distinction also marks the exact average between the exponential and de Vaucouleurs form of the Sérsic profile.

Fig. 1.3 shows the Sérsic profiles for fixed $I_{e}$ and $r_{e}$ and also the, with different $n$ correlating characteristics. The decreasing central concentration of the light distribution is evident, as well as the steepening slope of the profile at larger radii with smaller $n$. Hence early-type galaxies with usually a larger $n$, show a high density
core, but a significant outer wing, extending outwards to large radii, making it difficult to find the edge of a galaxy, as its light blurs into the sky background noise. For an $n=4$ profile, 84.7 percent of the flux are within $4 r_{e}$, for $5 r_{e}$ the contained flux amounts to 88.4 percent (Graham \& Driver, 2005).

To acquire the luminosity $L$ within


Figure 1.3.: Different Sérsic profiles with same $I_{e}$ and $r_{e}$. The distinctions in the center and large radii depending on $n$ can easily be seen. Diagram from (Peng et al., 2010). a certain radius $R$, one can integrate Equation (1.4) over the area of a disk, $A=\pi R^{2}$ :

$$
\begin{equation*}
L(<R)=\int_{0}^{R} I\left(R^{\prime}\right) 2 \pi R^{\prime} d R^{\prime} . \tag{1.5}
\end{equation*}
$$

After substituting $x=b_{n}\left(\frac{R}{r_{e}}\right)^{1 / n}$ this results into

$$
\begin{equation*}
L(<R)=I_{e} r_{e}^{2} 2 \pi n \frac{e^{b_{n}}}{\left(b_{n}\right)^{2 n}} \gamma(2 n, x), \tag{1.6}
\end{equation*}
$$

with $\gamma(2 n, x)$ being the incomplete gamma function. Replacing $\gamma(2 n, x)$ with the complete gamma function $\Gamma(2 n)=\gamma(2 n, \infty)$ equation (1.6) yields
the value for the total luminosity $L_{\text {tot }}$ (Ciotti, 1991).
Thus considering the definition of $r_{e}, b_{n}$ is the solution of the equation

$$
\begin{equation*}
\gamma\left(2 n, b_{n}\right)=\frac{\Gamma(2 n)}{2} \tag{1.7}
\end{equation*}
$$

which can only be solved numerically. Ciotti \& Bertin (1999) did an asymptotic expansion for $b_{n}$ and concluded that even a truncation after the first four terms shows better results than previous solutions, with errors being smaller than $\sim 10^{-6}$. Therefore their equation for $b_{n}$

$$
\begin{equation*}
b_{n}(n)=2 n-\frac{1}{3}+\frac{4}{405 n}+\frac{46}{25525 n^{2}} \tag{1.8}
\end{equation*}
$$

is used throughout this thesis.

### 1.3.2. Surface brightness profile

Taking the negative logarithm of equation (1.4) and multiplying the result with 2.5 the result is the surface brightness profile

$$
\begin{equation*}
\mu(R)=\mu_{e}+\frac{2.5 b_{n}}{\ln (10)}\left[\left(\frac{R}{r_{e}}\right)^{1 / n}-1\right], \tag{1.9}
\end{equation*}
$$

where $\mu_{e}$ is the surface brightness at $r_{e}$, i.e. the 'mean effective surface brightness'. Letting $R=0$ yields then the central surface brightness $\mu_{0}$

$$
\begin{equation*}
\mu(R=0)=\mu_{0}=\mu_{e}-\frac{2.5 b_{n}}{\ln (10)} . \tag{1.10}
\end{equation*}
$$

To obtain the mean surface brightness $\langle\mu\rangle_{e}$, defining the surface brightness within $r_{e}$ one has to consider the average intensity, $\langle I\rangle_{e}$, which is acquired by integrating the intensity over the area $A=\pi r_{e}^{2}$, such that

$$
\begin{equation*}
\langle I\rangle_{e}=\frac{\int I \mathrm{~d} A}{A}=\frac{I_{e} \mathrm{e}^{b_{n}} \int_{0}^{r_{e}} \mathrm{e}^{-b_{n}\left(R / r_{e}\right)^{1 / n}} 2 \pi R \mathrm{~d} R}{\pi r_{e}^{2}} . \tag{1.11}
\end{equation*}
$$

Applying again the same substitution $x=b_{n}\left(\frac{R}{r_{e}}\right)^{1 / n},\langle I\rangle_{e}$ can be written as

$$
\begin{equation*}
\langle I\rangle_{e}=I_{e} \frac{2 n \mathrm{e}^{b_{n}}}{b_{n}^{2 n}} \int_{0}^{b_{n}} \mathrm{e}^{-x} x^{2 n-1} \mathrm{~d} x=I_{e} f(n) . \tag{1.12}
\end{equation*}
$$

Reminding that $b_{n}$ ensures that $r_{e}$ is indeed the radius at which half of the total light is contained, one obtains

$$
\begin{equation*}
f(n)=\frac{n \mathrm{e}_{n}^{b}}{b_{n}^{2 n}} \int_{0}^{\infty} \mathrm{e}^{-x} x^{2 n-1} \mathrm{~d} x=\frac{n \mathrm{e}^{b_{n}}}{b_{n}^{2 n}} \Gamma(2 n), \tag{1.13}
\end{equation*}
$$

which finally leads to

$$
\begin{equation*}
\langle\mu\rangle_{e}=\mu_{e}-2.5 \log (f(n))=\mu_{e}-2.5 \log \left(\frac{n e^{b_{n}}}{b_{n}^{2 n}} \Gamma(2 n)\right) . \tag{1.14}
\end{equation*}
$$

Now considering again equation (1.6) and multiplying the negative logarithmic equivalent of this equation with 2.5 , one obtains the enclosed magnitude within a radius $R$

$$
\begin{equation*}
m(<R)=\mu_{e}-5 \log \left(r_{e}\right)-2.5 \log \left[2 \pi n \frac{\mathrm{e}^{b_{n}}}{\left(b_{n}\right)^{2 n}} \gamma(2 n, x)\right] \tag{1.15}
\end{equation*}
$$

This provides the opportunity to substitute equation (1.14) into equation (1.15), which yields at $R=r_{e}$

$$
\begin{equation*}
m\left(<r_{e}\right)=\langle\mu\rangle_{e}-2.5 \log \left(\pi r_{e}^{2}\right) \tag{1.16}
\end{equation*}
$$

or for the total apparent magnitude $m_{\text {tot }}$

$$
\begin{equation*}
m_{t o t}=\langle\mu\rangle_{e}-2.5 \log \left(2 \pi r_{e}^{2}\right) . \tag{1.17}
\end{equation*}
$$

### 1.3.3. Colors of galaxies

The $H S T$ and also the emerging of other powerful telescopes in the last two decades, be it in space or on the ground, opened up new chances to resolve more and also even more distant galaxies, not only in the optical, but also other wavelength ranges like UV or IR.

By using these multi-band images and fitting a Sérsic profile (see section 2.1) to the galaxies one has now the possibility to derive the fundamental structural parameters $n$ and $r_{e}$ at different wavelengths and is therefore capable to trace different stellar populations of the galaxy. Shorter, bluer wavelengths are an indicator for the distribution of recently formed young stars, optical wavelengths, which are on the redder side of the Balmer break ( $\sim 4000$ Ångström), detect a mixture of stars at different ages and at longer wavelengths only old, red and long-living stars are pinpointed. Thus it is also possible to reconstruct the evolutionary path the galaxy has taken over its lifespan and its internal components.

Indeed there have been many recent studies describing variations in $n$ and $r_{e}$ with wavelength at all redshifts suggesting the existence of different stellar populations in one galaxy, supporting the hierarchical scenario. So La Barbera et al. (2002) found for example that $r_{e}$ in NIR is $\sim 40$ percent smaller than in the optical wavelength regime for cluster early-type galaxies with $n \gtrsim 4$ at intermediate redshift ( $z=0.31$ ). Subsequently this has been confirmed by using a larger sample at different intermediate redshifts ( $z=0.21-0.64$ ) (La Barbera et al., 2003), showing that on average early-type galaxies are more concentrated at longer wavelengths, which reveals $r_{e}$ becoming smaller moving from UV to NIR, while $n$ increases.

The same $r_{e}$ evolution can be observed from local, low-redshift early-type galaxies, with $r_{e}$ decreasing $\sim 40$ percent going from g to K band (La Barbera et al., 2010; Kelvin et al., 2012), while $n$ shows no dependence on wavelength (Kennedy et al.,

2015; Vulcani et al., 2014), which confirms the assumption that spheroids can be described by a single Sérsic profile in all wavelengths.

Finally also at high redshifts $(z \sim 1.5)$ a decrease in size by approximately 20 percent going from the i to the H band can be observed (Chan et al., 2016).

However in addition to structural parameter measurements, a more thorough inspection of galaxies can be achieved by examining the colors respectively the color profiles and the change of these profiles with radius, i.e. the color gradients, disentangling the underlying physical processes associated with the evolution of passive galaxies, represented by age and metallicity gradients.

The colors of a galaxy are simply defined as the differences of the surface brightness in two different wavelengths at a certain radius, being the result of different structural parameters for the two passbands. The color profile is therefore the radial trend in color of a galaxy. Following equation (1.9) the color profile $\left(\mu_{\lambda_{1}}-\mu_{\lambda_{2}}\right)(R)$, with $\lambda_{1}$ being the bluer band results in

$$
\begin{align*}
\left(\mu_{\lambda_{1}}-\mu_{\lambda_{2}}\right)(R) & =\mu_{e, \lambda_{1}}+\frac{2.5 b_{n, \lambda_{1}}}{\ln 10}\left[\left(\frac{R}{r_{e, \lambda_{1}}}\right)^{1 / n_{\lambda_{1}}}-1\right]  \tag{1.18}\\
& -\mu_{e, \lambda_{2}}+\frac{2.5 b_{n, \lambda_{2}}}{\ln 10}\left[\left(\frac{R}{r_{e, \lambda_{2}}}\right)^{1 / n_{\lambda_{2}}}-1\right] .
\end{align*}
$$

The color gradient is then the slope of the color profile between a chosen radius interval, commonly scaled logarithmically.

$$
\begin{equation*}
\nabla_{\lambda_{1}-\lambda_{2}}=\frac{\Delta\left(\left(\mu_{\lambda_{1}}-\mu_{\lambda_{2}}\right)(R)\right)}{\Delta \log R} \tag{1.19}
\end{equation*}
$$

One common way to derive the slope is to fit a least squares fit onto the color profile with the slope of the fit being an estimate for the gradient or one of the alternative options is to do this 'analytically' following La Barbera et al. (2002), where the parameters $r_{m}$ and $r_{M}$ are the inner and outer boundaries in units of $r_{e}$ of the chosen radius interval to calculate the color gradient in.:

$$
\begin{align*}
\nabla_{\lambda_{1}-\lambda_{2}} & =\frac{2.5 \log (\mathrm{e})}{\log \left(r_{M}\right)-\log \left(r_{m}\right)} \\
& \times\left\{b_{n, \lambda_{2}}\left[\left(\frac{r_{m}}{r_{e, \lambda_{2}}}\right)^{1 / n_{\lambda_{2}}}-\left(\frac{r_{M}}{r_{e, \lambda_{2}}}\right)^{1 / n_{\lambda_{2}}}\right]\right.  \tag{1.20}\\
& \left.+b_{n, \lambda_{1}}\left[\left(\frac{r_{M}}{r_{e, \lambda_{1}}}\right)^{1 / n_{\lambda_{1}}}-\left(\frac{r_{m}}{r_{e, \lambda_{1}}}\right)^{1 / n_{\lambda_{1}}}\right]\right\} .
\end{align*}
$$

Using now simulated models describing the formation and developmental history of galaxies (see section 1.4), one is in the position to decipher, which processes and formation scenarios are responsible for and matching the observed colors and color gradients best, as the assembly history is imprinted in its stellar populations, with colors reflecting the time and conditions of the star formation, whereas the spatial distribution gives insights into the structures in which those stars were formed or accreted.

Considering the two major formation scenarios, the monolithic model and the hierarchical scheme, it becomes clear, that both models prognosticate different color gradients.

So the monolithic dissipative model predicts steep metallicity and constant or mildly positive age gradients, with stars in the center of a galaxy being younger and therefore bluer than the average population in the outskirts. The reason for those appearances lies in the underlying physics of the scenario itself. During the gravitational collapse, so the formation of the galaxy, gas flows towards the center, which leads to a more metal-rich population in the center than in the external regions, characterized by a negative metallicity gradient. Additionally this condition is enhanced due to the central deeper gravitational potential well, causing gas to keep flowing inwards, triggering a star formation and hence chemical enrichment, which is lasting longer in the center than in the outer regions, which is the reason for a null or positive age gradients, but a steep negative metallicity gradient. Thus the latter should dominate the color profile revealing a strong negative color gradient with values greater than $-1 \mathrm{mag} \mathrm{dex}^{-1}$ in radius for the classical and -0.5 to -0.3 mag dex ${ }^{-1}$ for the revised monolithic model (Montes et al., 2014, and references therein).

Applying the hierarchical dry merger and also the inside-out-growth scenarios on the other hand will result in much shallower gradients, due to a dilution of pre-
existing gradients as a result of the merger(s). The end result would depend on whether the stellar populations were mixed, according to the hierarchical model or the accreted stellar content just being added as an envelope to the existing galaxy, following the preferred inside-out-scheme.

Color gradients have been measured already almost 30 years ago (e.g. Vader et al., 1988), even then supporting negative color gradients, which are corresponding to a decrease in metallicity with radius for local, early-type galaxies (Peletier et al., 1990a,b). Using data from the extensive $S D S S$ survey (La Barbera \& de Carvalho, 2009) substantiated this picture of negative, metallicity-driven gradients, but also found a small positive age gradient implying early-type galaxies having smaller $r_{e}$ at higher redshifts (Montes et al., 2014; Wu et al., 2005).

With the advancement in technology newer studies followed allowing to expand the color gradient measurements to more distant galaxies. Those investigations confirmed metallicity as the dominant driver for color gradients for passive galaxies in the field, as well as in clusters, with the origin of the color gradients being independent of the respective galaxy environment (Saglia et al., 2000; Smail et al., 2001; Tamura \& Ohta, 2000; Tamura et al., 2000). La Barbera et al. (2003) found a mean UV-optical color gradient of $-0.18 \pm 0.04 \mathrm{mag} \mathrm{dex}^{-1}$, respectively a mean opticalNIR color gradient of $-0.4 \pm 0.1 \mathrm{mag} \mathrm{dex}^{-1}$ for passive galaxies at intermediate redshifts ( $z \sim 0.21-0.64$ ).

At high redshifts $(z \sim 1-2)$ the color gradients of elliptical galaxies appear significantly steeper with the values reaching twice the local values (Gargiulo et al., 2011, 2012; Guo et al., 2011), implying an important contribution by an age gradient at such high redshifts (Chan et al., 2016). The reason for that behaviour may be justified with the galaxies just reaching the phase of passive evolution (De Propris et al., 2015). In the following $\sim 8-9$ Gyrs the galaxies evolve and the age gradient tends to null, as the young, blue stars vanish and only the old, red, long-living stars remain in every region of the galaxies, with the metallicity gradient prevailing to be the reason for the color gradients at $z \sim 0$.

All these results regarding the evolution of the color gradients over cosmic time in combination with the size evolution and dependency on wavelength give reason to discount the monolithic scenario and to endorse the hierarchical/inside-out-growth model, which is in agreement with chemodynamical simulations (Kobayashi, 2004).

### 1.4. Stellar population synthesis

Since most galaxies can not be resolved into single stars one must resort to the use of independent stellar population diagnostics in form of evolutionary population synthesis (EPS) modelling to derive definite or at least well estimated values for the ages and elemental abundances which are correlating to the observed and measurable colors and their gradients. Once cosmological parameters, such as $H_{0}, \Omega_{\Lambda}$ and $\Omega_{M}$ are chosen, ages and metallicities derived from local observations can be used to determine the changing properties of galaxies over cosmic time (see Maraston, 2003).

The first works in this field happened already around five decades ago with the perception that colors due to their sensitivity to age can be applied to date extragalactic objects (Crampin \& Hoyle, 1961), followed by the comprehensive works of Beatrice Tinsley (e.g. Tinsley, 1972), who established evolutionary population synthesis as a tool to model spectroscopic and photometric properties of galaxies. One respectively two decades later the fuel consumption theorem and the isochrone synthesis technique arose, two methods to compute the 'simplest' form of an EPS, the Simple Stellar Population (SSP) models.

The underlying assumption of these models is that all stars are formed at one epoch and hence display the same chemical composition and age. Since for globular clusters belonging to our Milky Way, ages and metallicities are independently well known, those clusters can be used to calibrate SSP models to examine extragalactic stellar properties. The fuel consumption theorem and the isochrone synthesis technique differ according to the integration variable selected in the post main sequence (PMS) phase (see Maraston, 2005). With the first method the integration variable is represented by the 'fuel', serving as a parameter for the amount of hydrogen and/or helium, which is used during the PMS phase, whereas with the isochrone synthesis the integration is done over the contributions in all wavelengths to the overall emission from all mass bins along one isochrone. Both approaches deliver similar results, although models still sometimes depend on insufficient spectral libraries, especially for non-solar computations. The only distinction in the outcome affects ages $\sim 1$ Gyr, based on the complicated process of mass loss of stars in advanced stages, like in the Asymptotic Giant Branch (AGB) phase, which shows models based on the fuel theorem calculations displaying higher fluxes in the NIR (Renzini, 2006).

Besides the stellar evolutionary tracks and stellar model atmospheres, with the former describing how stars with different masses and chemical compositions and
therefore their stellar parameters evolve with time and the latter determining the resulting flux, the initial mass function (IMF) is a key component for the modelling of stellar populations. This (empirical) function is an assessment of the distribution of the initial mass for a stellar population and is usually expressed in the form of a power-law with single (e.g. Salpeter, 1955) or multiple slopes (e.g. Chabrier, 2003; Kroupa, 2001). For the calculations of colors however the choice of the IMF plays a negligible role, as the majority of the emission stems from stars in a narrow mass bin, residing near the main sequence turn-off point (Renzini, 2006).

What must be considered by all means on the other hand is the degeneracy of age and metallicity, with a reduction in age and an increase in metallicity leading to a very similar spectral energy distribution in the optical range (Worthey, 1994), which can be quantified as

$$
\begin{equation*}
\frac{\Delta \log (t)}{\Delta \log (Z)} \sim \frac{3}{2} \tag{1.21}
\end{equation*}
$$

One photometric approach to break this degeneracy is to compare one's measurement in two different colors, one being an optical color, whereas the second color is a combination of an optical and an IR band. The first one is almost equally sensitive to age and metallicity, while the second color traces the old stellar population ( $\gtrsim 1 \mathrm{Gyr}$ ), which is much more perceptive to changes in metallicity than age. Using bluer optical colors, including maybe even the NUV range, would result in a setback of breaking the degeneracy, as such colors are more sensitive to the metallicity, because of enclosing the 4000 Ångstrom break (Smail et al., 2001).

## 2. Tools \& Software

In this chapter the tools are presented, which have been used to derive results outlined in the course of this thesis. The first section describes software to detect, measure and extract the needed data. The second part outlines the means used to model the evolutionary path of galaxies using SSPs.

### 2.1. Parametric fitting: Source Extractor, GALFIT and Megamorph

### 2.1.1. Source Extractor

With the emergence of extensive surveys the need arose to detect, examine, measure and catalogue the final data products in an automatic fashion. One of today's most used software to do this task is called Source Extractor (SExtractor) (Bertin \& Arnouts, 1996).
Given an input file, where the user can define important parameters, such as detection conditions, background measurements constraints or photometric apertures and zeropoints (see appendix for an example file, used in this work), the programme is capable of automatically detecting, deblending, measuring and classifying objects, including robust sky background measurements. It also provides a very reliable separation of galaxies and stars, as well as, if the user wishes, inspection images, with the detected objects being masked out or depicted with apertures around them.
The order of steps performed start with an estimation of the sky background, followed by thresholding and deblending, proceeding to filtering detections, executing measurements, separating galaxies and stars and finally of course writing all these results into a catalogue with optional check fits-files as well.

Due to its large automation this software is optimally suited for the analysis of comprehensive extragalactic surveys, as the evaluations include amongst other results photometric assessments in user defined apertures and determinations on positions, position angles and axis ratios.

### 2.1.2. GALFIT

Applying apertures directly on an image to derive photometric measurements is always a challenging method as the probability is very high that either some flux gets missed or the sources' emission gets overestimated due to the contribution of the background as well as maybe neighbouring sources. Both is caused by the fact that the extent of a galaxy is rarely known, which is especially valid for early-type galaxies, due to their intrinsic light profile, which leads to an improper choice of apertures. The whole effort gets even more complicated, if one wants to compare the same object in different wavelengths, as the extent of the source can and will presumably vary, as well as the $\mathrm{S} / \mathrm{N}$. A solution of course would be to chose one or more apertures in one band and apply them to all images, showing the source in different passbands. Unfortunately the fluxes will be under- or overestimated, depending on the wavelength in which the aperture has been chosen. The same conclusions can be drawn for size measurements, as they depend on the photometric results.

One way to evade such obstacles would be to classify galaxies with parameters, defined as a ratio of two radii containing a portion of the total emission. A widely used example would be the parameters summarized in the $C A S$-system (Conselice, 2003), describing the compactness $(C)$, asymmetry $(A)$ and clumpiness $(S)$ of a galaxy. For example the compactness $C$ is defined as the ratio of the radii containing 20 and 80 percent of the total flux, $C=5 \log \left(r_{80} / r_{20}\right)$.

Nevertheless as also for these derivations and similar non-parametric methods the photometric outcomes are vital, these estimations suffer from the same biases found in aperture photometry. A completely alternative approach to avoid such problems is to fit the 2 D surface brightness profiles of galaxies with parametric models. Two widely spread tools to do this task are GIM2D (Simard et al., 2002) and GALFIT (Peng et al., 2002, 2010). The advantages of such a strategy is that all measurements are done in a consistent way and especially that the problem of missing flux is greatly minimized, due to the possibility to extrapolate the best-fitting model to large radii. The disadvantage obviously is that a specific model (or several) has to be assumed a priori for the observed light profile.

As GIM2D results show significant higher errors in the modelling of spheroidal galaxies with close neighbouring objects, due to the inability to fit the main galaxy as well as the additional sources simultaneously, and is just rather masking the companions, which leads to overestimated values for $n$ and is also showing a need
for more computational time (Häussler et al., 2007), GALFIT has been chosen to model the galaxies analysed in this thesis.

Given a simple input file GALFIT has the ability to extract structural parameters like the total magnitude, $r_{e}$ and $n$ via finding the best fit of a predetermined light profile to the observed one. The selection of possible 2D models include Nuker, Moffat or Sérsic profiles. The latter was used in this thesis, including the special types, Gaussian ( $n=0.5$ ), exponential $(n=1)$ and de Vaucouleurs $(n=4)$. One of the features GALFIT provides is that also a combination of profiles can be easily achieved. Consequently it is possible to specify in the input file that a source should comprise two components, one showing a de Vaucouleurs' profile, the second an exponential one, which would be a typical composition for late-type galaxies. Since only early-type galaxies are considered in this thesis, only a one-component Sérsic profile fit has been applied to the galaxies.

To find the parameters which are describing the best-fitting model GALFIT first convolves the model profile with a, from the user defined, point spread function (PSF), multiplying the Fourier transforms of the PSF and the models followed by an inverse transformation of them. Subsequently the convolved profiles are compared with the observed data and $\chi^{2}$ minimized using the Levenberg-Marquardt algorithm. This process is iterated until convergence is reached.

The output parameters for a Sérsic profile fit are $x$ and $y$, denoting the position of the source in pixels, the total magnitude, converted from the flux using the exposure time given in the FITS header of the image file and the profile being integrated to $R=\infty$, the effective radius $r_{e}$, the Sérsic index $n$, the axial ratio between major and minor semiaxis $Q$ and the position angle $P A$. Also uncertainties for those parameters are given but they should only be recognized as lower limits as these errors only reflect the fluctuations in the residual image induced by Poisson noise (Peng et al., 2010). This aspect is mostly valid only for idealized situations, such as modelling image simulations. For real images the errors rather stem from structures within a galaxy, which have not been modelled, like spiral arms or bars.

Additional to the output catalogue multi-frame fits files are provided, including the original image, the model and the residual, which offers a visual verification of the quality of the fit.

### 2.1.3. Megamorph

Data products from big surveys provide information on a vast number of sources multiplied by the number of observed passbands. Of course applications and algorithms like SExtractor and GALFIT are simplifying the efforts to extract valuable scientific information out of the data. Nevertheless it would be a tedious, very time consuming and due to the sheer number count almost not manageable work to analyse every single object in every filter. Hence for the specific case regarding SExtractor and GALFIT Barden et al. (2012) developed a framework for users to - once everything has been set up - execute SExtractor and GALFIT in an automatic fashion, called Galapagos ('Galaxy Analysis over Large Areas: Parameter Assessment by GALFITing Objects from SExtractor').

As a result this wrapper is not only capable of executing first SExtractor and subsequently using these preliminary outcomes regarding structural parameters as input variables for GALFIT, but it also includes some optimizing modifications in regard to the analysis of large images with sources in a wide magnitude range. So the detection of objects with SExtractor is done in a two-way approach according to two by the user defined input files. At first a so-called 'cold' mode is performed to detect bright sources, followed by a 'hot' mode to detect the more numerous faint objects. To do this in one SExtractor run would not be possible as either the bright sources would be split up and detected as several individual ones or the fainter objects would simply not be detected. To estimate the extent of a galaxy the Kron radius (Kron, 1980) is used, which as a consequence means that a 'hot', so faint source is only included to the final catalogue of objects, if it is not included in any other Kron ellipse.

Furthermore Galapagos has the ability to create postage stamp images of the objects to minimize the time and memory needed for the fitting process and eventually also possesses an own sky estimation routine. This is done by averaging the background flux in elliptical annuli, which are centered on the respective object, while it ignores other sources or image defects contaminating the flux in the respective annuli. Hence the background flux can be expressed as a function of radius and the definitive sky value is measured at the position where the profile's slope turns positive the second time.

Finally Galapagos can also judge if an object can be masked in the fitting process or must be fitted simultaneously with the primary source. If the Kron ellipse of an object overlaps with the one of the desired galaxy, this neighbour will be fitted as
well, as it may have a significant influence on the light profile of the primary source. If not it will be masked and will not be considered in the fitting process.

The output is finally summarized in a catalogue listing all detected and fitted sources with their modelled parameters including uncertainties derived with GALFIT. Furthermore also the usual multi-frame fits files are delivered, depicting the original image, the model and the residual.

It should be noted that, while the original version of Galapagos needs an $I D L$ runtime environment, Hiemer et al. (2014) developed a version, which can run in $C$ and has been optimized in regard of computational speed and efficiency, as it is capable to parallelize the processes, which is a much needed feature for using supercomputers.

The big disadvantage of Galapagos is that, although it can be utilized for multiband images, each of the passbands have to be considered separately. Therefore a lot of information gets lost, which holds especially true for images in a wavelength with a low $\mathrm{S} / \mathrm{N}$, or also clues are missed concerning the wavelength dependence of the flux of a source. To make use of this valuable information Galapagos has been further developed into MegaMorph (Bamford et al., 2012; Häußler et al., 2013).

To produce such an application a modified version of GALFIT has been designed, GALFIT-M, which is able to accept multiple input images and for which the model parameters are replaced by wavelength dependent Chebyshev polynomials, with the Chebyshev coefficients being the fitted parameters in the respective passbands. The degree of the polynomials may be defined by the user, ranging from the parameters being constant over all wavelengths to being completely independent. Hence the user can decide to which extent the values for each fitted parameter are connected to each other.

The required input file(s) (see appendix for an example) are very similar to the ones needed to run the original version of Galapagos. The only differences are obviously that now images for all bands must be declared and that additionally on one hand one must define one extra band for the SExtractor run, which could also be a coadded image to increase the number of detections, and on the other hand one image for deblending and masking.

Also the efficiency has been improved, which results in a factor 4 of less needed CPU time. This has been partially achieved by refining the loop routine. In the original version of Galapagos the algorithm postpones the fitting procedure for the next object, if it is contaminated by the object, which is currently fitted. MegaMorph
avoids this by skipping this particular source and returning to it as soon as it is safe to fit it.

The most improvement with modelling the galaxy parameters as a function of wavelength is that the user can select a balance between consistency and flexibility. With well chosen input parameters images in bands with low $\mathrm{S} / \mathrm{N}$ can profit from constraints of images with higher $\mathrm{S} / \mathrm{N}$. This leads to more robust measurements, which are conveyed with more useful magnitude assessments for faint galaxies, as well as improved results for sizes and $n$, especially in comparison to single-band fits. Hence it is possible to conduct a more precise analysis of colors and color gradients within a galaxy, which is particularly of great importance for this thesis.

### 2.2. Stellar population modeling: EzGal

To be in the position to compare the observed parameters like magnitudes and colors to the modelled ages predicted by SSP calculations, the Python wrapper EzGal (Mancone \& Gonzalez, 2012) has been used in the course of this thesis.

It offers the possibilities to calculate the evolution of a stellar population over redshift, based on different models (e.g. BC03, Bruzual \& Charlot (2003) or M05, Maraston (2005)) and varying initial conditions in regard to star formation history, IMF, metallicity or formation redshift. Moreover it is also possible to interpolate between metallicities for a certain set of models. The included filter set contains passbands ranging from the UV, such as GALEX filters to optical ones, used by $H S T$ or SDSS to IR bands, utilized for example by Spitzer or WISE. However the addition of extra filters poses no problem as the user can simply add filters by providing EzGal the respective response curves. Magnitudes can be produced in the Vega and AB photometric system, the latter being deployed in this thesis.

A comparison of the synthesis results of an SED evolution between the different models shows that there is a good agreement in the optical wavelength range for old models possessing solar metallicities, with a difference of about 0.1 mag. Turning however to younger ages ( $<2 \mathrm{Gyr}$ ) and longer wavelengths ( $>7500$ Ångstrom) the differences increase significantly, presumably due to thermally pulsating AGB stars (Mancone \& Gonzalez, 2012), which are differently treated in the various models.

EzGal derives the magnitude evolution over redshift from models, which are describing the SED evolution over age. So based on the default or user defined cosmological parameters, besides the absolute and apparent magnitudes in rest-frame,

EzGal can also derive those magnitudes in the observed frame, which are calculated using the redshift-dependent $k$-correction (Hogg et al., 2002a). Hence the measured apparent magnitudes in a certain wavelength and colors of the observed galaxies can directly be compared to the modelled values in the same filters at a certain redshift, which would be for this particular analysis the respective cluster redshifts. This in turn makes it possible to infer the corresponding ages and metallicities of the stellar population residing in the galaxies, after careful consideration of the age-metallicity degeneracy.
The concordance cosmology is adopted in this thesis with $H_{0}=70 \mathrm{~km} \mathrm{~s}^{-1} \mathrm{Mpc}^{-1}$, $\Omega_{\Lambda}=0.7$ and $\Omega_{M}=0.3$ (Spergel et al., 2003).

## 3. Data

### 3.1. CLASH

The major part of the data used in this thesis is provided by the 524 orbit HST Multi-Cycle Treasury Program The Cluster Lensing And Supernova survey with Hubble (CLASH) (Postman et al., 2012). By observing and examining 25 massive galaxy clusters with 16 HST filters ranging from near-UV to the near-IR four main goals should be achieved with this program:

- The derivation of accurate measurements of the mass distribution and substructures of dark matter in those clusters. These are obtained by combining the $H S T$ measurements with weak lensing maps derived from wide-field Subaru images produced by Suprime-Cam in the bands $B, V, R c, I$ and $z^{\prime}$, as well as mm and X-ray images. This wide spectrum of wavelengths, as well as the large field of view allows constraining the mass profiles of the clusters out to large radii, including precise assessments of the central density concentrations together with constraints in regard of the dark matter substructure distribution derived by galaxy-scale lensing.
- The detection of SN Ia out to a redshift $z \sim 2.5$ in order to measure the time dependence of the dark energy equation as well as a possible redshift dependence in the SNe themselves. To accomplish this goal the observations of each cluster are stretched over eight epochs with each cluster being observed at two orientations, which are approximately $30^{\circ}$ apart to reduce the overlap between the parallel ACS pointings (Fig. 3.1). At one point ACS observes the respective cluster core, whereas WFC3 views the surrounding parallel fields followed by a reversed observational setup. Thus the observations of the clusters and the search for SNe are done in a parallel mode.
- The utilization of the cluster's lensing capabilities to identify high redshift ( $z>7$ ) galaxies.
- The gain of additional information about the internal structures of the clusters along with clues regarding the evolution and therefore also possible formation scenarios of the cluster galaxies, as well as galaxies lying behind those clusters in redshift space. This can be achieved by utilizing the 16 different bands, resolving besides possible particular features especially different stellar populations. Hence with this knowledge of the internal structures of galaxies their evolution can be inferred.


Figure 3.1.: Fields covered by the CLASH observations, which are done twice at two different orientations $\sim$ $30^{\circ}$ apart. While the cluster core is observed with ACS, the parallel fields are observed with WFC3 and vice versa. Taken from (Postman et al., 2012).

All 25 clusters in the sample are massive, with masses of $5 \times 10^{14}<$ $M_{v i r} / M_{\odot}<3 \times 10^{15}$, with the clustercentric redshifts being in the range 0.18 $\lesssim z \lesssim 0.9$ with a median redshift $z \sim$ 0.4 . Out of those 25 clusters 20 were exclusively X-ray selected, with $k T>$ 5 keV , and five additional clusters were included solely due to their high lensing strength, revealing Einstein radii up to $55^{\prime \prime}$, to enhance the likelihood of finding high redshift galaxies. Yet at least in 18 of the 20 X-ray selected clusters one or more giant arcs are visible, corresponding to Einstein radii in the range $15^{\prime \prime}$ to $30^{\prime \prime}$. The supplementary five clusters are not necessarily relaxed, which is definitely the case for one of the clusters discussed in this thesis, MACS J0416.1-2403.

However most clusters of the X-ray selected sample appear as dynamically relaxed structures as can be confirmed by images done with the Chandra X-ray Observatory. They exhibit a well-defined central surface brightness peak and nearly concentric isophotes. Furthermore these observations are supported by X-ray pressure maps, showing almost no deviation from a hydrostatic equilibrium. In fact 14 clusters, including those two, which are analysed in this thesis, were initially detected by the X-ray cluster survey MAssive Cluster Survey (MACS) (Ebeling et al., 2001), intended to compile a large, unbiased sample of X-ray luminous, distant clusters with luminosities $L_{X} \gtrsim 5 \times 10^{44} \mathrm{erg} \mathrm{s}^{-1}(0.1-2.4 \mathrm{keV})$ and redshifts $z \gtrsim 0.3$.

To accomplish the ambitious goals of CLASH a total of 16 broadband filters (Fig. 3.2), ranging from near-UV to near-IR ( $\sim 2000-17000 \AA$ ) using the $A C S(W F C)$
for its greater throughput efficiency, especially regarding the red wavelength regime, and WFC3 (UVIS/IR) detectors are deployed. The region covered by the ACS filters spans $4.08 \mathrm{arcmin}^{2}$, with the $W F C 3 / I R$ field of view being 88 percent of this area.


Figure 3.2.: Response curves for the filters used in the CLASH survey ranging from nearUV to near-IR. Just for clarity some filter curves show an offset of 0.2 (dashed line). Plot taken from (Postman et al., 2012).

Table 3.1 lists the central filter wavelengths as well as the limiting magnitudes in the AB photometric system (Oke (1974); Oke \& Gunn (1983)) for extended sources at $5 \sigma$, which are the mean limits averaged over all extended sources (i.e. galaxies) not considering the sizes of the respective objects (Jouvel et al., 2014).

Table 3.1.: The central wavelengths of the HST filter set used in the CLASH survey and the respective $5 \sigma$ magnitude limits for extended sources. (Magnitude limits taken from (Jouvel et al., 2014))

| Detector | Filter | $\lambda_{c}$ <br> $[\AA]$ | $(5 \sigma)$ <br> $[\mathrm{AB} \mathrm{mag}]$ | Detector | Filter | $\lambda_{c}$ <br> $[\AA]$ | $(5 \sigma)$ <br> $[\mathrm{AB} \mathrm{mag}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| WFC3/UVIS | F225W | 2359 | 24.8 | ACS/WFC | F775W | 7764 | 26.5 |
| WFC3/UVIS | F275W | 2704 | 24.9 | ACS/WFC | F814W | 8333 | 27.5 |
| WFC3/UVIS | F336W | 3355 | 25.2 | ACS/WFC | F850LP | 9445 | 26.2 |
| WFC3/UVIS | F390W | 3921 | 26.1 | WFC3/IR | F105W | 10552 | 27.0 |
| ACS/WFC | F435W | 4297 | 26.2 | WFC3/IR | F110W | 11534 | 27.6 |
| ACS/WFC | F475W | 4760 | 26.7 | WFC3/IR | F125W | 12486 | 27.1 |
| ACS/WFC | F606W | 5907 | 27.2 | WFC3/IR | F140W | 13923 | 27.3 |
| ACS/WFC | F625W | 6318 | 26.6 | WFC3/IR | F160W | 15369 | 27.2 |

The images were subsequently cleaned, processed, aligned and co-added using the MosaicDrizzle software (Koekemoer et al., 2011) fixing for cosmic ray contamination and working out shifts and rotations. Additionally the exposures were corrected for galactic extinction using IR dust emission maps (Schlegel et al., 1998). The images have a scale of $0.065^{\prime \prime}$ per pixel with north being upwards and are normalized to counts s ${ }^{-1}$.

With this abundance of photometric measurements it is possible to derive precise and reliable photometric redshift measurements. This has been done using Le Phare as well as BPZ packages, with the result that both codes show a similar assessment of the photometric redshifts for galaxies in the cluster as well as in the foreground (Jouvel et al., 2014).

Finally it must be noted, that for this thesis the WFC3/UVIS bands were discarded, due to their low $\mathrm{S} / \mathrm{N}$, which left 12 bands for the analysis.

### 3.2. CLASH-VLT

CLASH-VLT is an ESO Large Program, titled 'Dark Matter Mass Distributions of Hubble Treasury Clusters and The Foundations of $\Lambda$ CDM Structure Formation Models' (Rosati et al., 2014) and a follow-up of the CLASH survey using the VIsible MultiObject Spectrograph (VIMOS) instrument (Le Fèvre et al., 2003) at the ESOVLT on 13 clusters already observed photometrically with CLASH at $z \sim 0.2-0.6$ (median $z \sim 0.4$ ).

The intention has been to obtain spectroscopic data to be able to identify 500 to 1000 member galaxies for each cluster and over 200 background galaxies with $z \lesssim 7$. This grants the opportunity to determine the cluster's mass density profiles by dynamical analysis of the cluster members and to compare and combine those results with independent mass profile measurements obtained from strong and weak gravitational lensing using Subaru Suprime-Cam images or X-ray studies out to a radius of approximately 3 Mpc . Furthermore this dataset also offers the possibility to probe the composition of the inner dark matter structure of the cluster halos.

Twelve masks were utilized, eight with a low resolution ( $R=180$ ) blue grism between $370-670 \mathrm{~nm}$ and four with medium resolution ( $R=580$ ) grism in the range of $480-1000 \mathrm{~nm}$ for high-redshifted lensed sources. The observations were done with 8-12 separate pointings for each cluster, with one quadrant being fixed focusing on the respective cluster cores to increase the exposure time in order to
detect also the faintest lensed objects. The field of view covers $15-20 \operatorname{arcmin}^{2}$, which corresponds to $\sim 10 \mathrm{Mpc}$ at $z \sim 0.4$. The total exposure time amounts to approximately 200 hours, which leads to the final dataset for all clusters consisting of $\sim 30000$ spectra with $\sim 7000$ being cluster members.

The obtained dataset has been reduced using the VIMOS Interactive Pipeline and Graphical Interface (VIPGI).

### 3.3. MACS 1206

The primary object and first cluster of study in this thesis is MACS J1206.2-0847 (MACS 1206 in the following). Inferring from X-ray observations revealing an X-ray luminosity of $L_{X} \sim 2.4 \times 10^{45} \mathrm{erg} \mathrm{s}^{-1}$, the cluster appears to be in a relaxed state, with the central BCG at $\mathrm{RA}_{2000}=12^{h} 06^{m} 12^{s} .15$ and $\mathrm{Dec}_{2000}=-8^{\circ} 48^{\prime} 3 .^{\prime \prime} 48$ located at the same position as the peak of the X-ray emission and the, from gravitational lensing determined mass center (Umetsu et al., 2012). Furthermore the mass profiles derived on one hand from dynamical and on the other hand from strong- and weaklensing methods (Zitrin et al., 2012) are in excellent agreement, with MACS 1206 revealing a virial mass of $M_{200} \sim 1.4 \times 10^{15} M_{\odot}$ and a virial radius of $r_{200}=1.98 \mathrm{Mpc}$ (Biviano et al., 2013), with the ' 200 ' subscript referring to the radius where the average density equals 200 times the cosmic mean density at the cluster redshift.

Nevertheless MACS 1206 exhibits a significant, not centrally concentrated, WNWESE elongated intracluster light (ICL) component, indicating that interactions between galaxies are still ongoing, resulting in tidal disruptions feeding the ICL, despite the overall relaxed appearance (Eichner et al., 2013). This assessment can be supported by the examination of the stellar mass profile, with the ratio of giant $\left(M_{*} / M_{\odot}>10^{10.5}\right)$ to subgiant galaxies $\left(10^{9.5}<M_{*} / M_{\odot}<10^{10.5}\right)$ being at the highest in the innermost regions ( $>0.5 \mathrm{Mpc}$ ) just to drop to its minimum in the adjoining region ( $0.5-1 \mathrm{Mpc}$ ) (Annunziatella et al., 2014). Also the presence of red galaxies showing a strong $\mathrm{H} \delta$ absorption, indicating a recent star formation event $1-2$ Gyr ago, in the center and along the ICL, implies recent or even ongoing interactions on galaxy scales (Girardi et al., 2015; Mercurio et al., 2015; Presotto et al., 2014).

### 3.3.1. Selection

The initial selection of the member galaxies (Fig 3.4) of MACS 1206 is based on the analysis and identification by Biviano et al. (2013).

Using the spectroscopic results obtained by the CLASH-VLT VIMOS


Figure 3.3.: Redshift distribution of the initial selection for MACS 1206. Red shows galaxies with spectroscopic redshift determinations, yellow the ones with photometric redshift estimates. The gray dashed line indicates the cluster redshift at $z \sim 0.44$ program it was possible to determine 600 member galaxies in the Subaru field belonging to the cluster via the analysis of their locations in the projected phase space, $R$ and $v_{r f}$ and applying a 'clean' as well as for comparison a 'peak + gap' algorithm.
Aside from that, by utilizing Subaru's Suprime-Cam's 5 bands additional photometric redshifts have been derived. Possible cluster membership of galaxies with just photometric but no spectroscopic redshift determinations has been investigated by using a comparison between galaxies possessing $z_{\text {spec }}$ and $z_{\text {phot }}$ measurements as well as color-color cuts in $R_{C}-I_{C}$ vs. $B-V$ (see also Annunziatella et al., 2014), which yields a redshift range of $0.34<z_{p h o t}<0.54$ for galaxies to be in to be defined as cluster members.
Using these membership determinations, combined with the photometric Le Phare redshift determinations by Jouvel et al. (2014) and the SExtractor catalogue for the ACS F814W band provided by the CLASH collaboration ${ }^{1}$, resulted in an initial selection of galaxies visible in the HST footprint for the analysis of MACS 1206.

This selection contains 216 sources, but as the field of view of WFC3/IR is significant smaller, all galaxies, which are just being detected in the ACS field of view have been discarded to optimize the accuracy of the output of the modeling process.

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Figure 3.4.: MACS 1206 in the WFC3 F105W band, depicting the initial selection of galaxies to be fitted, distinguished in a sample with spectroscopic redshift determinations (red circles) and photometric redshift estimates (yellow circles).

Doing this cut leaves a total of 110 galaxies (Fig. 3.4) for the fitting process, with 74 possessing a spectroscopic derived redshift, while 36 only have photometric redshift determinations. Fig. 3.3 depicts the redshift distributions of those galaxies, with red showing the spectroscopic and yellow the photometric determined member galaxies. As expected the peak of the distribution is at the cluster redshift of $z \sim 0.44$ indicated by the gray dashed line.

As the field of view of WFC3/IR encompasses less than half of the virial radius $r_{200}$ the galaxies are well located in the innermost regions of the cluster and are evenly distributed as can be seen in Fig. 3.4. This should avoid any biases, i.e. contamination by the ICL, orientation or location in the cluster. At the cluster redshift one arcsecond corresponds to $\sim 5.67 \mathrm{kpc}$.

### 3.4. MACS 0416

The second cluster examined in this thesis and being one of the five particular clusters in the CLASH survey selected due to its high magnification abilities is MACS J0416.1-2403 (MACS 0416).


Figure 3.5.: Distribution of sample galaxies in redshift for MACS 0416. The double peak structure is clearly visible. The cluster redshift $z \sim 0.397$ is depicted as the gray dashed line.

Lying at a redshift z $\sim 0.397$, this system shows a complex structure revealing two subclusters with two BCGs separated by $\sim 250 \mathrm{kpc}$ (Ogrean et al., 2015). Despite the clearly, in X-ray visible unrelaxed state of the cluster, the mass profiles derived from galaxy dynamics are in good agreement with the results obtained by strong and weak lensing, as well as X-ray measurements, displaying an X-ray luminosity of $L_{X} \sim$ $7.4 \times 10^{44} \mathrm{erg} \mathrm{s}^{-1}$. The overall virial mass amounts to $M_{200} \sim 0.9 \times 10^{15} M_{\odot}$, the virial radius to $r_{200}=1.82 \pm$ 0.11 Mpc (Balestra et al., 2015). With the complex dynamical state of the cluster but its dark matter and gas component being well aligned, it can be assumed that the cluster is in a premerging phase. The mass center appears to lie nearly halfway between the two BCGs, nevertheless, following Balestra et al. (2015) the cluster center is adopted to be the NE-BCG at $\mathrm{RA}_{2000}=04^{h} 16^{m} 09^{s} .14$ and $\mathrm{Dec}_{2000}=$ $-24^{\circ} 04^{\prime} 03 . .^{\prime \prime} 1$ coinciding with the peak position of the X-ray emission.

### 3.4.1. Selection

As is the case for MACS 1206 also for MACS 0416 the basis of the membership determinations for galaxies are the results of the $C L A S H-V L T$ program.

After comparing the observed spectra with template spectra using the software
$E Z$ and visual inspection of outliers resulting in a final dataset of 4386 redshift determined sources, distributed over the whole Subaru field, Balestra et al. (2015), following the approach of Biviano et al. (2013) for MACS 1206, applied the 'peak + gap' method, which yields two overlapping peaks in redshift distribution at $z \sim 0.396$ and $z \sim 0.400$, confirming the complex structure of the cluster.


Figure 3.6.: Initial sample selection for MACS 0416. Spectroscopic members are shown with red circles, galaxies with photometric redshifts with yellow circles. The NE BCG, also possessing a spectroscopic redshift determination, is indicated by a cyan circle.

Changing the position of the cluster center from the location of the NE BCG to the position of the barycenter $\left(\mathrm{RA}_{2000}=04^{h} 16^{m} 08^{s} .60\right.$ and $\left.\mathrm{Dec}_{2000}=-24^{\circ} 04^{\prime} 25 .^{\prime \prime} 2\right)$ leads to a negligible impact on the selection of galaxies belonging to the cluster, as only galaxies at radii larger than $5^{\prime}$ are affected. Hence this factor can be ignored in
this thesis, as only galaxies within the $H S T$ field of view, which corresponds to less than half of $r_{200}$, are considered.

Finally employing a 'simple isothermal sphere' model with the center of the cluster being the position of the NE BCG results into a final redshift range of 0.381 $\lesssim z \lesssim 0.411$. Combining these constraints with SExtractor catalogues including photometric redshifts (Jouvel et al., 2014) from the CLASH pipeline results into a preliminary selection of 179 galaxies, with 134 possessing spectroscopic data and 45 only photometric redshifts.

Eliminating the galaxies, which are only visible in the ACS field leads to an initial sample containing 92 galaxies, 81 of them having spectroscopic derived redshifts, while 11 , mostly very faint sources, got only photometric measurements.

The collection of galaxies chosen to be modelled are shown in Fig. 3.6, with sources surrounded by a red circle including the NE BCG indicated by a cyan circle being spectroscopic measured, whereas the yellow circles denote the photometric assessed galaxies. Evident is the cluster's orientation in NE-SW direction, as well as the diffuse intra cluster-light following this elongation.

The redshift distribution of the dataset is depicted in Fig. 3.5, illustrating the apparent double peak in redshift-space, with the gray dashed line indicating the cluster redshift at $z \sim 0.397$. The color coding is still the same, for galaxies with spectroscopic measurements being represented by the red bars and photometric ones by yellow bars. Due to the slightly lesser distance of MACS 0416 in comparison to MACS 1206 one arcsecond corresponds in this case to $\sim 5.39 \mathrm{kpc}$.

## 4. Modelling of galaxies

As the focus lies on early-type galaxies and to keep the computational time to a reasonable amount, the choice has been made to fit the sample galaxies with a one-component fit, i.e. a singular Sérsic profile. Although this may seem to be an oversimplification, single component fits have already proven to be sufficient to describe elliptical galaxies (e.g. Häußler et al., 2013), usually by adopting on average $n \sim 4$ for the Sérsic profile.
Due to the big variety even within the spheroidal galaxy population, the modelling of galaxies can be a very delicate task. Using MegaMorph, a wrapper for SExtractor and GALFIT (see Chapter 2), the user has the opportunities to specify and adjust numerous parameters and input variables in regard to a particular scientific scenario in order to optimize the robustness of the output.
The starting point is to detect all the desired sources, which should enter the fitting process. Due to the big dynamical range in magnitude this can pose a challenging exercise, although MegaMorph provides a two-way method performing SExtractor twice with different setups. The first run intends to detect the bright sources, whereas the second run is used to include the faint ones. Even with this feature this method has to be executed several times in order to identify all the faint objects aimed for, without splitting the brighter ones into numerous sources. An example of a SExtractor input file is given in the appendix. The most influential parameter in this regard is DETECT_MINAREA, which holds the value of coherent pixels of a source needed, to be identified as such.
Unfortunately for MACS 0416 nine sources could not be included in the subsequent fitting process due to their low surface brightness and also partially due to the very prominent ICL. This leads to a reduced sample for this cluster with 83 objects. However those nine rejected sources were nevertheless much too faint to derive any reasonable results and would be therefore discarded in the subsequent analysis anyway.
Before finally commencing the multi-band fitting routine, PSF models have to be provided for each band. The PSFs used for the modelling processes outlined in this
thesis have been contributed by Alberto Molino, who produced them as part of his PhD thesis ${ }^{1}$. In order to create them he has examined all $25 C L A S H$ clusters looking for potential stars and eventually ended up with $\sim 200$, visually selected possible candidates. After making sure that each individual star is visible in all $16 H S T$ filters, not saturated in any of these bands and within an area of 25 pixels free of any contaminating neighbours, he obtained a final sample of $\sim 70$ stars. Ultimately they have been combined and normalized to compose an average model for each band by combining stars from different positions in the image. This procedure is adequate, as the PSF of $H S T$ is almost invariant in regard of time, as well as position, so one PSF is sufficient for each band.

The last part of the preparations concerned the setup of the GALFIT M algorithm itself. A complete input file is included in the appendix, showing the different parts of the setup, i.e. inputs for the SExtractor runs, the creation of postage stamps, the sky determination, the actual modelling using GALFIT $M$ and finally the output catalogue. As it would exceed the scope of this chapter to explain every single input variable, only the most important variables, which are influencing more or less the output and are apparently mainly but not only situated in the GALFIT M part of the setup, are described.

Although GALFIT offers the possibilities to fix or even to fit the sky background by itself, MegaMorph makes no use of this option. It rather calculates the background flux using a different method and keeps it fixed during the subsequent fitting process, which appears to be the more reliable approach (Häussler et al., 2007). By averaging the emitted background flux in increasing elliptical isophotes around the current object with masking interfering neighbouring objects, MegaMorph is able to derive the background flux in dependence to the radius. The final sky estimate is eventually the value, where the slope of this function turns positive for the second time. Of course these assessments can be manipulated by several parameters, such as the number of the sky measurements included to calculate the slope, in the setup file defined in D12, the distance between, as well as the width of the annuli specified in D0 6 and D07, the initial radius D08 or also the minimum separation between two sources set in D19.

The most important parameters however are placed in the GALFIT part, being indicated by an ' E ' in front. E11 constrains the maximum attainable $r_{e}$ and is very generously defined with 400 pixels in this thesis, a value which is never even nearly

[^1]reached. E14 and E15 represent the minimum and maximum limits of $n$, for this work set to 0.2 and 8 , as galaxies with higher $n$ are hardly observed, and E19 holds the value of maximum minutes allowed to model a galaxy.

Nonetheless the big feature distinguishing MegaMorph is its ability to utilize the images in all input bands simultaneously in order to derive a wavelength dependent model for each galaxy. This is achieved by replacing the model parameters by wavelength dependent Chebyshev functions. The degrees of freedom of those polynomials can be specified (E20), enabling the user to set parameters from constant to completely independent in regard to wavelength. For the fits carried out in this thesis the centre position of a galaxy in the form of $x$ and $y$ coordinates, as well as its axial ratio and position angle are considered constant over wavelength. The reason for fixing the latter two parameters is to avoid artificial color gradients, which may be introduced due to different shapes and/or orientation. The structural parameters $r_{e}$ and $n$ are assumed to be best described by a second order polynomial, although a linear function is expected to be sufficient for the brightest elliptical galaxies (Häußler et al., 2013). To obtain the most accurate measurements for magnitudes no constraints are given, hence the magnitudes in the different bands can vary completely independent from the wavelength.

It is obvious that due to this multitude in variables the values have to be chosen carefully and in consideration of their impact on the end result. Hence several test runs, with smaller sample sizes were performed in order to obtain the best models. Unfortunately due to the amount of computational time with a cluster requiring several weeks to be completely done (and eventually also licensing issues) a comparison with two-component fits could not be achieved, but would pose a very interesting project.

Eventually the fitting process has been successful for 103 of the 110 galaxies selected in MACS 1206. For the 7 galaxies left, 4 of them possessing spectroscopic measurements, the remaining 3 only photometric assessments, the maximum time allowed to derive a model has been exceeded, which resulted in an omission of those objects.

Unfortunately for MACS 0416 the results are much worse, with only 42 galaxies being successfully fitted in the initial run. The reason for GALFIT failing to model such a high number of galaxies lies in the prominent ICL in the center of the cluster. Due to the intense background an extensive part, if not the entire image with all the objects residing in it is fitted by MegaMorph, due to their contribution to the
light distribution of the current processed source. This increases the computational time for even just one galaxy to a minimum of several days, which in turn exceeds by far the maximum time allowed to fit a galaxy. In order to raise the number of galaxy models this maximum time parameter had been increased several times in subsequent runs, as well as the number of objects selected for each run had been constraint to a handful. This method has allowed to obtain additional 16 model fits, resulting in a final number of 58 galaxies. Regrettably it has not been possible to finish the analysis over the total sample due to a missing $I D L$ license.

An example output by GALFIT is shown in Fig. 4.1. The left column displays the example galaxy in the respective postage stamps of the different filters, in the middle column the corresponding models are shown, while eventually the right column depicts the various residuals. Clearly visible is the increase in flux with increasing wavelength, indicating that an old stellar population is observed. As the data images further possess the same flux scaling, also the growth in $\mathrm{S} / \mathrm{N}$ towards redder filters is obvious.

The results of the fits are shown in the form of a color-magnitude diagram in Fig. 4.2. The coloured points represent the fitted galaxies, depicting their apparent integrated magnitude in the AB photometric system in $Y_{105}$ on the x axis and their $r_{625}-Y_{105}$ color, corresponding to $B-I$ in rest-frame, on the ordinate. The points themselves are color coded according to the derived Sérsic index $n$ of the galaxies, with blue indicating the lower end, while red denotes the high- $n$ objects.

The underlying gray dots are the corresponding measurements taken from the CLASH SExtractor catalogue, depicting all sources, independent of redshift, which are visible in the WFC3/IR field of view. The good agreement in regard of the location of the clearly recognizable red sequence proves the good quality of the MegaMorph results. Furthermore it is fairly obvious that the majority of the fitted objects lies in the red sequence and possesses $n$ values corresponding to early-type galaxies. This observation meets the expectations, as the focus lies on the innermost part of the cluster, which is supposed to be populated mainly by bright, massive elliptical galaxies.

The subsequent analysis required a separation in early and late type galaxies, which has been achieved by combining two different methods. Following several publications (e.g. Shen et al., 2003; Barden et al., 2005; Vulcani et al., 2014) the galaxies have been separated at first by their $n$ measurements in F814W, which has been the reference band for the fit, with the distinction being drawn at $n=2.5$.


Figure 4.1.: Example of a multi-band GALFIT fitting output for a galaxy of the MACS 1206 sample. From left to right: Original input images, models and residuals for the respective fits. The scaling of the images is matched, which in turn also illustrates the different detection levels in the different filters.


Figure 4.1. (cont.)


Figure 4.1. (cont.)


Figure 4.2.: Color-magnitude diagrams for MACS 1206 (left) and MACS 0416 (right). The gray dots represent all detected sources in the WFC3/IR field of view, overlayed by the fitted galaxies, color coded according to their Sérsic index $n$.

Hence galaxies with $n<2.5$ have been discarded from further considerations. Additionally a visual inspection of the sources in the F814W band has been performed three times for each cluster, with the objects being examined each time in a different order. As a result all galaxies displaying $n \gtrsim 2.5$, but showing spiral or peculiar features are being rejected as well. Applying these constraints leads to two final samples, one for each cluster, consisting of 79 galaxies ( 57 of them spectroscopically analyzed) for MACS 1206 and 38 galaxies ( 37 of them possessing spectroscopic measurements) for MACS 0416. The tables containing the determinations in total magnitude, $r_{e}$ and $n$ for both clusters can be found in the appendix.

With the deduction of the magnitude and the structural parameters $r_{e}$ and $n$ for each galaxy in each band it is now possible to analyze the variation of those values over wavelength. Additionally it is feasible to produce the respective surface brightness profiles based on the model arguments for each selected galaxy in each observed filter. This eventually has been achieved by utilizing the formula (1.9) and all the pre-calculations necessary presented in section 3 of chapter 1 (Fig. 4.3). The $5 \sigma$ magnitude assessments for extended sources by Annunziatella et al. (2014) (Tab. 3.1) serve as limiting magnitudes to calculate the surface brightness profiles for the respective filters. However as these profiles are purely modelled being based on the galaxy properties, there is no real need to truncate the profiles at exactly these limits and any restraints can be chosen, as long as the desired radius intervals
are within those constraints. Hence due to the completely modelled nature of the profile, the precision in radius can be chosen as accurately as desired as well as the differences of the two profiles, e.g. the colors (see formula (1.18)), can be calculated at any predetermined radii (Fig. 4.3). The slopes of the color profiles, i.e. the color gradient in a chosen interval, are then derived by applying a least squares approximation to the calculated color profiles on a logarithmic radius scale.

The combinations chosen to proceed are the optical color $r_{625}-Y_{105}$ and the optical-infrared color $Y_{105}-H_{160}$, which correspond to rest-frame $B-I$ and $I-Y$. Since the $4000 \AA$ break at $z \sim 0.4$ lies approximately at $5500 \AA$, both colors are well situated on the redder wavelength side of the break and are tracing therefore a stellar population of intermediate and older ages, representing the majority of stars in a typical elliptical galaxy. Such a combination of rest-frame optical and optical-IR colors appears to be best suitable to break the age-metallicity degeneracy, with the first color tracing age as well as metallicity, whereas the latter color traces primarily metallicity (Smail et al., 2001; Worthey, 1994).
To be consistent with other studies examining the color gradients of early-type galaxies (e.g. Guo et al., 2011; La Barbera et al., 2002, 2003; Saglia et al., 2000; Tamura \& Ohta, 2000), but also to benefit from the high quality and deep imaging of the HST, which is capable of observing galaxies out to large radii, the gradients are derived in an interval between $0.1 r_{e}$ and $2 r_{e}$ of the respective bluer band. The subsequent discrete color measurements in respect to the SSP comparisons are performed at five different, predetermined radii, $0.1,0.5,1$ and $2 r_{e}$, again in regard to the bluer band in order to probe the potentially, different stellar subpopulations. Furthermore just for the color gradient analysis the colors $g_{475}-I_{814}$ and $I_{814}-H_{160}$, resembling in rest-frame $U-V$ and $V-Y$, are included.
In Fig. 4.3 an example output for one of the galaxies of the MACS 1206 sample is shown. The upper plot depicts in blue the surface brightness profile for $r_{625}$ as well as the position of the respective $0.1,1$ and $2 r_{e}$, represented by the dashed lines. The same, just with red colors applies for $Y_{105}$. As expected the magnitude in the redder band is higher, since the old stellar population dominates the longer wavelength regime. In the lower part of the figure with the x -axis being logarithmically scaled the according color profile in black as well as the respective least squares fit in magenta are displayed. The coloured dashed lines ranging from red to blue illustrate the location of $0.1,0.2,0.5,1$ and $2 r_{e}$ of the bluer band, in this case the $r_{625}$ filter.

The second method to derive the color gradients is a more 'analytical' one following La Barbera et al. (2002) (see formula (1.20)). This approach quantifies the ratios of the effective radii of the processed bands in respect to the interval boundaries $r_{m}$ and $r_{M}$, in which the color profiles and gradients should be calculated in. To be consistent with La Barbera et al. (2002), $r_{m}=0.1 r_{e, 625}$ and $r_{M}=r_{e, 625}$ are adopted.


Figure 4.3.: Surface brightness profiles (upper plot) and color profile with overlayed fit (lower plot) for a MACS 1206 early-type galaxy. The blue and red dotted and dashed lines in the upper diagram correspond to $0.1,1$ and $2 r_{e}$ of the respective bands. The, from red to blue coloured, dashed lines in the lower figure correspond to $0.1,0.2,0.5,1$ and $2 r_{e}$ of the bluer band, in this particular case the $r_{625}$ filter.

## 5. Comparison with SSP models

In order to infer information about the ages and metallicities of the underlying stellar population from the derived colors, it is necessary to compare them to predictions produced by stellar population synthesis models. To achieve that goal in this thesis SSP models have been adopted in several ways making use of the tool EzGal (Mancone \& Gonzalez, 2012).

This Python wrapper allows to produce synthetic magnitudes and colors depending on formation redshift $z_{f}$, metallicity $Z$, starformation history and IMF based on different input models, such as the ones published by Bruzual \& Charlot (2003) (BC03) or Maraston (2005). Additionally it is capable of calculating these values in an observed frame for the desired filters at a predetermined redshift by redshifting the model spectra and convolving it with the filter curves and applying the corresponding $k$-corrections. Thus it is possible to compare directly the measured colors to several modelled ones, differing due to varying input ages and metallicities, with these parameters in turn being linked to the formation and evolution conditions of a galaxy.

The models chosen were the ones by BC 03 , as only those are supplied with different metallicities, which makes it also possible to interpolate additional SED evolutions with other metallicities than the predetermined ones. Furthermore a Chabrier IMF (Chabrier, 2003) has been adopted. As only early-type galaxies are analyzed, simple stellar population (SSP) models are utilized, characterized by a singular important star formation epoch - a so-called star formation burst - at the formation redshift and a subsequent passive evolution. The following cosmological parameters $H_{0}=$ $70 \mathrm{~km} \mathrm{~s}^{-1} \mathrm{Mpc}^{-1}, \Omega_{\Lambda}=0.7$ and $\Omega_{M}=0.3$ (Spergel et al., 2003) are used, being in agreement with other studies examining the CLASH clusters (e.g. Annunziatella et al., 2014; Balestra et al., 2015; Biviano et al., 2013; Girardi et al., 2015; Jouvel et al., 2014; Ogrean et al., 2015; Presotto et al., 2014; Zitrin et al., 2012).

The first approach to derive indications about the development and properties of the stellar population of the selected elliptical galaxies consisted of using models with their metallicities at cluster redshift being fixed to the solar value, so $Z_{\odot}=0.02$.

The formation redshifts employed range from $z=0.5$ to $z=10$ with eleven steps in between, corresponding to ages between $\sim 5 \mathrm{Gyr}$ and $\sim 13 \mathrm{Gyr}$ at $z=0$. The modelled colors are consequently matched with the measured ones and the corresponding age of the best fit allows drawing conclusions about the formation epoch of the respective galaxy, assuming solar metallicity.

The second method applied is very much alike to the first one, only this time metallicity is allowed to vary and the age of the models is fixed to an age of 12 Gyr at $z=0$, which corresponds to a formation redshift $z_{f} \sim 4$. The supplied BC 03 models provide predictions for $Z=[0.008,0.02,0.05]$, which corresponds to $Z \sim[-0.4,0.0,0.4]$ $\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$. Additionally the interpolation capabilities of EzGal are exploited by adding three models with $Z=0.015, Z=0.03$ and $Z=0.04$ equaling to $Z=-0.1$ $\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right], Z=0.2\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$ and $Z=0.3\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$. By matching again the predicted and the observed colors, constraints on the metallicities of the underlying stellar populations can be applied with the limitation of an assumed age.

The issue with both of these approaches is that either the current metallicities or the ages must be fixed for the galaxies. Even with deploying these methods on several colors to minimize the effects of the age-metallicity degeneracy, a significant impact of this phenomenon cannot be ruled out.

To break this degeneracy one can create a grid of two colors at a certain redshift, whose predicted values depend on different SSPs based on varying ages and metallicities. Such a grid is not unlike the one spanned by the right choice of Lick abundance indices and is therefore also capable to trace age in one direction and metallicity in the other, provided the right color combination is chosen. Following Smail et al. (2001) and also La Barbera et al. (2002) a combination of optical and optical-infrared colors appears best suited to resolve this issue. Hence the two colors chosen are $r_{625}-Y_{105}$ for the optical part and $Y_{105}-H_{160}$ for the optical-IR one. Those colors translate at the redshift $z \sim 0.4$, which is applicable for both clusters, to $B-I$ and $I-Y$ in the rest-frame.

Once the grid is finally generated, one can match the individual measured colors with the nearest point of the grid and hence can infer the best fitting age and metallicity for this particular color at this distinct redshift.

## 6. Results

This chapter outlines the main results of this thesis, split into three sections. The first part considers potential wavelength dependencies of the structural parameters $r_{e}$ and $n$. The second part presents the derived color gradients, whereas in the last section the comparison of the color measurements at different radii with the colors predicted by the SSP models and the effect on age and metallicity determinations of the stellar populations of the galaxies are described.

### 6.1. Wavelength-dependence of parameters

The eventually obtained samples of galaxy models for both clusters consists of 117 early-type galaxies in total, fitted in each of the 12 bands, providing - among other things - measurements for $n$ and $r_{e}$ for each filter. Hence it is possible to examine potential trends of these parameters with wavelength, which are already observed by numerous studies (Chan et al., 2016; La Barbera et al., 2002, 2003, 2010; Kelvin et al., 2012; Kennedy et al., 2015; Vulcani et al., 2014).
The trends presented in this thesis consider the average across both cluster samples for each passband with the uncertainties being estimated as $1.253 \sigma / \sqrt{N}$, with $\sigma$ denoting the standard deviation, whereas $N$ is the number of galaxies in the respective sample, following the evaluation by Kennedy et al. (2015) and Vulcani et al. (2014).

The distributions in form of histograms for both clusters and both parameters are shown in the Figures $6.1-6.4$. It is evident that a relatively high number of MACS 1206 galaxies in $B_{435}$ adopt the maximum attainable $n=8$ (Fig. 6.1). This is presumably due to the low $\mathrm{S} / \mathrm{N}$ in this particular band, causing MegaMorph to encounter difficulties determining the right profile. Nevertheless the lack of galaxies obtaining such a high value for $n$ in the other bands, proves that those fits can be treated with a reasonable confidence. The distribution of $r_{e}$ of the MACS 1206 sample (Fig. 6.2) shows no conspicuous features in all bands and, as mentioned earlier, the maximum allowed value for a fit of 400 pixels, which corresponds to $\sim 150 \mathrm{kpc}$
at $\mathrm{z} \sim 0.4$, is not even nearly reached. Instead the median values for $r_{e}$ adopt expectable values of $\sim 2 \mathrm{kpc}$.


Figure 6.1.: Distribution of the Sérsic index $n$ for the MACS 1206 sample in the individual bands. The solid magenta line indicates the median, while the dashed magenta lines denotes the uncertainty interval. The respective filters are given in the upper left corners.


Figure 6.2.: Distribution of the effective radius $r_{e}$ for the MACS 1206 sample in the individual bands. The solid magenta line indicates the median, while the dashed magenta lines denotes the uncertainty interval. The respective filters are given in the upper left corners.

Unfortunately the results for MACS 0416 do not look to be as trustworthy as the assessments for MACS 1206, with approximately 15 percent of the galaxies in the sample attaining an $n=8$ in the $B_{435}$ and $g_{475}$ bands (Fig. 6.3). However the $r_{e}$ measurements (Fig. 6.4) present themselves very similar to the $r_{e}$ distributions of MACS 1206, showing also average effective radii of $r_{e} \sim 2 \mathrm{kpc}$ over all bands.


Figure 6.3.: Same as figure 6.1 only for the MACS 0416 sample.


Figure 6.4.: Same as figure 6.2 only for the MACS 0416 sample.

Still the distributions of $n$ and $r_{e}$ vary significantly with wavelength, proving thereby that the parameters were not too constrained in regard to the degrees of freedom for the Chebyshev polynomials during the fit. This is especially satisfying with respect to the results for MACS 1206, as the same setup as for MACS 0416 has been used ${ }^{1}$, but with the distributions for $n$ and $r_{e}$ differing hardly over the different passbands.

Nevertheless it seems that due to the much more prominent ICL of MACS 0416, the modelling of the galaxies requires a much more comprehensive treatment. Also it can not be ruled out that a significant amount of galaxies in the MACS 0416 sample possess a much more complex stellar population, due to the merging process of the subclusters and the resulting galaxy interactions. For such galaxies a onecomponent fit would not be sufficient and a fit with at least two Sérsic profiles would perhaps clarify this issue. Another approach to gain insights on this matter is to obtain information on the ages of those galaxies at different radii, which coincides with the main goal of this thesis.

To evaluate the measurements, the tables 6.1 and 6.2 list the medians of $n$ and $r_{e}$ for both clusters and for each band including the mentioned error estimates. Already with a first glimpse it is noticeable that for the early-type galaxies in MACS 1206 as well as to a lesser extent in MACS 0416, $n$ appears to be relatively independent of wavelength, while $r_{e}$ definitely seems to decrease with increasing wavelength.

Table 6.1.: Median values for $n$ and $r_{e}$ for the MACS 1206 sample in each filter. The uncertainties are derived by using $1.253 \sigma / \sqrt{N}$.

|  | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{I}_{775}$ | $\mathrm{I}_{814}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $n$ | $4.91 \pm 0.22$ | $4.91 \pm 0.21$ | $4.72 \pm 0.19$ | $4.63 \pm 0.18$ | $4.38 \pm 0.17$ | $4.37 \pm 0.17$ |
| $r_{e}$ | $2.53 \pm 0.29$ | $2.41 \pm 0.29$ | $2.23 \pm 0.28$ | $2.17 \pm 0.27$ | $2 \pm 0.26$ | $1.96 \pm 0.26$ |
|  |  |  |  |  |  |  |
|  | $\mathrm{z}_{850}$ | $\mathrm{Y}_{105}$ | $\mathrm{YJ}_{110}$ | $\mathrm{~J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| $n$ | $4.26 \pm 0.17$ | $4.2 \pm 0.17$ | $4.14 \pm 0.17$ | $4.17 \pm 0.17$ | $4.23 \pm 0.18$ | $4.35 \pm 0.18$ |
| $r_{e}$ | $1.93 \pm 0.26$ | $1.89 \pm 0.25$ | $1.84 \pm 0.25$ | $1.83 \pm 0.25$ | $1.82 \pm 0.25$ | $1.76 \pm 0.25$ |

[^2]Table 6.2.: The same as table 6.1, only for the MACS 0416 sample.

|  | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{I}_{775}$ | $\mathrm{I}_{814}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $n$ | $5.54 \pm 0.35$ | $5.27 \pm 0.32$ | $4.85 \pm 0.26$ | $4.9 \pm 0.25$ | $4.2 \pm 0.23$ | $4.07 \pm 0.23$ |
| $r_{e}$ | $2.88 \pm 0.34$ | $2.73 \pm 0.33$ | $2.43 \pm 0.32$ | $2.37 \pm 0.32$ | $2.18 \pm 0.34$ | $2.12 \pm 0.34$ |


|  | $\mathrm{z}_{850}$ | $\mathrm{Y}_{105}$ | $\mathrm{YJ}_{110}$ | $\mathrm{~J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $n$ | $3.77 \pm 0.24$ | $3.58 \pm 0.25$ | $3.54 \pm 0.26$ | $3.51 \pm 0.26$ | $3.54 \pm 0.27$ | $3.65 \pm 0.29$ |
| $r_{e}$ | $2.06 \pm 0.36$ | $2.02 \pm 0.37$ | $2 \pm 0.38$ | $1.98 \pm 0.38$ | $1.98 \pm 0.38$ | $2.01 \pm 0.37$ |

To quantify the dependence on wavelength, $\mathcal{N}$ and $\mathcal{R}$ are introduced. These variables describe the medians of the ratios between the $n$ and $r_{e}$ measurements in the $H_{160}$ and the $r_{625}$ band and can therefore provide information on how those parameters vary with wavelength. These two filters have been selected on one hand because they are also used throughout the color gradient and subsequent SSP analysis and on the other hand due to their provision of the widest range in wavelength. This interval could only be increased by relying on the $B_{435}$ or $g_{475}$ band, but both of these filters possess a significant lower $\mathrm{S} / \mathrm{N}$ and, as already mentioned above, especially some of the assessments of $n$ in $B_{435}$ should be treated with caution.

Now as $n$ describes the shape of the light profile, an $\mathcal{N}<1$ in this context corresponds to a higher central light concentration in the bluer $r_{625}$ band than in the redder $H_{160}$ filter. Hence $\mathcal{N}>1$ obviously results in an opposite behavior, while an $\mathcal{N} \sim 1$ implies a similar shape of the light profile in those two passbands.

The parameter $\mathcal{R}$ describes the variation in size in these two filters, with $\mathcal{R}<1$ indicating a smaller size in the redder $H_{160}$ band. As a result the centers appear redder, while the outskirts present themselves bluer in comparison. For $\mathcal{R}>1$ the opposite is true, with the centers being bluer than the outskirts, whereas $\mathcal{R} \sim 1$ implies a similar size in both passbands. Hence, much like color gradients, $\mathcal{R}$ serves as an indicator of the color variation with radius within an object.

Calculating the medians for $\mathcal{N}$ and $\mathcal{R}$ for both cluster samples yield the results presented in table 6.3. The uncertainties are again calculated using the $1.253 \sigma / \sqrt{N}$ estimation.

Although the $n$ measurements for the MACS 0416 sample are not as robust as the assessments for galaxies in MACS 1206, it is feasible to consider $\mathcal{N} \sim 1$ for both populations, which means that the observed objects can be described indeed
by a one-component profile. These conclusions hold true especially for the MACS 1206 sample with $\left\langle n_{H, 160}\right\rangle \sim 0.9\left\langle n_{r, 625}\right\rangle$. For galaxies in MACS 0416 the value is with $\left\langle n_{H, 160}\right\rangle \sim 0.8\left\langle n_{r, 625}\right\rangle$ a little bit lower. This result can be explained by considering the decreased accuracy for $n$ measurements and a possible contamination by non-truly elliptical galaxies possessing two or more significant stellar components. Nevertheless the majority of the MACS 0416 galaxies can be also characterized by a singular light profile.

Table 6.3.: Medians of the $n_{H, 160} / n_{r, 625}(\mathcal{N})$ and $r_{e_{H, 160}} / r_{e_{r, 625}}(\mathcal{R})$ measurements for both cluster samples.

|  | $\mathcal{N}$ | $\mathcal{R}$ |
| :---: | :---: | :---: |
| MACS 1206 | $0.92 \pm 0.02$ | $0.75 \pm 0.02$ |
| MACS 0416 | $0.84 \pm 0.04$ | $0.77 \pm 0.04$ |

Examining $\mathcal{R}$, the results for MACS 1206 and MACS 0416 hold the same outcome with the effective radii being $\sim 25$ percent smaller in $H_{160}$ than in $r_{625}$. This already indicates that the majority of the scrutinized galaxies possess redder centers and bluer outskirts, hence negative color gradients.

To visualize the ratios for every individual galaxy, Fig. 6.5 displays


Figure 6.5.: $\mathcal{N}$ against $\mathcal{R}$ for every galaxy. The peak of $\mathcal{N}$ tends to $\sim 0.9$, while $\langle\mathcal{R}\rangle \sim$ 0.75 . the $\mathcal{N}=n_{H, 160} / n_{r, 625}$ versus the $\mathcal{R}=r_{e_{H, 160}} / r_{e_{r, 625}}$ for the entire elliptical galaxies sample, without making a distinction between MACS 1206 and MACS 0416 objects. The uncertainties returned by GALFIT for the $n$ and $r_{e}$ assessments are smaller than the size of the markers. The gray dashed lines are indicating unity in $\mathcal{N}$ and $\mathcal{R}$. The trend for $\mathcal{N}$ and especially $\mathcal{R}$ is fairly obvious with them being smaller in the redder band. This suggests the tracing of different superimposed stellar populations, with the bluer one extending to larger radii, but showing a similar light distribution.

Also a possible correlation between $n$ and $r_{e}$ has been examined. However independent of the inspected passbands or cluster samples no trends are observed.

### 6.2. Color gradients

To derive the color gradients two approaches have been used. The first one has been to fit linear least squares (lsq) fits to the derived color profiles with the radii being logarithmically scaled (Fig. 4.3). The slopes of these fits equal the change in color in a predetermined radius interval, e.g. the color gradient.

The second method applied has been the utilization of the formula by La Barbera et al. (2002) (see formula (1.20)) with $r_{m}=0.1 r_{e, 625}$ and $r_{M}=r_{e, 625}$.
The interval for which the color gradient has been derived ranges from $0.1-2 r_{e}$ in regard to the bluer band of the color for each galaxy. The colors examined are $g_{475}-I_{814}, r_{625}-Y_{105}, I_{814}-H_{160}$ and $Y_{105}-H_{160}$, which are corresponding at $z \sim 0.4$ to $U-V, B-R, V-Y$ and $R-Y$ in rest-frame. As the $g_{475}$ band lies on the bluer end of the 4000 Ångstrom break at the cluster redshifts, it is also possible to detect a potential younger stellar population and their subsequent probable impact on the color variation with increasing radius. The order of magnitude of the errors returned by GALFIT lies in the range of $10^{-3}-10^{-4}$ for all parameters and is therefore not considered in the subsequent analysis, as an inclusion of those uncertainties would result in a non discernable effect. To derive more accurate uncertainty estimates it would have been necessary to rely on simulations of the respective clusters containing similar galaxy populations as the observed ones. An ensuing MegaMorph analysis of those clusters would then deliver more information on the introduced errors, as the results of both modelling processes, one for the observed and one for the simulated galaxies, whose properties are known beforehand, can be compared. Unfortunately such a procedure has not been able to be completed in the given timeframe, but is currently conducted by Ulrike Kuchner as a part of her PhD thesis.

The individual results for each galaxy are presented in the tables, which can be found in the appendix, while the medians of the color gradients for each color and cluster are presented in the Tables 6.4 and 6.5 . The first columns display the respective colors, whereas the second and third columns contain the medians of the fitted respectively calculated color gradients. The uncertainties in regard to the medians are again calculated using $1.253 \sigma / \sqrt{N}$. The fourth columns indicate the results of Kolmogorov-Smirnov (KS) tests between the respective fitted and calcu-
lated color gradients assessments. Quantifying the distance between two distributions, the smaller the respective value derived by the KS test, the more equal are the two examined distributions.

Table 6.4.: Medians of the color gradients for the MACS 1206 sample in four different colors.

| Color | $\nabla_{\text {median,fit }}$ <br> $\left[\mathrm{mag} \mathrm{dex}^{-1}\right]$ | $\nabla_{\text {median,calc }}$ <br> $\left[\mathrm{mag} \mathrm{dex}^{-1}\right]$ | $K S$ |
| :---: | :---: | :---: | :---: |
| $g_{475}-I_{814}$ | $-0.17 \pm 0.01$ | $-0.13 \pm 0.01$ | 0.33 |
| $r_{625}-Y_{105}$ | $-0.18 \pm 0.01$ | $-0.15 \pm 0.01$ | 0.24 |
| $I_{814}-H_{160}$ | $-0.13 \pm 0.02$ | $-0.13 \pm 0.02$ | 0.08 |
| $Y_{105}-H_{160}$ | $-0.06 \pm 0.01$ | $-0.07 \pm 0.01$ | 0.08 |

Table 6.5.: Medians of the color gradients for the MACS 0416 sample in four different colors.

| Color | $\nabla_{\text {median,fit }}$ <br> $\left[\mathrm{mag} \mathrm{dex}^{-1}\right]$ | $\nabla_{\text {median,calc }}$ <br> $\left[\mathrm{mag} \mathrm{dex}^{-1}\right]$ | $K S$ |
| :---: | :---: | :---: | :---: |
| $g_{475}-I_{814}$ | $-0.11 \pm 0.02$ | $0.02 \pm 0.02$ | 0.45 |
| $r_{625}-Y_{105}$ | $-0.13 \pm 0.02$ | $-0.05 \pm 0.02$ | 0.32 |
| $I_{814}-H_{160}$ | $-0.12 \pm 0.02$ | $-0.11 \pm 0.02$ | 0.11 |
| $Y_{105}-H_{160}$ | $-0.07 \pm 0.02$ | $-0.08 \pm 0.02$ | 0.08 |

All averaged fitted color gradients show a negative result, i.e. a decrease in magnitude with radius, indicating that the centers of the majority of the galaxies are redder than their exterior regions. For MACS 1206 this holds also true for the calculated color gradient measurements, being in good agreement with the assessments by the lsq fitting method. However for MACS 0416 these observations are only completely valid for the two 'redder' colors $I_{814}-H_{160}$ and $Y_{105}-H_{160}$. While the calculated median color gradient for $r_{625}-Y_{105}$ is still negative, the discrepancy with a factor $\sim 2$ between the fitted and the calculated result is quite significant. Even more prominent is this effect for the $g_{475}-I_{814}$ color, with the calculated median color gradient being neutral or slightly positive. These implications are also reflected in the outcome of the KS tests, with the 'redder' colors for both clusters obtaining values implying both distributions to originate from the same sample.

The reason for the deviating results in regard to the calculated color gradients for the MACS 0416 sample lies presumably in the approach itself. While the fitting
method directly compares the two respective bands and returns expected results, the recipe by La Barbera et al. (2002) introduces additional parameters $r_{m}$ and $r_{M}$, with these being chosen to be 0.1 and $1 r_{e}$ from the $r_{625}$ band to be consistent with previous publications (La Barbera et al., 2002, and references therein). A subsequent utilization of this procedure on a color containing the $g_{475}$ band, amplifies potential misinterpretations, which can not be ruled out in these both bands due to the low $\mathrm{S} / \mathrm{N}$ because of the prominent ICL, and leads eventually to inaccurate results.

To visualize the assessments for the galaxies more individually Fig. 6.6 and 6.8 show the distributions of the gradients in form of histograms for the four respective colors, distinguished into the color gradients derived by the fitting method in red and by the calculative approach in blue. The corresponding colored, dashed lines denote the medians, the dotted, gray line the position of a neutral color gradient, e.g. $\nabla\left(\lambda_{1}-\lambda_{2}\right)=0$.


Figure 6.6.: Distributions of the color gradients for the four colors depending on the deriving method. Red indicates the color gradients assessed by lsq fitting, blue displays the calculated gradients. The corresponding dashed lines denotes the respective medians, whereas the gray dotted line indicates a color gradient $=0$.

For the MACS 1206 sample an overwhelming majority shows negative or in the case of $Y_{105}-H_{160}$ at least neutral color gradients ranging from $\sim 0.1$ to $\sim-0.3$ difference in magnitude per dex in radius. The distinction between the fitted and calculated color gradients decrease with wavelength of the bands constituting the respective colors. In this sense the differences are most striking for the $g_{475}-I_{814}$ color, which is not surprising considering the low $\mathrm{S} / \mathrm{N}$ in $g_{475}$ and the resulting potential miscalculation of the structural parameters for some of the galaxies.


Figure 6.7.: Differences in the color gradient results for the MACS 1206 galaxies in the four colors between the lsq approach and the calculations following La Barbera et al. (2002). The magenta dot represents the median value for both measurements, the gray dashed line unity and the gray dotted lines a color gradient $=0$.

Nevertheless the obvious general trend for the 79 galaxies is to have negative color gradients, with no more than two galaxies exhibiting at any point color gradients $>0.1$ mag dex ${ }^{-1}$, independent of color or method:

- For $g_{475}-I_{814}$ and $r_{625}-Y_{105}$ none of the galaxies possess a fitted color gradient $>0.1 \mathrm{mag}_{\mathrm{dex}}{ }^{-1}$ and only one galaxy does that in regard to the calculated gradients, which corresponds to $\sim 1.3$ percent out of the whole sample.
- In $I_{814}-H_{160}$ again only one galaxy $(\sim 1.3 \%)$ shows an analogous behavior having a color gradient $>0.1$ mag dex ${ }^{-1}$ with both approaches.
- Finally for $Y_{105}-H_{160}$, two galaxies $(\sim 2.6 \%)$ display such a color gradient in the fitted category as well as the calculated section, with one being the implied galaxy in $I_{814}-H_{160}$.

Fig. 6.7 gives a summary of the color gradient calculations for the elliptical galaxies in MACS 1206. The abscissa denotes the fitted color gradients, while the ordinate represents the calculated ones. The gray dots represent the individual galaxy measurements, the magenta dot the median of the color gradients of the respective assessment method, with its size indicating the uncertainty. The gray, dashed line indicates unity between the color gradients, while the dotted, gray lines distinguish the results into positive and negative measurements. Again the negligible amounts of galaxies possessing positive color gradients are evident, as well as the increasing scatter between the derivation methods with 'bluer' colors.


Figure 6.8.: Same as Fig. 6.6, only for the MACS 0416 sample.

Unfortunately as expected, the assessments are less coherent for the MACS 0416 sample, as can be seen in Fig. 6.8. Especially in $g_{475}-I_{814}$ the color gradients show a significant separation in the methods how they are derived, with the median of the calculated gradients being even slightly positive.

For $r_{625}-Y_{105}$ the separation is still distinct, however the majority of the galaxies already display a negative color gradient. For the two 'redder' bands $I_{814}-H_{160}$ and $Y_{105}-H_{160}$ the results are very similar to those for MACS 1206 with the medians being as good as identical and only a negligible amount of objects shows a positive color gradient. This overall behaviour reflects again the impact of $r_{m}$ and $r_{M}$ on the global result, especially with the measurements in the bluer bands $g_{475}$ and $r_{625}$ being significantly influenced by the ICL.


Figure 6.9.: Same as Fig. 6.7, only for the MACS 0416 sample.

Still also for these less precise measurements in comparison to MACS 1206 the number of galaxies showing a color gradient $>0.1 \mathrm{mag} \mathrm{dex}^{-1}$ is minor:

- The number of galaxies in $g_{475}-I_{814}$ with a fitted color gradient $>0.1 \mathrm{mag}$ $\mathrm{dex}^{-1}$ amounts only to 2 , which corresponds to $\sim 2.5$ percent of the sample,
while for the calculated sample 6 galaxies ( $\sim 7.6 \%$ ) show such a behavior.
- In $r_{625}-Y_{105}$ only one galaxy ( $\sim 1.3 \%$ ) surpasses this condition with its fitted color gradient, while none of the galaxies achieve that with the respective calculated color gradients.
- Finally for $I_{814}-H_{160}$ and $Y_{105}-H_{160}$ no galaxy shows a color gradient $>0.1$ mag dex ${ }^{-1}$, independent of the method.

Analogically to Fig. 6.7, Fig. 6.9 also summarizes the color gradient results for MACS 0416. The larger scatter in $g_{475}-I_{814}$ and $r_{625}-Y_{105}$ is clearly visible, but so is also the relative low number of positive color gradients, except for the calculated color gradients in $g_{475}-I_{814}$.

To determine any potential correlations of the color gradients Fig. 6.10 and Fig. 6.11 depict the gradients in all four colors versus different galactic properties.


Figure 6.10.: Fitted and calculated color gradients in the four colors for the MACS 1206 galaxies versus projected cluster-centric distance $R[\mathrm{Mpc}]$ (left), total magnitude in $I_{814}$ $\operatorname{Mag}_{I, 814}[\mathrm{AB} \operatorname{mag}]$ (middle) and stellar mass $\mathrm{M}_{*}\left[\mathrm{M}_{\odot}\right]$ (right).


Figure 6.11.: Same as Fig. 6.10, only for MACS 0416.

The fitted gradients are again indicated by the red dots, the calculated ones by the blue markers. Neither for the galaxies in MACS 1206 nor for the objects in MACS 0416 any dependency on the projected cluster centric distance (left columns) can be found. The distances are calculated by moving all galaxies into a plane situated at the respective mean cluster redshift.

Also plotting the gradients of both clusters versus the respective total integrated magnitudes in $I_{814}$, which has been the reference band for the fitting process, reveals no visible scaling relations.

Finally the stellar masses have been derived for 56 out of the 79 galaxies in MACS 1206 and 13 out of the 38 objects in MACS 0416 using the photometric measurements by Subaru. The mass range of the galaxies covers mainly an interval of $10^{10} \lesssim M_{*} \lesssim$ $10^{11}$, which are typical values for bright elliptical galaxies situated in the centers of clusters. Likewise no dependencies for the color gradients on the stellar mass of the respective galaxies can be found.

As the gradients are calculated out to $2 r_{e}$ it can be excluded that a lack of cor-
relations is caused by a truncation of the measurements of the color variations and a subsequent comparison with global galactic properties. Hence the position of an early-type galaxy in a cluster, as well as its luminosity and stellar mass does not influence the color gradients and therefore also the mechanisms responsible for the creation of different stellar populations.

### 6.3. Determination of age and metallicity

To infer information about the ages and metallicities of the individual galaxies and their stellar populations, SSP models have been employed in three different ways. Models from BC 03 have been utilized, as EzGal already provides them with three different metallicities ( $Z \sim[-0.4,0.0,0.4]\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$ ), which hence offers the possibility to interpolate additional models with metallicities in between those supplied values.

For the first two approaches each time one of the two parameters age and metallicity got fixed to a certain value, while the other has been allowed to vary in order to model the colors $r_{625}-Y_{105}$ and $Y_{105}-H_{160}$ in the observed frame at the respective cluster redshifts. After applying the extinction correction on those calculated colors, they were compared to the, at different radii measured colors to derive potential constraints on the ages and metallicities of the stellar populations at these particular locations within a galaxy. The radii applied are meant to encompass a large portion of a galaxy and are therefore chosen to be $0.1 r_{e}, 0.5 r_{e}, 1 r_{e}$ and $2 r_{e}$.

The first method used a fixed overall metallicity to the solar value $Z=0$ $\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$, while the age has been allowed to vary, which has been achieved by utilizing different formation redshifts, ranging from $z=0.5$ to $z=10^{2}$, resulting in ages between 5 and 13 Gyr at $z=0$. This leads to a number of discrete values for each respective color at the cluster redshifts, making it possible to compare them with the measured colors. The best matches between models and observations would then give clues about the ages of the stellar populations at the different positions within a galaxy.

The second approach is done analogously to the first one, just with the roles of age and metallicity swapped. Including the interpolated ones, the implemented metallicities are $Z \sim[-0.4,-0.1,0.0,0.2,0.3,0.4]\left[\log \left(Z / Z_{\odot}\right)\right]$, while a formation redshift of $z_{f}=4$ has been assumed, corresponding to an overall age for the galaxy

[^3]of 12 Gyr at $z=0$.
The individual assessments for each galaxy and color are attached in the appendix, while the tables 6.6 and 6.7 contain the results for the median ages and metallicities for both colors and clusters at the predetermined radii. As usual the errors are calculated using $1.253 \sigma / \sqrt{N}$.

These results should be treated with some caution due to several reasons. The most obvious one is of course, that either an age or a metallicity has to be assumed, therefore already introducing a bias. Also the extinction correction, which ranges from $\sim 0.2 \mathrm{mag}$ for $r_{625}$ to $\sim 0.03 \mathrm{mag}$ for $H_{160}$, as well as a possible dust presence can have a significant impact on the colors. Finally as shown in the previous sections the results in $r_{625}$ for the MACS 0416 sample appear to be error-prone.

Table 6.6.: Median results of the SSP analysis for ages and metallicities of the MACS 1206 galaxies at different radii and for both colors.

| Color | Age [Gyr] |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| $r_{625}-Y_{105}$ | $13 \pm 0.3$ | $12 \pm 0.3$ | $9.8 \pm 0.3$ | $7.5 \pm 0.3$ |
| $Y_{105}-H_{160}$ | $13 \pm 0.3$ | $10.2 \pm 0.3$ | $10.2 \pm 0.3$ | $8.7 \pm 0.3$ |
|  | $\mathrm{Z}\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$ |  |  |  |
| $r_{625}-Y_{105}$ | $0.2 \pm 0.03$ | $0.0 \pm 0.03$ | $-0.1 \pm 0.03$ | $-0.4 \pm 0.02$ |
| $Y_{105}-H_{160}$ | $0.2 \pm 0.02$ | $0.2 \pm 0.02$ | $0.0 \pm 0.02$ | $0.0 \pm 0.02$ |

Table 6.7.: Median results of the SSP analysis for ages and metallicities of the MACS 0416 galaxies at different radii and for both colors.

| Color | $0.1 r_{e}$ | $0.5 r_{e}$ |  | $1 r_{e}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Age [Gyr] | $2 r_{e}$ |  |  |  |  |  |
| $r_{625}-Y_{105}$ | $8.7 \pm 0.4$ | $12 \pm 0.4$ | $10 \pm 0.4$ | $7.3 \pm 0.3$ |  |  |
| $Y_{105}-H_{160}$ | $12.5 \pm 0.5$ | $8.7 \pm 0.5$ | $7.5 \pm 0.4$ | $7.5 \pm 0.4$ |  |  |
|  |  |  |  |  |  |  |
|  | $\mathrm{Z}\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$ |  |  |  |  |  |
| $r_{625}-Y_{105}$ | $-0.4 \pm 0.04$ | $0.0 \pm 0.04$ | $-0.1 \pm 0.04$ | $-0.4 \pm 0.03$ |  |  |
| $Y_{105}-H_{160}$ | $0.2 \pm 0.03$ | $0.0 \pm 0.03$ | $0.0 \pm 0.04$ | $0.0 \pm 0.03$ |  |  |

One technique to avoid such issues for example would be to enhance the already applied constraints. For instance in the course of the first method one can assign a solar metallicity to the measured colors at $0.1 r_{e}$ for each individual galaxy as opposed to the entire objects. A corresponding approach for the second procedure would be to assume an average age of 12 Gyr at $z=0$ for the, at $0.1 r_{e}$ measured colors instead for the galaxy itself. Both of these premises would perhaps lead to improved results, but only at the cost of completely assuming the properties of (one of) the underlying stellar populations. So would possess an exterior population of a galaxy with a significant positive color gradient due to an age variance ages much older than 12 Gyrs, which would not agree with any of today's formation scenarios.
However despite all these setbacks, especially in regard to total values, it is still feasible to extract relative results for age and metallicity between the distinct regions. Consequently all assessments show a decrease of age and metallicity with radius with the exception of $r_{625}-Y_{105}$ in MACS 0416. For this particular color the median age difference in MACS 1206 between 0.1 and $2 r_{e}$ amounts to $\sim 4.3 \pm 0.3 \mathrm{Gyr}$, while the metallicity difference is $\sim 0.4 \pm 0.03\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$. For the same range and sample the median age and metallicity variations in $Y_{105}-H_{160}$ add up to $\sim 1.9 \pm 0.4 \mathrm{Gyr}$ and $0.1 \pm 0.03\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$. Regarding the MACS 0416 galaxies in the same range and the same color one finds an age discrepancy of $\sim 2.8 \pm 0.6 \mathrm{Gyr}$ and a metallicity difference of $0.2 \pm 0.04\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$.

To avoid the need of any assumptions at all and to break the age-metallicity degeneracy a third approach has been employed in the form of a grid spanned by two colors. While the optical-infrared color $r_{625}-Y_{105}$ is mainly tracing the age of the stars, the infrared color $Y_{105}-H_{160}$ is primarily influenced by the metallicity of the stellar populations. Other combinations of colors are not that well suited, as both or at least one can be affected equally by age and metallicity.

Fig. 6.12 depicts such a grid in addition to the colors of the individual galaxies of the MACS 1206 sample shown in green. The colors are calculated utilizing the outputs by MegaMorph for the total integrated magnitudes in the respective bands. The magenta dot denotes the median of the age and metallicity assessments with the error being indicated by its size.


Figure 6.12.: $r_{625}-Y_{105}$ vs. $Y_{105}-H_{160}$ diagram with the overall colors of the MACS 1206 galaxies overlayed in green. The magenta dot represents the median of the ages and metallicites of the sample, with the errors being of the order of the size of the marker.

The affiliated histograms (Fig. 6.13) show the distributions of the ages and metallicities of the galaxies, with the median and its errors being again illustrated as the solid respectively dashed magenta lines. The median age lies at $8.7 \pm 0.3 \mathrm{Gyr}$, while the median metallicity appears to attains a solar value of $Z=0.0 \pm 0.03$ $\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$.

However it must be noted that in addition to the significant impact of extinction and dust, the different age tracks between 7 and 9 Gyr are almost identical at solar metallicities, hence introducing an additional error source.

Nevertheless serving as a consistency check it is satisfying to see that, even though there is a high number of potential error contributors, the majority of galaxies lie within the grid and display feasible values.


Figure 6.13.: Distributions for the derived ages and metallicities for the MACS 1206 sample. The medians are indicated by the solid magenta line, the according uncertainties by the dashed lines.

The same analysis has been done for the MACS 0416 galaxies with Fig. 6.14 showing the color-color diagram and Fig. 6.15 the distributions of the ages and metallicities of the individual objects. In addition to the above mentioned uncertainty sources, the error-prone determinations in $r_{625}$ may also be an explanation for the handful of outliers in Fig. 6.14. Nevertheless as can also be seen in Fig. 6.15 the properties seem to be very similar to the MACS 1206 sample with the median age being $7.1 \pm 0.4$ Gyr and the median metallicity $Z=0.1 \pm 0.04$.


Figure 6.14.: $r_{625}-Y_{105}$ vs. $Y_{105}-H_{160}$ diagram with the overall colors of the MACS 0416 galaxies overlayed in green. The magenta dot represents the median of the ages and metallicites of the sample, with the errors being of the order of the size of the marker.


Figure 6.15.: Distributions for the derived ages and metallicities for the MACS 0416 sample. The medians are indicated by the solid magenta line, the according uncertainties by the dashed lines.

One of the main goals of this thesis has been to apply this method not only to the overall colors, but also to the colors at certain radii, namely $0.1,0.5,1$ and $2 r_{e}$ in order to at least reveal the age and metallicity variances within a galaxy.


Figure 6.16.: Color-color diagram for one of the galaxies in the MACS 1206 sample. The red dot represents the color values at $0.1 r_{e}$, the yellow one at $0.5 r_{e}$, the green one at $1 r_{e}$ and finally the blue one at $2 r_{e}$.

To illustrate this approach Fig. 6.16 shows the color values for one example galaxy belonging to the MACS 1206 sample. The measurements at the four different radii are denoted by the coloured markers with the red dot illustrating the assessment at $0.1 r_{e}$, the yellow one at $0.5 r_{e}$, the green one at $1 r_{e}$ and eventually the blue one at $2 r_{e}$. Such an approach reveals that for this particular galaxy the age as well as the metallicity decrease with radius and that therefore both parameters are responsible for the observed color gradient.

Subsequently this analysis has been applied to the entire MACS 1206 sample. Fig. 6.17 illustrates the results with the same color coding as in Fig. 6.16 and
the additional median values depicted by the magenta dots. A significant part of the measurements at $0.1 r_{e}$ show higher ages than the maximum age of 13 Gyr , which holds also true to a lesser extent for the determinations at 0.5 and $1 r_{e}$. The extinction correction can play a critical role as the difference in $r_{625}-Y_{105}$ accounts for $\sim 0.1 \mathrm{mag}$, yet it is not enough to explain this behavior in its entirety.


Figure 6.17.: Color-color diagrams for the color measurements in the MACS 1206 at the four different discrete radii. The color coding is the same as in Fig. 6.16, the magenta dots represent the respective median values with its size indicating the error. The grid points are determined by the same ages and metallicities as in Fig. 6.14.

Dust, errors in the measurements and also uncertainties in the models themselves cannot be ruled out. The need to constrain the possible reasons and their contributions to the errors leads back to the essential requirements of applicable uncertainties concerning the MegaMorph measurements, which are unfortunately not available. As expected the measurements at $2 r_{e}$ are very similar to the ones utilizing the total integrated magnitudes (Fig. 6.12), as the flux contribution at radii $>2 r_{e}$ gets insignificant.

Nevertheless even with the absolute values being problematic, one can still infer some constraints by considering the relative age and metallicity variations with radius. The histograms in Fig. 6.18 and Fig. 6.19 show the distributions in age respectively metallicity for the measurements at the four distinct radii.


Figure 6.18.: Distributions of the assessed ages for the measurements in the MACS 1206 sample at the four different radii, which are indicated in the upper left corner and accordingly color coded. The median and its errors are again depicted as solid respectively dashed magenta lines.


Figure 6.19.: Distributions of the assessed metallicities for the measurements in the MACS 1206 sample at the four different radii, which are indicated in the upper left corner and accordingly color coded. The median and its errors are again depicted as solid respectively dashed magenta lines.

While for $0.1 r_{e}$ approximately half of the galaxies are in the 13 Gyr bin, as this age constitutes as the maximum attainable age, only $\sim 20$ percent of the sample also display such an age at $2 r_{e}$. To quantify this assessment: By averaging the age differences of the individual galaxies the centers appear to be $1.2 \pm 0.4 \mathrm{Gyr}$ older than the exterior regions.

The variety in metallicities does not display such a behavior, nevertheless the metallicity decreases from supersolar values at $0.1 r_{e}$ to a solar regime already at $0.5 r_{e}$ and remains constant. The median metallicity difference therefore measures up to $0.1 \pm 0.03\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$.

As expected the results for MACS 0416 are less encouraging (Fig. 6.20). Especially the age distribution at $0.1 r_{e}$ (Fig. 6.21) with a median of $7.3 \pm 0.4 \mathrm{Gyr}$ seems to be significantly out of line with the other age determinations.


Figure 6.20.: Color-color diagrams for the color measurements in the MACS 0416 at the four different discrete radii. The color coding is the same as in Fig. 6.16, the magenta dots represent the respective median values with its size indicating the error. The grid points are determined by the same ages and metallicities as in Fig. 6.14

This demeanor is presumably again explained by the errors in the $r_{625}$ band. As $n$ determines the shape of the surface brightness profile, its influence is particularly considerable at small radii. Hence imprecise measurements of $n$ can lead to substan-
tially wrong colors, which appears to apply in this case. To verify this hypothesis the individual ages have been derived for a radius $r=0.2 r_{e}$. As predicted the impact of an inaccurate $n$ does not vanish but decreases, which results into a median age of $10.2 \pm 0.5$ Gyr. However neglecting the measurements at the most inner parts and considering a range of $0.5-2 r_{e}$ the sources show the expected trend with the galaxies being on average $3.0 \pm 0.6$ Gyr older in the centers than in the outskirts. Furthermore like it is the case for MACS 1206 also the MACS 0416 sample at $2 r_{e}$ shows a very similar distribution to the one, for which the colors are derived from the total integrated magnitudes (Fig. 6.14).


Figure 6.21.: Distributions of the assessed ages for the measurements in the MACS 0416 sample at the four different radii, which are indicated in the upper left corner and accordingly color coded. The median and its errors are again depicted as solid respectively dashed magenta lines.


Figure 6.22.: Distributions of the assessed metallicities for the measurements in the MACS 0416 sample at the four different radii, which are indicated in the upper left corner and accordingly color coded. The median and its errors are again depicted as solid respectively dashed magenta lines.

As the metallicity is mainly described by the infrared color $Y_{105}-H_{160}$, with both bands possessing much more accurate measurements, the distributions of the galaxies' metallicities at the different radii are in good agreement with the expectations and also the MACS 1206 sample. At $0.1 r_{e}$ more than 50 percent of the galaxies' stellar populations display supersolar metallicities, with additional $\sim 25$ percent revealing metallicities in the solar range.

Increasing the radius confirms the trend of decreasing metallicities, with the median of the individual metallicity variations within a galaxy being $0.1 \pm 0.06$ $\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$.

Table 6.8 summarizes the median assessments of the individual changes in age and
metallicity in both samples. The simultaneous decrease of age and metallicity with radius in both clusters is evident, with the already presented exception of the ages in MACS 0416 at $0.1 r_{e}$. This suggests that both galactic properties are responsible for the color gradients.

Table 6.8.: Medians of the individual measurements for ages and metallicities for both clusters at the four predetermined radii.

| Cluster | $0.1 r_{e}$ | $0.5 r_{e}$ |  | $1 r_{e}$ |
| :---: | :---: | :---: | :---: | :---: |
| Age [Gyr] | $2 r_{e}$ |  |  |  |
| MACS 1206 | $13 \pm 0.3$ | $12 \pm 0.3$ | $10.2 \pm 0.3$ | $10.2 \pm 0.3$ |
| MACS 0416 | $7.3 \pm 0.4$ | $12.5 \pm 0.5$ | $12 \pm 0.5$ | $7.5 \pm 0.5$ |
| $\mathrm{Z}\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$ |  |  |  |  |
| MACS 1206 | $0.2 \pm 0.03$ | $0.0 \pm 0.02$ | $0.0 \pm 0.02$ | $0.0 \pm 0.02$ |
| MACS 0416 | $0.2 \pm 0.04$ | $0.0 \pm 0.04$ | $-0.1 \pm 0.04$ | $0.0 \pm 0.05$ |

In comparison to the first two approaches with assuming a fixed overall metallicity respectively age the SSP grid results show a significant lower variation in both parameters, especially in the respective colors tracing the corresponding properties. This supports the assumption that differences in metallicity as well as age are origins for the color gradients.

## 7. Conclusions

The analyses of the early-type galaxies of the clusters MACS 1206 and MACS 0416 done in the course of this thesis revealed several characteristics of the examined galaxies, which appear to be mostly in good agreement with the current scientific consensus and also previous studies. However in order to derive more confident results it is vital to include in the future a robust error budget. As a consequence the impacts of galactic extinction as well as a potential dust presence in the observed galaxies could be better estimated.
Nevertheless a clear trend has been found for the effective radii to be $\sim 25$ percent smaller in the infrared $H_{160}$ band than in the optical $r_{625}$ filter, while the Sérsic index $n$ can be assumed constant over this range. The medians of $n$ over the whole examined wavelength range amounts to $n \sim 4.44 \pm 0.19$ for the MACS 1206 sample and $4.05 \pm 0.29$ for the MACS 0416 selection. These values turn out to be smaller than the median of $\langle n\rangle \sim 6 \mathrm{La}$ Barbera et al. (2010) found for local galaxies. However their analysis commenced from the $g$-band and extended down to the $K$-band, which could also explain their $r_{e}$ assessments in these two bands decreasing by 35 percent. On the other hand Vulcani et al. (2014) likewise found an $\langle n\rangle \sim 4$, which can be considered constant over all filters, by also using the multi-band capabilities of MegaMorph for their local galaxy sample ranging from $u$ to $K$ and being provided through the GAMA survey. Additionally they also detected a decrease of $\sim 40$ percent for $r_{e}$ going from $u$ to $H$. Similar trends have been found by Kelvin et al. (2012), who established a 38 percent reduction in $r_{e}$ for the same sample as well but by utilizing only single component fits. Finally Kennedy et al. (2015) confirms these results with examining the data supplied by GAMA II.
The same trends are discovered for early-type galaxies at intermediate redshifts, similar to the ones treated in this thesis. La Barbera et al. $(2002,2003)$ obtain comparable ratios of $r_{e}$ and $n$ for approximately 270 galaxies in three different clusters lying between $z=0.21$ and $z=0.64$. Utilizing filters corresponding to $U, V$ and $H$ in rest-frame their analyses reveal $r_{e, U V} / r_{e, o p t}=1.2 \pm 0.05, r_{e, o p t} / r_{e, I R}=1.26 \pm 0.06$ and $n_{U V} / n_{\text {opt }}=1.0 \pm 0.1, n_{\text {opt }} / n_{I R}=0.88 \pm 0.03$. Those determinations are in
excellent agreement with the assessments of this thesis and also consistent with the local values. Finally also for passive cluster galaxies at $z=1.39$ those trends are still valid, as Chan et al. (2016) derived in a range between $i_{775}$ and $H_{160}$ a $\sim 20$ percent decrease in size for those objects.

The decrease in $r_{e}$ with $n$ remaining constant over wavelength can be explained by intrinsic negative color gradients being inherent in the early-type galaxies. Hence the bulk of the treated galaxies possess redder centers and a bluer stellar population in the exterior regions. According to the derived decrease of $r_{e}$ the majority of the in this thesis examined galaxies show indeed negative color gradients with the medians of the gradients in four different colors being negative except for one, ranging between $-0.05 \pm 0.02$ and $-0.18 \pm 0.01 \mathrm{mag} \mathrm{dex}^{-1}$ for the entire data sample containing all galaxies.

These outcomes are perfectly consistent with the findings of Saglia et al. (2000), who also derived the color gradients for $z \sim 0.4$ cluster galaxies in the rest-frame $U-B, U-V$ and $B-V$ colors, as all their objects have at least in one color a negative gradient in the same range presented in this thesis. Other studies treating galaxies at similar redshifts, like Tamura \& Ohta (2000) with finding a median gradient of $\nabla\left(\lambda_{B, 450}-\lambda_{I, 814}\right)=-0.23 \pm 0.05 \mathrm{mag} \mathrm{dex}^{-1}$, La Barbera et al. (2002, 2003) derive likewise similar values for color gradients in elliptical galaxies and confirm the determinations. The values found in this thesis are especially consistent with the mean color gradients obtained by La Barbera et al. (2003) with $\nabla\left(\lambda_{U V}-\lambda_{\text {opt }}\right)=$ $-0.18 \pm 0.04 \mathrm{mag} \mathrm{dex}^{-1}$ and $\nabla\left(\lambda_{\text {opt }}-\lambda_{I R}\right)=-0.4 \pm 0.1 \mathrm{mag} \mathrm{dex}^{-1}$.

Similar results have already been described by Peletier et al. (1990a,b), who found mean color gradients of $-0.20 \pm 0.02 \mathrm{mag}^{\mathrm{dex}}{ }^{-1}$ in $U-R$, respectively $-0.09 \pm 0.02$ in $B-R$ for local early-type galaxies and also by La Barbera \& de Carvalho (2009), who expanded the analysis to seven different colors ranging from $g-r$ to $g-K$.

Many of the newest studies focus on higher redshift elliptical galaxies detecting on average steeper negative gradients than for local and intermediate- $z$ galaxies. Gargiulo et al. (2012) found gradients between -0.1 $\pm 0.1$ and $-1.0 \pm 0.1 \mathrm{mag} \mathrm{dex}^{-1}$ for their sample of early-type galaxies at $1.0<z<1.9$. Also Guo et al. (2011) reported for their $z \sim 2$ galaxies steeper color gradients, as well as De Propris et al. (2015), who derived a median color gradient of $-0.25 \mathrm{mag}_{\mathrm{dex}}{ }^{-1}$ for their highredshift ( $\langle z\rangle \sim 1.25$ ) selection. Finally very recently Chan et al. (2016) found a median value of $\sim 0.45 \mathrm{mag}^{\mathrm{dex}}{ }^{-1}$ for passive cluster galaxies at $z \sim 1.39$, which corresponds to being a factor 2 steeper than the local values.

Furthermore no correlation of the color gradients with cluster-centric distance, magnitude nor stellar mass can be found confirming that the variations in stellar populations are independent of those galactic properties. Saglia et al. (2000) came to the same conclusions in regard to the magnitudes, as well as Guo et al. (2011) and Gargiulo et al. (2012) concerning the stellar mass of the galaxies.

Interpreting the yield of the SSP models performed in the course of this thesis the origins of the color gradients, can be explained by a combination of age and metallicity gradients. The age differences from the innermost region to the outskirts of the galaxies appear rather shallow and add up on median to $1.2 \pm 0.4$ for MACS 1206 respectively $3.0 \pm 0.6$ Gyr for the less accurate modelled MACS 0416 sample, if the determinations at $0.1 r_{e}$ are not considered. The metallicity gradients are characterized for both clusters to be $\sim-0.2$ dex per decade in radius. Similar results are found by Saglia et al. (2000), La Barbera et al. (2003) and Tamura et al. (2000). Although the age variations found in this thesis appear to be stronger than the ones found in these publications, they are still well within the error estimates of these studies, especially considering that the age determinations done in this work required the results of the less precise $r_{625}$ band, in particular regarding the assessments in MACS 0416. Hence this confirms that variations in metallicity are the dominant reason for the observed color gradients in elliptical galaxies.

Consequently all this coherent findings including the ones presented in this thesis can be used to constrain possible evolutionary scenarios. Early-type galaxies showing a decrease in $r_{e}$ with wavelength and therefore also negative color gradients display differences in stellar populations. The order of these variations can support and disprove the different theories of galactic developments. In this context a monolithic collapse scenario would require much steeper negative gradients than observed, with logarithmic metallicity gradients between -0.5 and -1 dex per radial decade, due to the continuing inflow of metal rich gas towards the center during the formation process. As a result of this gas enrichment the younger stellar population must be situated in the interior, innermost regions of a galaxy, with ages of the stars increasing with radius. Hence a positive age gradient should be detectable, but rather the exact opposite is the case with the elliptical galaxies showing rather negative age gradients.

Also the assumption of gas rich mergers playing an important role in the recent evolution ( $z \lesssim 0.5-1$ ) can be omitted, as this would result into an old, red and extended component superimposed by a younger and bluer stellar population with
an in comparison smaller $r_{e}$. Also this would contradict the similar determinations of $n$ and $r_{e}$ for local galaxies and those at intermediate redshifts. The observations in color variations as well as the findings of decreasing age and metallicity with radius support rather the hierarchical scenario with dry mergers contributing to an insideout growth of the primary early-type galaxy. During this process the stars of lower mass galaxies with a younger and metal poorer population get accreted and settle into a more extended component compared to the original stellar distribution. This does not only lead to a significant increase in radius, whereas the mass is to a lesser extent enlarged, but also to a dilution of the age and metallicity variations, which subsequently match the observed gradients at different redshifts. This assumption is supported by the observations of very compact but massive quiescent galaxies at $z \sim 2$, with their size estimates being not an effect of errors in the modelling process and a resulting absence of flux detections at larger radii (Davari et al., 2014, and references therein).

## 8. Summary

In this thesis a total of 117 early-type galaxies at $z \sim 0.4$ are treated, with 79 belonging to the MACS 1206 cluster, while the remaining sources are part of the MACS 0416 structure. The photometric data is provided by the $C L A S H$ survey in 12 bands ranging from $B_{435}$ to $H_{160}$, while the spectroscopic measurements and hence also the membership determinations are supplied from the follow-up program $C L A S H-V L T$.

After determining an initial membership selection for the respective clusters, 2D models are produced for the observations in the 12 different filters utilizing the tool MegaMorph. A subsequent visual selection combined with the choice to only consider galaxies with $n>2.5$ yields the final sample of elliptical galaxies. Using the structural parameters being delivered by the models, surface brightness profiles are created in order to produce in turn the radial profiles for the colors $g_{475}-I_{814}$, $r_{625}-Y_{105}, I_{814}-H_{160}$ and $Y_{105}-H_{160}$. Consequently the color gradients are derived as the logarithmic slopes of these color profiles by either approximating an lsq fit or calculating the ratios of the effective radii in both respective bands. Finally the contributions of age and metallicity to the color gradients are analyzed by employing BC03 SSP models each time adopting a Chabrier IMF, with the models being based on different formation redshifts, ranging from $z_{f}=0.5$ to $z_{f}=10$ and metallicities with a super-solar value $Z=0.05$ at the one end and sub-solar $Z=0.008$ on the other one. An ensuing comparison between the simulated colors at the cluster redshift with the observed ones, delivers the aimed for constraints.

Although better error estimates are needed to obtain improved insights on the impacts of dust and extinction, as well as better constraints on absolute values, the relative results acquired for elliptical galaxies in the course of this thesis are consistent with those of previous studies:

- For the whole sample the effective radius decreases on median with wavelength, with $r_{e, 160} \sim 0.75 r_{e, 625}$, describing already indirectly a negative color gradient, due to the presence of distinct stellar populations at different radii.
- The Sérsic index $n$ can be assumed on median constant over the same wavelength range, restricting already the potential development scenarios to those, which leave $n$ unchanged.
- The majority of the galaxies and the overall median show negative color gradients lying between $-0.05 \pm 0.02$ and $-0.18 \pm 0.01 \mathrm{mag} \mathrm{dex}^{-1}$, indicating a redder, older stellar population in the centers and bluer, younger stars in the outskirts. The results are independent of the derivation method, lsq fit or calculation, and also differ insignificantly between the employed colors.
- No correlation or dependence of the color gradients on stellar mass, magnitude or cluster-centric distance can be found, implying that all those properties are not directly responsible for the radial color variations.
- Age as well as more dominantly metallicity are driving the color gradients. While the centers appear to be $\sim 1$ (MACS 1206) to $\sim 3$ (MACS 0416) Gyr older than the exterior regions, the metallicity drops from super-solar $Z \sim 0.03$ in the innermost parts to values similar to the solar metallicity in the outer domains.
- Considering all those observations the most probable evolutionary scenario for elliptical galaxies is that they already assemble a majority of their stellar mass at very high redshifts ( $z \gtrsim 2$ ), followed by a passive evolution of the initial stellar population. With time the galaxies undergo several dry and minor mergers, resulting in a much more notable increase in size than in mass, as the metal poorer and younger populations of the accreted galaxies settle into more extended structures.

The options to improve and/or broaden the determinations done in this thesis are manifold. One choice may be to distinguish the galaxies into late- and earlytype galaxies not only by morphological criteria, but also to consider spectroscopic measurements. Consequently a much better differentiation even within those two main branches would be possible, offering the chance to examine potential trends with morphological type. In general it would also be intriguing to determine the results for field galaxies and to analyze the differences between galaxies in diverging environments. Finally an expansion of these investigations to different redshifts, in particular going to greater distances, could deliver compelling new insights on the formation and evolution of galaxies.

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## Appendix

## A. Source Extractor input file

```
# Default configuration file for SExtractor 2.8.6
# ____________________log
CATALOGNAME /no_need_for_a_name.cat
# Will be overwritten anyway! name of the output catalog
CATALOG_TYPE ASCII_HEAD
# NONE, ASCII, ASCII_HEAD, ASCII_SKYCAT, ASCIL_VOTABLE, FITS_1.0
#or FITS_LDAC
PARAMETERSNAME / no_need.params
# Will be overwritten anyway! name of the file
#containing catalog contents
#-_ Extraction
DETECT_TYPE CCD
# CCD (linear) or PHOTO (with gamma correction)
DETECT_MINAREA 80
# minimum number of pixels above threshold default: 100
THRESHLTYPE RELATIVE
# threshold type: RELATIVE (in sigmas) or ABSOLUTE (in ADUs)
DETECT_THRESH 2.5
# < <sigmas> or <threshold >,<ZP> in mag.arcsec-2 #default: 2.8
2 ANALYSIS_THRESH 2.5
23 # <sigmas> or <threshold >,<ZP> in mag. arcsec -2 #default: 2.8
2 4 ~ F I L T E R ~ Y ~
25 # apply filter for detection (Y or N)?
6 FILTER_NAME tophat_3.0_3x3.conv
27 # name of the file containing the filter
8 FILTER_THRESH
29 # Threshold[s] for retina filtering
0 DEBLEND_NTHRESH 32
# # Number of deblending sub-thresholds #default: 64
DEBLEND_MINCONT 0.002
33 # Minimum contrast parameter for deblending
```

```
CLEAN Y
# Clean spurious detections? (Y or N)?
CLEAN_PARAM 0.5
# Cleaning efficiency
MASK_TYPE CORRECT
# type of detection MASKing: can be one of NONE, BLANK or CORRECT
#-__ Weighting
WEIGHT_TYPE NONE
# type of WEIGHTing: NONE, BACKGROUND, MAP_RMS,
#MAP_VAR or MAP_WEIGHT
WEIGHT_IMAGE /no_need.fits
# Will be overwritten anyway! weight-map filename
WEIGHT_GAIN Y
# modulate gain (E/ADU) with weights? (Y/N)
WEIGHT_THRESH
# weight threshold[s] for bad pixels
```



```
FLAG_IMAGE flag.fits
# filename for an input FLAG-image
FLAG_TYPE OR
# flag pixel combination: OR, AND, MIN, MAX or MOST
#- Photometry
PHOT_APERTURES 60
# MAGAPER aperture diameter(s) in pixels
PHOT_AUTOPARAMS 2.5, 3.5
# MAG_AUTO parameters: <Kron_fact>,<min_radius>
PHOT_PETROPARAMS 2.0, 3.5
# MAG_PETRO parameters: <Petrosian_fact >, <min_radius>
PHOT_AUTOAPERS 0.0,0.0
#<estimation >,<<measurement> minimum apertures for MAGAUTO
#and MAGPETRO
PHOT_FLUXFRAC 0.5
# flux fraction[s] used for FLUX_RADIUS
SATURLEVEL 60000.0
# level (in ADUs) at which arises saturation
#SATUR_KEY NOEXIT
# keyword for saturation level (in ADUs) - don't trust SATURATE
MAG_ZEROPOINT 25.9433
# magnitude zero-point
MAGGAMMA 4.0
```

```
7 # gamma of emulsion (for photographic scans)
GAIN 2.0
# detector gain in e-/ADU
#GAIN_KEY GAIN
# keyword for detector gain in e-/ADU
PIXEL_SCALE 0.065
# size of pixel in arcsec ( 0=use FITS WCS info)
#_ Star/Galaxy Separation
SEEING_FWHM 0.1
# stellar FWHM in arcsec
STARNNWNAME default.nnw
# Neural-Network_Weight table filename
#___ Background
BACK_TYPE AUTO
# AUTO or MANUAL
#BACKTYPE MANUAL
# AUTO or MANUAL
BACKVALUE 3951.
# Default background value in MANUAL mode
BACK_SIZE 214
# Background mesh: <size> or <width>,<height>
BACK_FILTERSIZE 5
# Background filter: <size> or <width>,<height>
BACKPHOTO_TYPE LOCAL
# can be GLOBAL or LOCAL
BACKPHOTO_THICK 60
# thickness of the background LOCAL annulus
BACK_FILTTHRESH 0.0
# Threshold above which the background-map filter operates
#-_ Check Image
CHECKIMAGE_TYPE NONE
# can be NONE, BACKGROUND, BACKGROUND_RMS, MINIBACKGROUND,
# MINIBACK_RMS, -BACKGROUND, nFILTERED, OBJECTS, -OBJECTS,
# SEGMENTATION or APERTURES
CHECKIMAGENAME / check.fits
# Filename for the check-image
#-_ Memory (change with caution!)
MEMORY_OBJSTACK 30000
# 3000 number of objects in stack
```

Appendix

```
MEMORYPIXSTACK 10000000
# 1200000 number of pixels in stack - x4 from default
MEMORY_BUFSIZE 4096
# 1024 number of lines in buffer
#-_ Association
ASSOCNAME sky.list
# name of the ASCII file to Associate
ASSOC_DATA 2,3,4
# columns of the data to replicate (0= all)
ASSOC_PARAMS 2,3,4
# columns of xpos,ypos[,mag]
ASSOC_RADIUS 2.0
# cross-matching radius (pixels)
ASSOC_TYPE NEAREST
# Association method: FIRST, NEAREST, MEAN, MAG_MEAN, SUM,
# MAG_SUM, MIN or MAX
ASSOCSELEC_TYPE MATCHED
# ASSOC selection type: ALL, MATCHED or -MATCHED
#-_M_ Miscellaneous
VERBOSE_TYPE QUIET
# can be QUIET, NORMAL or FULL
WRITEXML N
# Write XML file (Y/N)?
XMLNAME sex.xml
# Filename for XML output
XSL_URL file:///usr/share/sextractor/sextractor.xsl
# Filename for XSL style-sheet
NTHREADS 1
# Number of simultaneous threads for the SMP version of SExtractor
# 0 = automatic
FITS_UNSIGNED N
# Treat FITS integer values as unsigned (Y/N)?
INTERP_MAXXLAG 16
# Max. lag along X for 0-weight interpolation
INTERP_MAXYLAG 16
# Max. lag along Y for 0-weight interpolation
INTERP_TYPE ALL
# Interpolation type: NONE, VAR_ONLY or ALL
```


## B. MegaMorph setup file

```
1 # FILE LOCATIONS
A00) /home/marian/data/member_run_1/images.gala
#file containing ALL input files, including SExtractor
#(1st line) and all bands, weight,...)
A01) /home/marian/data/member_run_1/output_tier_1
#output directory for catalogues
#
# SEXTRACTOR SETUP\
B00) execute
#execute the SExtractor block
B01) /home/marian/programs/sextractor/usr/bin/sex
#SExtractor executable including path
B02) /home/marian/data/member_run_1/gala.param
#output parameters in .param-format
B03) /home/marian/data/member_run_1/M1206_cold.sex
#SExtractor setup file (cold),
#it has to properly deblend all brighter objects
B04) coldcat
#output catalogue (cold)
B05) coldseg.fits
#output segmentation map (cold)
B06 ) /home/marian/data/member_run_1/M1206_hot.sex
#SExtractor setup file (hot), tuned to pick up faint objects
B07) hotcat
#output catalogue (hot)
B08) hotseg.fits
#output segmentation map (hot)
B09) 1.05
#enlarge the cold isophotes for catalogue combination by a factor.
#1.1 = by 10%. If the source in hot mode lies "sufficiently"
#outside of all cold Kron ellipses, it enters the output catalogue
B10) outcat
#output combined catalogue
B11) outseg.fits
#output combined segmentation map
B12) outparam
#output parameter file
B13) check.fits
9 #check image filename
0 B14) apertures
1 #check image type
```

```
B15) none
# ('rms' = true) is the SEXtractor 'weight' map a RMS map
#that should be used for SExtractor?
B16) none
#list of positions to be excluded from the individual SExtractor
#catalogues (to remove bad detections; format: x y)
B17) 1.5
#radius in pix used to exclude objects: gala rejects any
#detection within a certain radius B16 automatically from the
#SEx catalogue on a subsequent run of the code.
#Thus, if one wants to refine the cat, the sex section of the
#code has to run twice, i.e. gala needs to be started first
#with only the sex section activated and then run a 2nd time with
#the SEx section and optionally others enables as well.
#the first exec is for identifying bad detections,
#the 2nd run then treats them.
B18)
#if set "outonly": hot/cold catalogues/segmaps are deleted,
#else: all catalogues/segmaps are kept
B19) none
#list of positions to be excluded from the final
B20) combined_se_cat
#only gets created with C)
#combined sextractor catalogue put in A01)
#
# STAMP SETUP=
C00) execute
#execute the Stamps creation block
C01) stamps
#descriptor file for postage stamps
C02 )
#postposition for postage stamps (only used for single image run,
#NOT for mutli-wavelength) #preposition for postage stamps
#only works if left blank! (can't find rms image otherwise)
C03) 1.5
#scale factor by which the sextractor isophotes are enlarged
#
# SKY PREPARATION SETUP
D00) execute
#execute the sky preparation block
#always perform D, E and F together
D01) skymap
#output object/sky-mapfile
```

```
D02) outsky
#output filename for sky values
D03) 1.1
#scale factor by which SEx isophote is enlarged (for skymap)
D04) 1.1
#scale factor by which SEx isophote is enlarged (for neighbours)
D05) }1
#additional offset to scale factor
D06) 15
#distance between individual sky isophotes
D07) 30
#width of individual sky isophotes
D08) 15
#gap between sextractor isophote and inner sky isophote
D09) 2.5
#cut below which objects are considered as contributing
D10) 3
#nobj_max; max number of allowed contributing sources
D11) 1.4
#power by which the flux_radius is raised to convert to Re
D12) 15
#calculate the slope of the sky from the x last determinations
D13) -0.25
#slope in fwhm_image vs. mag_best below which object is star
D14) 2.75
# here: Flux-radius vs. Ic_aper
#zeropoint in fwhm_image vs. mag_best below which object is star
D15) 5
#magnitude faint end limit for secondaries when fitting galaxies
D16) }
#magnitude faint end limit for secondaries when fitting stars
D17) 3
#number of neighbouring frames (see text for details)
D18) 16 #!!
#maximum number of parallel processes (see text for details)
D19) }20
#minimum distance (in arcseconds) between sources processed
D20) 60
#minimum distance (in arcseconds) to already blocked objects
#(to make sure that no faint object next to a very bright one is
# fit. standard value: 0.3*D20 )
D21) /home/marian/data/member_run_1/all_members.coords
#optional list containing primary targets
```

```
D22) 1.0
#search/correlation radius for the above list [arcsec]
# GALFIT SETUP=#
E00) /home/marian/programs/megamorph/galfitm/galfitm - 1.2.0-linux -x86_64
#Galfit executable including path
E01 )
#filename for list of tiles in current batch (format: image)
#feature switched off when file does not exist, all images are done
E02) obj
#object file preposition
E03) galfit
#preposition for GALFIT output files
E04) /home/marian/data/original_data/psf/55 psf_f814w_norm.fits ,
/home/marian/data/original_data/psf/45 psf_f435w_norm.fits ,
/home/marian/data/original_data/psf/15 psf_f475w_norm.fits ,
/home/marian/data/original_data/psf/f606w.psfmos.fits ,
/home/marian/data/original_data/psf/55 psf_f625w_norm.fits ,
/home/marian/data/original_data/psf/55 psf_f775w_norm.fits ,
/home/marian/data/original_data/psf/55 psf_f850LPw_norm.fits ,
/home/marian/data/original_data/psf/15 psf_f105w_norm.fits ,
/home/marian/data/original_data/psf/15 psf_f110w_norm.fits ,
/home/marian/data/original_data/psf/15 psf_f125w_norm.fits ,
/home/marian/data/original_data/psf/15 psf_f140w_norm.fits ,
/home/marian/data/original_data/psf/15 psf_f160w_norm.fits
#PSF file names including paths
E05 ) mask
#mask file preposition
E06) constr
#constraint file preposition
E07) 150
#convolution box size
E08) 0
#zeropoint, #here without extinction correction
E09) 0.065
#plate scale of the images [arcsec/pixel]
E10) 1
#exposure time
E11) 400
#constraint max Re
E12) -5
#constraint min magnitude deviation (minus)
E13) 5
```

```
#constraint max magnitude deviation (plus)
E14) 0.2
#constraint min sersic index
E15) 8
#constraint max sersic index
E16)
#nice?
E17) 4.0
#GALFIT version string. E.g. 2.0.3c; 4.0 = GalfitM
E18) input,model,residual, psf
E19)}45
#time (in minutes) after which galfit processes are killed
#feature switched off when time =0
E20) 0,0,11,2,2,0,0
# x,y,mag,re,n,AR,PA
#order of Chebyshev polynomials in the individual parameters
#to be used in GALFIT
# 0 = constant over all wavelength,
#1 = linear over wavelength ,...
#band-1 = free
E21) galfit
# outputnfolder name for all galfit output files
E22) restrict
# restrict the maxiumum number of degrees of freedom to the
#number of band for an object? ("restrict" = true)
E23) }3
# fraction of data pixels ==0 which restricts the polynomial (in %),
# larger value allows more pixels to be 0 before taking action
E24) 3
# minimum number of good images that are needed for the fit
#(is set to be >1 by the code),
#e.g. objects without any good data are NEVER fit)
B/D DECOMPOSITION SETUP
F00)
# execute the B/D fitting block. Mostly assumes that single
#sersic fitting (block E) is also or had previsouly been run
F01) 0,0,11,0,-1,0,0
#similar to E18 for BULGE
F02) 0, 0, 11,0, - 1,0,0
#similar to E18 for DISK
F03) bdfix
# output label
```

Appendix

```
F04 )
#optional list containing primary targets
F05 ) 1.0
#search/correlation radius for the above list
F06 ) 40.0
F07) input, model, residual, component, psf
# OUTPUT CATALOGUE SETUP
G00) execute
#execute catalogue combination block (read single_sersic)
G01) execute
#execute catalogue combination block (read single-sersic & B/D)
G02) final_catalog.fits
#filename for output catalogue in A01)
```


## C. MegaMorph results

This section contains the results provided by the MegaMorph fitting process. The assessments are split into twelve tables separated by the parameters magnitude, $r_{e}$ and $n$, the seven ACS and five WFC3/IR bands and finally the two clusters MACS 1206 and MACS 0416. The uncertainties are the ones returned by GALFIT.
Table C.1.: Magnitude assessments in the ACS bands for the fitted early-type galaxies in MACS 1206.

| ID | MACS 1206 ACS AB Magnitudes |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{B}_{435}$ | $\mathrm{g}_{475}$ | $\mathrm{V}_{606}$ | $\mathrm{r}_{625}$ | 1775 | $\mathrm{I}_{814}$ | Z850 |
| 1 | $22.6 \pm 0.0027$ | $22.2 \pm 0.0041$ | $21.0 \pm 0.0015$ | $20.7 \pm 0.0015$ | $20.1 \pm 0.0008$ | $19.9 \pm 0.0007$ | $19.6 \pm 0.0004$ |
| 3 | $22.6 \pm 0.0033$ | $22.4 \pm 0.0021$ | $21.1 \pm 0.0019$ | $20.9 \pm 0.0019$ | $20.2 \pm 0.0016$ | $20.0 \pm 0.0013$ | $19.6 \pm 0.0009$ |
| 4 | $24.3 \pm 0.0067$ | $24.0 \pm 0.0057$ | $23.0 \pm 0.0036$ | $22.8 \pm 0.0032$ | $22.3 \pm 0.0017$ | $22.0 \pm 0.0014$ | $21.7 \pm 0.0021$ |
| 5 | $22.2 \pm 0.0343$ | $21.6 \pm 0.0193$ | $20.4 \pm 0.0096$ | $20.1 \pm 0.0089$ | $19.7 \pm 0.0054$ | $19.4 \pm 0.0045$ | $19.1 \pm 0.0035$ |
| 6 | $21.7 \pm 0.0009$ | $21.2 \pm 0.0006$ | $20.0 \pm 0.0005$ | $19.6 \pm 0.0004$ | $19.1 \pm 0.0003$ | $18.8 \pm 0.0003$ | $18.5 \pm 0.0003$ |
| 8 | $24.8 \pm 0.0884$ | $24.4 \pm 0.062$ | $23.3 \pm 0.0333$ | $23.2 \pm 0.0318$ | $22.5 \pm 0.0199$ | $22.3 \pm 0.0157$ | $22.0 \pm 0.0155$ |
| 10 | $22.2 \pm 0.0006$ | $21.7 \pm 0.0004$ | $20.5 \pm 0.0003$ | $20.1 \pm 0.0003$ | $19.6 \pm 0.0002$ | $19.3 \pm 0.0003$ | $19.0 \pm 0.0002$ |
| 11 | $23.0 \pm 0.0048$ | $22.5 \pm 0.0032$ | $21.5 \pm 0.0019$ | $21.1 \pm 0.002$ | $20.7 \pm 0.0015$ | $20.4 \pm 0.0014$ | $20.1 \pm 0.0017$ |
| 12 | $22.5 \pm 0.0$ | $22.1 \pm 0.0$ | $20.8 \pm 0.0$ | $20.5 \pm 0.0$ | $19.9 \pm 0.0$ | $19.6 \pm 0.0$ | $19.3 \pm 0.0$ |
| 15 | $22.9 \pm 0.0025$ | $22.4 \pm 0.0018$ | $21.5 \pm 0.0017$ | $21.1 \pm 0.0017$ | $20.7 \pm 0.0014$ | $20.4 \pm 0.0011$ | $20.1 \pm 0.0008$ |
| 17 | $22.0 \pm 0.0015$ | $21.4 \pm 0.0013$ | $20.2 \pm 0.001$ | $19.8 \pm 0.0011$ | $19.4 \pm 0.0005$ | $19.0 \pm 0.0004$ | $18.7 \pm 0.0002$ |
| 19 | $23.6 \pm 0.0563$ | $23.2 \pm 0.0419$ | $22.4 \pm 0.0293$ | $21.9 \pm 0.0269$ | $21.5 \pm 0.0197$ | $21.2 \pm 0.0165$ | $20.9 \pm 0.0118$ |
| 20 | $24.1 \pm 0.0096$ | $23.8 \pm 0.0097$ | $23.0 \pm 0.0059$ | $22.8 \pm 0.0056$ | $22.5 \pm 0.0024$ | $22.2 \pm 0.0023$ | $21.9 \pm 0.0018$ |
| 21 | $22.9 \pm 0.0074$ | $22.4 \pm 0.0064$ | $21.4 \pm 0.0045$ | $21.1 \pm 0.004$ | $20.6 \pm 0.0028$ | $20.4 \pm 0.002$ | $20.1 \pm 0.0013$ |
| 22 | $23.0 \pm 0.0043$ | $22.7 \pm 0.0042$ | $21.8 \pm 0.0035$ | $21.3 \pm 0.0033$ | $20.8 \pm 0.0021$ | $20.6 \pm 0.0016$ | $20.3 \pm 0.0011$ |
| 23 | $24.0 \pm 0.0022$ | $23.4 \pm 0.0021$ | $22.4 \pm 0.0018$ | $22.1 \pm 0.0016$ | $21.6 \pm 0.0011$ | $21.3 \pm 0.001$ | $21.0 \pm 0.0005$ |
| 24 | $24.2 \pm 0.0059$ | $23.6 \pm 0.0046$ | $22.5 \pm 0.0032$ | $22.1 \pm 0.0029$ | $21.6 \pm 0.0024$ | $21.3 \pm 0.0026$ | $21.1 \pm 0.0028$ |
| 25 | $23.6 \pm 0.0024$ | $23.1 \pm 0.0018$ | $21.9 \pm 0.0013$ | $21.5 \pm 0.0012$ | $21.0 \pm 0.0008$ | $20.7 \pm 0.0007$ | $20.4 \pm 0.0004$ |
| 26 | $23.6 \pm 0.0017$ | $22.9 \pm 0.0013$ | $21.7 \pm 0.0009$ | $21.3 \pm 0.001$ | $20.8 \pm 0.0005$ | $20.5 \pm 0.0005$ | $20.2 \pm 0.0004$ |
| 28 | $23.2 \pm 0.0034$ | $22.6 \pm 0.0012$ | $21.4 \pm 0.0007$ | $21.0 \pm 0.0006$ | $20.5 \pm 0.0004$ | $20.2 \pm 0.0003$ | $20.0 \pm 0.0003$ |
| 29 | $23.0 \pm 0.0012$ | $22.2 \pm 0.0011$ | $21.0 \pm 0.001$ | $20.6 \pm 0.0006$ | $20.2 \pm 0.0005$ | $19.8 \pm 0.0004$ | $19.6 \pm 0.0003$ |
| 30 | $22.6 \pm 0.005$ | $22.4 \pm 0.0044$ | $21.4 \pm 0.0038$ | $21.2 \pm 0.0036$ | $20.7 \pm 0.0031$ | $20.4 \pm 0.0023$ | $20.1 \pm 0.0019$ |
| 31 | $22.8 \pm 0.0043$ | $22.3 \pm 0.0038$ | $21.3 \pm 0.0022$ | $20.8 \pm 0.0019$ | $20.3 \pm 0.0012$ | $20.1 \pm 0.0011$ | $19.7 \pm 0.0008$ |


| ID | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{z}_{850}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $22.4 \pm 0.0032$ | $21.7 \pm 0.0028$ | $20.5 \pm 0.0023$ | $20.1 \pm 0.0021$ | $19.6 \pm 0.0012$ | $19.3 \pm 0.0011$ | $19.0 \pm 0.0005$ |
| 32 | 22.4 |  |  |  |  |  |  |
| 33 | $24.2 \pm 0.0505$ | $23.8 \pm 0.0404$ | $22.6 \pm 0.0297$ | $22.4 \pm 0.0273$ | $21.7 \pm 0.0196$ | $21.5 \pm 0.0165$ | $21.2 \pm 0.0127$ |
| 35 | $24.0 \pm 0.0026$ | $23.5 \pm 0.0022$ | $22.3 \pm 0.0017$ | $21.9 \pm 0.0015$ | $21.4 \pm 0.0005$ | $21.1 \pm 0.0005$ | $20.8 \pm 0.0011$ |
| 37 | $23.4 \pm 0.0702$ | $23.0 \pm 0.0575$ | $22.2 \pm 0.0398$ | $21.8 \pm 0.0356$ | $21.3 \pm 0.0232$ | $21.0 \pm 0.0197$ | $20.7 \pm 0.0144$ |
| 38 | $23.8 \pm 0.003$ | $23.5 \pm 0.0026$ | $22.8 \pm 0.002$ | $22.4 \pm 0.0019$ | $21.9 \pm 0.0019$ | $21.7 \pm 0.0012$ | $21.4 \pm 0.0009$ |
| 39 | $24.0 \pm 0.0061$ | $23.3 \pm 0.0046$ | $22.3 \pm 0.0034$ | $21.8 \pm 0.0034$ | $21.3 \pm 0.0021$ | $21.0 \pm 0.002$ | $20.8 \pm 0.001$ |
| 40 | $23.8 \pm 0.0451$ | $23.4 \pm 0.0319$ | $22.3 \pm 0.0185$ | $22.0 \pm 0.0173$ | $21.4 \pm 0.0119$ | $21.2 \pm 0.0101$ | $20.9 \pm 0.0076$ |
| 41 | $23.9 \pm 0.0014$ | $23.4 \pm 0.0012$ | $22.3 \pm 0.0007$ | $21.9 \pm 0.0012$ | $21.3 \pm 0.0004$ | $21.0 \pm 0.0002$ | $20.7 \pm 0.0004$ |
| 42 | $23.1 \pm 0.0002$ | $22.6 \pm 0.0002$ | $21.3 \pm 0.0001$ | $20.9 \pm 0.0001$ | $20.4 \pm 0.0001$ | $20.1 \pm 0.0001$ | $19.8 \pm 0.0001$ |
| 43 | $21.4 \pm 0.0012$ | $20.9 \pm 0.0009$ | $19.7 \pm 0.0005$ | $19.4 \pm 0.0005$ | $18.9 \pm 0.0004$ | $18.6 \pm 0.0002$ | $18.3 \pm 0.0002$ |
| 45 | $23.4 \pm 0.0026$ | $22.9 \pm 0.0021$ | $21.8 \pm 0.0016$ | $21.4 \pm 0.0014$ | $21.0 \pm 0.0011$ | $20.7 \pm 0.001$ | $20.4 \pm 0.001$ |
| 46 | $23.4 \pm 0.0019$ | $22.9 \pm 0.0017$ | $21.8 \pm 0.0014$ | $21.4 \pm 0.0014$ | $20.9 \pm 0.0012$ | $20.6 \pm 0.001$ | $20.4 \pm 0.0008$ |
| 47 | $21.9 \pm 0.0016$ | $21.5 \pm 0.0012$ | $20.1 \pm 0.0008$ | $19.9 \pm 0.0008$ | $19.2 \pm 0.0005$ | $19.0 \pm 0.0005$ | $18.7 \pm 0.0004$ |
| 48 | $23.3 \pm 0.0017$ | $23.0 \pm 0.0014$ | $22.1 \pm 0.0011$ | $21.7 \pm 0.001$ | $21.1 \pm 0.0007$ | $20.8 \pm 0.0006$ | $20.5 \pm 0.0005$ |
| 50 | $23.7 \pm 0.0012$ | $23.3 \pm 0.0005$ | $22.0 \pm 0.0003$ | $21.7 \pm 0.0002$ | $21.2 \pm 0.0$ | $20.9 \pm 0.0$ | $20.6 \pm 0.0001$ |
| 51 | $22.4 \pm 0.0004$ | $21.8 \pm 0.0001$ | $20.7 \pm 0.0001$ | $20.3 \pm 0.0001$ | $19.8 \pm 0.0$ | $19.5 \pm 0.0$ | $19.2 \pm 0.0$ |
| 52 | $21.9 \pm 0.0004$ | $21.3 \pm 0.0002$ | $20.1 \pm 0.0002$ | $19.7 \pm 0.0001$ | $19.2 \pm 0.0$ | $18.9 \pm 0.0001$ | $18.6 \pm 0.0001$ |
| 55 | $23.5 \pm 0.0027$ | $23.1 \pm 0.0025$ | $22.0 \pm 0.0021$ | $21.7 \pm 0.002$ | $21.2 \pm 0.0014$ | $20.9 \pm 0.0012$ | $20.6 \pm 0.0007$ |
| 56 | $24.1 \pm 0.0022$ | $23.7 \pm 0.0016$ | $22.9 \pm 0.0009$ | $22.5 \pm 0.0009$ | $21.9 \pm 0.0006$ | $21.7 \pm 0.0005$ | $21.4 \pm 0.0004$ |
| 57 | $21.8 \pm 0.001$ | $21.3 \pm 0.0008$ | $20.2 \pm 0.0006$ | $19.8 \pm 0.0004$ | $19.3 \pm 0.0003$ | $19.1 \pm 0.0002$ | $18.8 \pm 0.0001$ |
| 58 | $24.1 \pm 0.0134$ | $23.5 \pm 0.0106$ | $22.4 \pm 0.0082$ | $22.0 \pm 0.0075$ | $21.5 \pm 0.0057$ | $21.2 \pm 0.005$ | $20.9 \pm 0.0044$ |
| 59 | $24.2 \pm 0.0094$ | $23.6 \pm 0.0053$ | $22.6 \pm 0.0027$ | $22.2 \pm 0.0027$ | $21.6 \pm 0.0018$ | $21.3 \pm 0.0019$ | $21.1 \pm 0.0011$ |
| 61 | $23.8 \pm 0.0307$ | $23.3 \pm 0.0201$ | $22.1 \pm 0.0105$ | $21.7 \pm 0.0092$ | $21.2 \pm 0.006$ | $20.9 \pm 0.0049$ | $20.6 \pm 0.0042$ |
| 62 | $23.9 \pm 0.0056$ | $23.5 \pm 0.0052$ | $22.6 \pm 0.004$ | $22.2 \pm 0.0038$ | $21.7 \pm 0.0031$ | $21.4 \pm 0.0021$ | $21.1 \pm 0.0015$ |


| ID | MACS 1206 ACS AB Magnitudes cont. |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{B}_{435}$ | $\mathrm{g}_{475}$ | $\mathrm{V}_{606}$ | $\mathrm{r}_{625}$ | 1775 | $\mathrm{I}_{814}$ | Z850 |
| 63 | $25.0 \pm 0.0718$ | $24.4 \pm 0.0485$ | $23.2 \pm 0.0242$ | $22.8 \pm 0.0217$ | $22.3 \pm 0.012$ | $21.9 \pm 0.0107$ | $21.6 \pm 0.0089$ |
| 64 | $23.7 \pm 0.0119$ | $23.2 \pm 0.0108$ | $22.4 \pm 0.0089$ | $21.9 \pm 0.0082$ | $21.6 \pm 0.0063$ | $21.3 \pm 0.0042$ | $21.1 \pm 0.0025$ |
| 65 | $23.3 \pm 0.0066$ | $23.1 \pm 0.005$ | $21.8 \pm 0.0034$ | $21.5 \pm 0.0029$ | $21.0 \pm 0.0026$ | $20.7 \pm 0.002$ | $20.5 \pm 0.0015$ |
| 68 | $23.3 \pm 0.0139$ | $22.7 \pm 0.0102$ | $21.6 \pm 0.0074$ | $21.2 \pm 0.0067$ | $20.7 \pm 0.0038$ | $20.5 \pm 0.0036$ | $20.2 \pm 0.0022$ |
| 69 | $22.7 \pm 0.003$ | $22.3 \pm 0.0025$ | $21.2 \pm 0.0022$ | $20.9 \pm 0.002$ | $20.5 \pm 0.0015$ | $20.2 \pm 0.0014$ | $19.9 \pm 0.0013$ |
| 70 | $23.1 \pm 0.0052$ | $22.7 \pm 0.0041$ | $21.7 \pm 0.003$ | $21.3 \pm 0.0027$ | $20.8 \pm 0.0015$ | $20.5 \pm 0.0012$ | $20.3 \pm 0.0008$ |
| 71 | $23.3 \pm 0.0009$ | $22.8 \pm 0.0008$ | $21.8 \pm 0.0006$ | $21.4 \pm 0.0006$ | $20.9 \pm 0.0004$ | $20.6 \pm 0.0004$ | $20.4 \pm 0.0004$ |
| 72 | $23.7 \pm 0.0061$ | $23.3 \pm 0.0051$ | $22.3 \pm 0.0035$ | $22.0 \pm 0.0041$ | $21.4 \pm 0.0021$ | $21.1 \pm 0.0018$ | $20.9 \pm 0.0013$ |
| 73 | $24.4 \pm 0.0099$ | $23.9 \pm 0.0059$ | $22.8 \pm 0.0033$ | $22.5 \pm 0.0029$ | $22.0 \pm 0.0028$ | $21.7 \pm 0.0013$ | $21.4 \pm 0.001$ |
| 74 | $24.4 \pm 0.0574$ | $24.1 \pm 0.0395$ | $23.0 \pm 0.0254$ | $22.9 \pm 0.0233$ | $22.2 \pm 0.0166$ | $22.0 \pm 0.0131$ | $21.8 \pm 0.0112$ |
| 76 | $22.6 \pm 0.0014$ | $22.0 \pm 0.0011$ | $20.7 \pm 0.001$ | $20.3 \pm 0.0009$ | $19.8 \pm 0.0008$ | $19.5 \pm 0.0008$ | $19.3 \pm 0.0007$ |
| 77 | $22.3 \pm 0.0307$ | $22.0 \pm 0.0262$ | $21.3 \pm 0.018$ | $20.9 \pm 0.016$ | $20.5 \pm 0.009$ | $20.2 \pm 0.009$ | $19.8 \pm 0.0061$ |
| 80 | $24.9 \pm 0.0044$ | $24.4 \pm 0.0033$ | $23.4 \pm 0.0021$ | $23.1 \pm 0.0019$ | $22.7 \pm 0.0013$ | $22.4 \pm 0.0011$ | $22.1 \pm 0.0011$ |
| 81 | $23.9 \pm 0.0075$ | $23.5 \pm 0.0065$ | $22.6 \pm 0.0045$ | $22.3 \pm 0.0043$ | $21.9 \pm 0.0027$ | $21.6 \pm 0.0022$ | $21.3 \pm 0.0019$ |
| 82 | $24.6 \pm 0.0064$ | $24.2 \pm 0.0051$ | $23.4 \pm 0.0038$ | $23.0 \pm 0.0034$ | $22.4 \pm 0.0025$ | $22.2 \pm 0.0022$ | $21.9 \pm 0.0016$ |
| 83 | $23.6 \pm 0.034$ | $22.9 \pm 0.0242$ | $21.8 \pm 0.0164$ | $21.5 \pm 0.0151$ | $20.9 \pm 0.0102$ | $20.6 \pm 0.0088$ | $20.3 \pm 0.0064$ |
| 84 | $24.2 \pm 0.001$ | $23.6 \pm 0.0026$ | $22.4 \pm 0.0005$ | $22.0 \pm 0.0004$ | $21.4 \pm 0.0003$ | $21.1 \pm 0.0003$ | $20.8 \pm 0.0002$ |
| 85 | $23.8 \pm 0.0079$ | $23.3 \pm 0.0062$ | $22.4 \pm 0.0035$ | $21.9 \pm 0.0031$ | $21.4 \pm 0.0023$ | $21.2 \pm 0.0018$ | $20.9 \pm 0.0018$ |
| 87 | $23.8 \pm 0.0042$ | $23.1 \pm 0.0034$ | $22.0 \pm 0.0024$ | $21.5 \pm 0.0021$ | $21.0 \pm 0.0016$ | $20.8 \pm 0.0016$ | $20.5 \pm 0.0016$ |
| 88 | $23.9 \pm 0.0309$ | $23.4 \pm 0.0204$ | $22.3 \pm 0.0095$ | $21.9 \pm 0.0087$ | $21.3 \pm 0.0052$ | $21.0 \pm 0.0047$ | $20.7 \pm 0.0038$ |
| 89 | $24.1 \pm 0.0039$ | $23.6 \pm 0.0026$ | $22.6 \pm 0.0018$ | $22.2 \pm 0.0016$ | $21.7 \pm 0.0014$ | $21.4 \pm 0.0009$ | $21.2 \pm 0.0006$ |
| 90 | $24.5 \pm 0.0453$ | $24.0 \pm 0.0313$ | $22.9 \pm 0.0174$ | $22.5 \pm 0.0157$ | $22.0 \pm 0.0104$ | $21.7 \pm 0.0086$ | $21.4 \pm 0.0066$ |
| 91 | $24.4 \pm 0.003$ | $24.1 \pm 0.0016$ | $22.9 \pm 0.0011$ | $22.7 \pm 0.001$ | $22.1 \pm 0.0007$ | $21.9 \pm 0.0006$ | $21.6 \pm 0.0004$ |
| 93 | $23.9 \pm 0.0102$ | $23.8 \pm 0.0076$ | $22.5 \pm 0.0047$ | $22.2 \pm 0.0047$ | $21.6 \pm 0.0035$ | $21.3 \pm 0.0037$ | $21.0 \pm 0.0033$ |


| ID |  | MACS 1206 ACS AB Magnitudes cont. |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{Z}_{850}$ |  |
| 94 | $24.5 \pm 0.0038$ | $24.0 \pm 0.0029$ | $22.9 \pm 0.0021$ | $22.5 \pm 0.0012$ | $21.9 \pm 0.0011$ | $21.6 \pm 0.0007$ | $21.4 \pm 0.0004$ |  |
| 96 | $24.1 \pm 0.0038$ | $23.7 \pm 0.0034$ | $22.3 \pm 0.0028$ | $22.0 \pm 0.0024$ | $21.4 \pm 0.0021$ | $21.1 \pm 0.0016$ | $20.8 \pm 0.0012$ |  |
| 97 | $22.9 \pm 0.0076$ | $22.2 \pm 0.0045$ | $21.0 \pm 0.0028$ | $20.7 \pm 0.0024$ | $20.1 \pm 0.0016$ | $19.8 \pm 0.0013$ | $19.5 \pm 0.001$ |  |
| 98 | $22.1 \pm 0.013$ | $21.6 \pm 0.0091$ | $20.3 \pm 0.0057$ | $20.0 \pm 0.0049$ | $19.3 \pm 0.0026$ | $19.1 \pm 0.0023$ | $18.7 \pm 0.0029$ |  |
| 99 | $23.8 \pm 0.0236$ | $23.4 \pm 0.0083$ | $22.4 \pm 0.0068$ | $22.1 \pm 0.0066$ | $21.6 \pm 0.0031$ | $21.4 \pm 0.0027$ | $21.1 \pm 0.0019$ |  |
| 100 | $23.4 \pm 0.0019$ | $23.0 \pm 0.0014$ | $21.9 \pm 0.0009$ | $21.5 \pm 0.0009$ | $20.9 \pm 0.0006$ | $20.6 \pm 0.0006$ | $20.4 \pm 0.0004$ |  |
| 101 | $23.3 \pm 0.0018$ | $23.0 \pm 0.0015$ | $22.1 \pm 0.0015$ | $21.8 \pm 0.0016$ | $21.3 \pm 0.0011$ | $21.0 \pm 0.0012$ | $20.8 \pm 0.0011$ |  |
| 104 | $24.3 \pm 0.0146$ | $23.9 \pm 0.0098$ | $22.7 \pm 0.0052$ | $22.4 \pm 0.0046$ | $21.9 \pm 0.003$ | $21.6 \pm 0.0024$ | $21.3 \pm 0.0022$ |  |

Table C.2.: Magnitude assessments in the WFC3/IR bands for the fitted early-type galaxies in MACS 1206.

| ID | MACS 1206 WFC3/IR AB Magnitudes |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $Y J_{110}$ | $\mathrm{J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| 1 | $19.6 \pm 0.0004$ | $19.5 \pm 0.0003$ | $19.1 \pm 0.0059$ | $19.2 \pm 0.0003$ | $19.1 \pm 0.0004$ |
| 3 | $19.5 \pm 0.0006$ | $19.3 \pm 0.0004$ | $19.1 \pm 0.0026$ | $19.1 \pm 0.0002$ | $18.9 \pm 0.0002$ |
| 4 | $21.7 \pm 0.0016$ | $21.6 \pm 0.008$ | $21.5 \pm 0.002$ | $21.4 \pm 0.0025$ | $21.2 \pm 0.0024$ |
| 5 | $18.9 \pm 0.0027$ | $18.8 \pm 0.0022$ | $18.7 \pm 0.0024$ | $18.6 \pm 0.0026$ | $18.4 \pm 0.0044$ |
| 6 | $18.3 \pm 0.0003$ | $18.2 \pm 0.0003$ | $18.1 \pm 0.0002$ | $17.9 \pm 0.0002$ | $17.7 \pm 0.0004$ |
| 8 | $21.9 \pm 0.0134$ | $21.8 \pm 0.0787$ | $21.7 \pm 0.0119$ | $21.6 \pm 0.0113$ | $21.5 \pm 0.0162$ |
| 10 | $18.9 \pm 0.0002$ | $18.8 \pm 0.0003$ | $18.6 \pm 0.0002$ | $18.4 \pm 0.0001$ | $18.2 \pm 0.0004$ |
| 11 | $20.0 \pm 0.0015$ | $19.9 \pm 0.0016$ | $19.8 \pm 0.0021$ | $19.7 \pm 0.0031$ | $19.5 \pm 0.0038$ |
| 12 | $19.2 \pm 0.0$ | $19.3 \pm 0.0$ | $19.0 \pm 0.0$ | $18.9 \pm 0.0$ | $18.8 \pm 0.0$ |
| 15 | $20.0 \pm 0.0006$ | $19.9 \pm 0.0006$ | $19.8 \pm 0.0007$ | $19.7 \pm 0.0013$ | $19.6 \pm 0.0025$ |
| 17 | $18.6 \pm 0.0001$ | $18.5 \pm 0.0001$ | $18.4 \pm 0.0002$ | $18.2 \pm 0.0002$ | $18.1 \pm 0.0002$ |
| 19 | $20.9 \pm 0.0081$ | $20.8 \pm 0.0068$ | $20.7 \pm 0.0069$ | $20.6 \pm 0.0076$ | $20.5 \pm 0.0137$ |
| 20 | $21.8 \pm 0.0015$ | $21.8 \pm 0.0022$ | $21.6 \pm 0.0027$ | $21.5 \pm 0.0035$ | $21.4 \pm 0.0057$ |
| 21 | $20.0 \pm 0.0012$ | $19.9 \pm 0.0014$ | $19.8 \pm 0.0021$ | $19.7 \pm 0.0042$ | $19.5 \pm 0.0089$ |
| 22 | $20.1 \pm 0.0006$ | $20.0 \pm 0.0005$ | $19.9 \pm 0.0005$ | $19.8 \pm 0.0011$ | $19.6 \pm 0.0024$ |
| 23 | $20.9 \pm 0.0002$ | $20.9 \pm 0.0004$ | $20.7 \pm 0.0008$ | $20.6 \pm 0.0014$ | $20.5 \pm 0.0018$ |
| 24 | $21.0 \pm 0.0043$ | $20.9 \pm 0.0028$ | $20.9 \pm 0.0029$ | $20.8 \pm 0.0023$ | $20.7 \pm 0.0033$ |
| 25 | $20.3 \pm 0.0004$ | $20.2 \pm 0.0005$ | $20.1 \pm 0.0005$ | $19.9 \pm 0.0007$ | $19.7 \pm 0.0009$ |
| 26 | $20.1 \pm 0.0005$ | $19.9 \pm 0.0004$ | $19.8 \pm 0.0005$ | $19.7 \pm 0.0005$ | $19.5 \pm 0.0007$ |
| 28 | $19.8 \pm 0.0002$ | $19.7 \pm 0.0001$ | $19.6 \pm 0.0002$ | $19.5 \pm 0.0002$ | $19.3 \pm 0.0004$ |
| 29 | $19.4 \pm 0.0002$ | $19.2 \pm 0.0001$ | $19.1 \pm 0.0001$ | $18.9 \pm 0.0002$ | $18.8 \pm 0.0003$ |
| 30 | $20.0 \pm 0.0012$ | $19.9 \pm 0.0011$ | $19.7 \pm 0.0008$ | $19.6 \pm 0.0009$ | $19.3 \pm 0.0014$ |
| 31 | $19.6 \pm 0.0007$ | $19.4 \pm 0.0012$ | $19.3 \pm 0.0021$ | $19.1 \pm 0.0042$ | $19.0 \pm 0.0069$ |
| 32 | $18.9 \pm 0.0003$ | $18.7 \pm 0.0006$ | $18.6 \pm 0.0009$ | $18.4 \pm 0.0011$ | $18.2 \pm 0.0012$ |
| 33 | $21.2 \pm 0.0104$ | $21.1 \pm 0.01$ | $20.7 \pm 0.1308$ | $20.9 \pm 0.0098$ | $20.7 \pm 0.0146$ |
| 35 | $20.6 \pm 0.0012$ | $20.5 \pm 0.0012$ | $20.4 \pm 0.0012$ | $20.3 \pm 0.0017$ | $20.1 \pm 0.0014$ |
| 37 | $20.6 \pm 0.0122$ | $20.5 \pm 0.011$ | $20.5 \pm 0.0106$ | $20.4 \pm 0.0099$ | $20.3 \pm 0.0132$ |
| 38 | $21.3 \pm 0.0007$ | $21.3 \pm 0.0005$ | $21.2 \pm 0.0005$ | $21.1 \pm 0.0006$ | $21.0 \pm 0.0009$ |
| 39 | $20.7 \pm 0.0008$ | $20.6 \pm 0.0012$ | $20.5 \pm 0.0014$ | $20.4 \pm 0.0011$ | $20.2 \pm 0.0011$ |
| 40 | $20.7 \pm 0.0057$ | $20.6 \pm 0.0048$ | $20.5 \pm 0.0048$ | $20.4 \pm 0.0052$ | $20.2 \pm 0.0079$ |
| 41 | $20.6 \pm 0.0004$ | $20.5 \pm 0.0002$ | $20.4 \pm 0.0002$ | $20.2 \pm 0.0004$ | $20.1 \pm 0.0004$ |
| 42 | $19.6 \pm 0.0001$ | $19.5 \pm 0.0$ | $19.4 \pm 0.0001$ | $19.2 \pm 0.0004$ | $19.1 \pm 0.0009$ |
| 43 | $18.2 \pm 0.0001$ | $18.0 \pm 0.0001$ | $17.9 \pm 0.0001$ | $17.7 \pm 0.0001$ | $17.5 \pm 0.0001$ |
| 45 | $20.2 \pm 0.0005$ | $20.2 \pm 0.0004$ | $20.0 \pm 0.0005$ | $19.9 \pm 0.0006$ | $19.7 \pm 0.0009$ |
| 46 | $20.2 \pm 0.0007$ | $20.1 \pm 0.0005$ | $20.0 \pm 0.0004$ | $19.9 \pm 0.0004$ | $19.7 \pm 0.0006$ |
| 47 | $18.6 \pm 0.0004$ | $18.4 \pm 0.0004$ | $18.1 \pm 0.0045$ | $18.2 \pm 0.0004$ | $18.0 \pm 0.0006$ |


| ID | $\mathrm{MACS} 1206 \mathrm{WFC} / \mathrm{IR} \mathrm{AB} \mathrm{Magnitudes} \mathrm{cont}$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |$]$


| ID |  | MACS 1206 WFC3/IR AB Magnitudes cont. |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $\mathrm{YJ}_{110}$ | $\mathrm{~J}_{125}$ | $\mathrm{HH}_{140}$ | $\mathrm{H}_{160}$ |  |
| 100 | $20.2 \pm 0.0003$ | $20.1 \pm 0.0003$ | $20.0 \pm 0.0003$ | $19.9 \pm 0.0003$ | $19.7 \pm 0.0004$ |  |
| 101 | $20.6 \pm 0.0009$ | $20.5 \pm 0.0007$ | $20.4 \pm 0.0006$ | $20.3 \pm 0.0004$ | $20.1 \pm 0.0003$ |  |
| 104 | $21.2 \pm 0.0013$ | $21.1 \pm 0.0011$ | $21.0 \pm 0.0012$ | $20.9 \pm 0.0011$ | $20.7 \pm 0.0017$ |  |

Table C.3.: $r_{e}$ assessments in the ACS bands for the fitted early-type galaxies in MACS 1206. To calculate the values in kpc the cluster redshift $z \sim 0.44$

| ID | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{Z}_{850}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $5.9 \pm 0.04$ | $5.6 \pm 0.036$ | $5.0 \pm 0.026$ | $4.9 \pm 0.023$ | $4.3 \pm 0.014$ | $4.1 \pm 0.011$ | $3.8 \pm 0.007$ |
| 1 | $5.9 \pm 5$ | $1206 \mathrm{ACS} r_{e}[\mathrm{kpc}]$ |  |  |  |  |  |
| 3 | $4.9 \pm 0.021$ | $4.7 \pm 0.019$ | $4.4 \pm 0.014$ | $4.3 \pm 0.013$ | $3.9 \pm 0.009$ | $3.8 \pm 0.007$ | $3.5 \pm 0.005$ |
| 4 | $4.0 \pm 0.087$ | $4.0 \pm 0.072$ | $3.9 \pm 0.042$ | $3.9 \pm 0.034$ | $3.9 \pm 0.024$ | $3.8 \pm 0.026$ | $3.8 \pm 0.033$ |
| 5 | $9.5 \pm 0.139$ | $9.3 \pm 0.123$ | $8.9 \pm 0.089$ | $8.7 \pm 0.078$ | $8.3 \pm 0.047$ | $8.1 \pm 0.038$ | $7.9 \pm 0.024$ |
| 6 | $6.4 \pm 0.034$ | $6.2 \pm 0.03$ | $5.7 \pm 0.023$ | $5.6 \pm 0.021$ | $5.1 \pm 0.013$ | $5.0 \pm 0.011$ | $4.7 \pm 0.008$ |
| 8 | $1.5 \pm 0.067$ | $1.5 \pm 0.059$ | $1.5 \pm 0.042$ | $1.4 \pm 0.037$ | $1.4 \pm 0.023$ | $1.4 \pm 0.019$ | $1.4 \pm 0.014$ |
| 10 | $4.8 \pm 0.011$ | $4.6 \pm 0.01$ | $4.2 \pm 0.008$ | $4.0 \pm 0.007$ | $3.6 \pm 0.005$ | $3.5 \pm 0.004$ | $3.3 \pm 0.003$ |
| 11 | $5.8 \pm 0.055$ | $5.6 \pm 0.049$ | $5.1 \pm 0.036$ | $4.9 \pm 0.031$ | $4.4 \pm 0.019$ | $4.2 \pm 0.015$ | $3.9 \pm 0.01$ |
| 12 | $5.3 \pm 0.0$ | $5.5 \pm 0.0$ | $5.7 \pm 0.0$ | $5.8 \pm 0.0$ | $5.9 \pm 0.0$ | $6.0 \pm 0.0$ | $6.0 \pm 0.0$ |
| 15 | $6.0 \pm 0.049$ | $5.8 \pm 0.044$ | $5.4 \pm 0.032$ | $5.3 \pm 0.028$ | $4.8 \pm 0.017$ | $4.6 \pm 0.014$ | $4.3 \pm 0.009$ |
| 17 | $4.9 \pm 0.011$ | $4.8 \pm 0.01$ | $4.6 \pm 0.007$ | $4.5 \pm 0.007$ | $4.3 \pm 0.004$ | $4.2 \pm 0.004$ | $4.1 \pm 0.003$ |
| 19 | $4.2 \pm 0.034$ | $4.1 \pm 0.029$ | $3.7 \pm 0.019$ | $3.5 \pm 0.016$ | $3.1 \pm 0.007$ | $3.0 \pm 0.006$ | $2.8 \pm 0.006$ |
| 20 | $3.8 \pm 0.159$ | $3.7 \pm 0.136$ | $3.5 \pm 0.088$ | $3.4 \pm 0.074$ | $3.2 \pm 0.041$ | $3.2 \pm 0.036$ | $3.1 \pm 0.034$ |
| 21 | $4.1 \pm 0.11$ | $4.0 \pm 0.098$ | $3.9 \pm 0.071$ | $3.8 \pm 0.063$ | $3.6 \pm 0.039$ | $3.5 \pm 0.032$ | $3.4 \pm 0.024$ |
| 22 | $2.4 \pm 0.028$ | $2.4 \pm 0.025$ | $2.2 \pm 0.018$ | $2.1 \pm 0.016$ | $1.9 \pm 0.01$ | $1.9 \pm 0.008$ | $1.8 \pm 0.005$ |
| 23 | $2.8 \pm 0.035$ | $2.7 \pm 0.032$ | $2.5 \pm 0.024$ | $2.4 \pm 0.021$ | $2.2 \pm 0.014$ | $2.1 \pm 0.012$ | $2.0 \pm 0.008$ |
| 24 | $1.5 \pm 0.052$ | $1.5 \pm 0.046$ | $1.4 \pm 0.034$ | $1.3 \pm 0.03$ | $1.2 \pm 0.02$ | $1.1 \pm 0.018$ | $1.0 \pm 0.015$ |
| 25 | $1.9 \pm 0.017$ | $1.8 \pm 0.015$ | $1.6 \pm 0.01$ | $1.6 \pm 0.009$ | $1.4 \pm 0.005$ | $1.3 \pm 0.004$ | $1.2 \pm 0.004$ |
| 26 | $1.6 \pm 0.006$ | $1.6 \pm 0.005$ | $1.5 \pm 0.003$ | $1.5 \pm 0.003$ | $1.4 \pm 0.002$ | $1.3 \pm 0.002$ | $1.3 \pm 0.002$ |
| 28 | $4.5 \pm 0.017$ | $4.2 \pm 0.015$ | $3.8 \pm 0.01$ | $3.6 \pm 0.009$ | $3.1 \pm 0.004$ | $2.9 \pm 0.003$ | $2.7 \pm 0.002$ |
| 29 | $3.6 \pm 0.007$ | $3.4 \pm 0.006$ | $3.2 \pm 0.004$ | $3.1 \pm 0.004$ | $2.8 \pm 0.002$ | $2.7 \pm 0.002$ | $2.6 \pm 0.001$ |


| ID | MACS 1206 ACS $r_{e}[\mathrm{kpc}]$ cont. |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{B}_{435}$ | $\mathrm{g}_{475}$ | $\mathrm{V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | Z850 |
| 30 | $3.5 \pm 0.05$ | $3.3 \pm 0.045$ | $3.0 \pm 0.033$ | $2.9 \pm 0.029$ | $2.5 \pm 0.018$ | $2.4 \pm 0.015$ | $2.3 \pm 0.01$ |
| 31 | $4.1 \pm 0.075$ | $4.1 \pm 0.064$ | $4.0 \pm 0.041$ | $4.0 \pm 0.035$ | $4.1 \pm 0.021$ | $4.1 \pm 0.02$ | $4.1 \pm 0.02$ |
| 32 | $4.2 \pm 0.03$ | $4.1 \pm 0.027$ | $3.7 \pm 0.021$ | $3.6 \pm 0.02$ | $3.3 \pm 0.014$ | $3.2 \pm 0.012$ | $3.0 \pm 0.009$ |
| 33 | $2.4 \pm 0.147$ | $2.3 \pm 0.131$ | $2.1 \pm 0.095$ | $2.0 \pm 0.083$ | $1.8 \pm 0.05$ | $1.7 \pm 0.04$ | $1.5 \pm 0.026$ |
| 35 | $1.1 \pm 0.008$ | $1.1 \pm 0.007$ | $1.0 \pm 0.006$ | $1.0 \pm 0.005$ | $0.9 \pm 0.003$ | $0.9 \pm 0.003$ | $0.8 \pm 0.002$ |
| 37 | $4.4 \pm 0.238$ | $4.4 \pm 0.215$ | $4.4 \pm 0.163$ | $4.4 \pm 0.146$ | $4.2 \pm 0.097$ | $4.2 \pm 0.082$ | $4.0 \pm 0.058$ |
| 38 | $2.5 \pm 0.02$ | $2.4 \pm 0.018$ | $2.2 \pm 0.013$ | $2.1 \pm 0.012$ | $1.8 \pm 0.007$ | $1.7 \pm 0.006$ | $1.5 \pm 0.004$ |
| 39 | $2.3 \pm 0.034$ | $2.2 \pm 0.03$ | $2.0 \pm 0.023$ | $1.9 \pm 0.021$ | $1.6 \pm 0.014$ | $1.5 \pm 0.012$ | $1.4 \pm 0.008$ |
| 40 | $2.1 \pm 0.06$ | $2.0 \pm 0.054$ | $1.8 \pm 0.04$ | $1.8 \pm 0.035$ | $1.6 \pm 0.022$ | $1.5 \pm 0.018$ | $1.4 \pm 0.012$ |
| 41 | $1.5 \pm 0.003$ | $1.5 \pm 0.003$ | $1.4 \pm 0.002$ | $1.4 \pm 0.002$ | $1.3 \pm 0.001$ | $1.2 \pm 0.001$ | $1.2 \pm 0.001$ |
| 42 | $3.7 \pm 0.009$ | $3.6 \pm 0.008$ | $3.4 \pm 0.005$ | $3.3 \pm 0.005$ | $3.1 \pm 0.002$ | $3.0 \pm 0.002$ | $2.9 \pm 0.002$ |
| 43 | $7.4 \pm 0.005$ | $7.2 \pm 0.005$ | $6.7 \pm 0.004$ | $6.5 \pm 0.003$ | $6.0 \pm 0.002$ | $5.9 \pm 0.002$ | $5.7 \pm 0.001$ |
| 45 | $2.9 \pm 0.031$ | $2.8 \pm 0.027$ | $2.6 \pm 0.019$ | $2.5 \pm 0.016$ | $2.3 \pm 0.01$ | $2.3 \pm 0.008$ | $2.2 \pm 0.006$ |
| 46 | $2.5 \pm 0.009$ | $2.4 \pm 0.008$ | $2.2 \pm 0.006$ | $2.1 \pm 0.006$ | $1.9 \pm 0.004$ | $1.8 \pm 0.003$ | $1.7 \pm 0.002$ |
| 47 | $9.2 \pm 0.092$ | $8.9 \pm 0.083$ | $8.2 \pm 0.063$ | $8.0 \pm 0.057$ | $7.4 \pm 0.04$ | $7.2 \pm 0.035$ | $6.8 \pm 0.028$ |
| 48 | $1.6 \pm 0.009$ | $1.5 \pm 0.008$ | $1.3 \pm 0.006$ | $1.3 \pm 0.006$ | $1.1 \pm 0.004$ | $1.0 \pm 0.003$ | $0.9 \pm 0.002$ |
| 50 | $2.2 \pm 0.004$ | $2.2 \pm 0.004$ | $2.1 \pm 0.002$ | $2.1 \pm 0.002$ | $2.0 \pm 0.001$ | $1.9 \pm 0.001$ | $1.9 \pm 0.001$ |
| 51 | $4.3 \pm 0.003$ | $4.2 \pm 0.002$ | $3.9 \pm 0.002$ | $3.8 \pm 0.002$ | $3.5 \pm 0.001$ | $3.4 \pm 0.001$ | $3.2 \pm 0.0$ |
| 52 | $8.4 \pm 0.008$ | $8.1 \pm 0.007$ | $7.5 \pm 0.004$ | $7.4 \pm 0.003$ | $6.8 \pm 0.002$ | $6.7 \pm 0.002$ | $6.5 \pm 0.003$ |
| 55 | $2.9 \pm 0.027$ | $2.8 \pm 0.023$ | $2.6 \pm 0.016$ | $2.5 \pm 0.013$ | $2.3 \pm 0.008$ | $2.2 \pm 0.006$ | $2.1 \pm 0.006$ |
| 56 | $1.4 \pm 0.006$ | $1.3 \pm 0.005$ | $1.2 \pm 0.003$ | $1.2 \pm 0.003$ | $1.0 \pm 0.001$ | $1.0 \pm 0.001$ | $0.9 \pm 0.001$ |
| 57 | $2.1 \pm 0.016$ | $2.1 \pm 0.014$ | $2.1 \pm 0.009$ | $2.0 \pm 0.008$ | $2.0 \pm 0.005$ | $2.0 \pm 0.004$ | $1.9 \pm 0.004$ |
| 58 | $1.7 \pm 0.072$ | $1.7 \pm 0.065$ | $1.5 \pm 0.05$ | $1.5 \pm 0.046$ | $1.3 \pm 0.032$ | $1.3 \pm 0.028$ | $1.2 \pm 0.022$ |
| 59 | $1.4 \pm 0.02$ | $1.4 \pm 0.018$ | $1.2 \pm 0.013$ | $1.2 \pm 0.012$ | $1.1 \pm 0.008$ | $1.0 \pm 0.006$ | $0.9 \pm 0.004$ |


| ID | MACS 1206 ACS $r_{e}[\mathrm{kpc}]$ cont. |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{B}_{435}$ | $\mathrm{g}_{475}$ | $\mathrm{V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | Z850 |
| 61 | $2.0 \pm 0.029$ | $1.9 \pm 0.026$ | $1.8 \pm 0.018$ | $1.8 \pm 0.016$ | $1.7 \pm 0.01$ | $1.7 \pm 0.008$ | $1.6 \pm 0.006$ |
| 62 | $2.1 \pm 0.025$ | $2.0 \pm 0.023$ | $1.8 \pm 0.017$ | $1.8 \pm 0.015$ | $1.6 \pm 0.01$ | $1.5 \pm 0.008$ | $1.4 \pm 0.005$ |
| 63 | $1.3 \pm 0.079$ | $1.4 \pm 0.07$ | $1.5 \pm 0.05$ | $1.5 \pm 0.044$ | $1.6 \pm 0.028$ | $1.6 \pm 0.024$ | $1.6 \pm 0.019$ |
| 64 | $2.0 \pm 0.043$ | $1.9 \pm 0.038$ | $1.7 \pm 0.028$ | $1.7 \pm 0.025$ | $1.5 \pm 0.016$ | $1.5 \pm 0.013$ | $1.4 \pm 0.008$ |
| 65 | $4.9 \pm 0.18$ | $4.7 \pm 0.162$ | $4.1 \pm 0.124$ | $3.9 \pm 0.112$ | $3.4 \pm 0.075$ | $3.2 \pm 0.063$ | $2.8 \pm 0.044$ |
| 68 | $3.4 \pm 0.058$ | $3.2 \pm 0.051$ | $2.9 \pm 0.036$ | $2.8 \pm 0.032$ | $2.6 \pm 0.02$ | $2.5 \pm 0.017$ | $2.3 \pm 0.014$ |
| 69 | $4.9 \pm 0.05$ | $4.7 \pm 0.042$ | $4.1 \pm 0.025$ | $3.9 \pm 0.02$ | $3.4 \pm 0.011$ | $3.2 \pm 0.01$ | $2.9 \pm 0.013$ |
| 70 | $3.2 \pm 0.064$ | $3.1 \pm 0.055$ | $2.8 \pm 0.035$ | $2.8 \pm 0.03$ | $2.6 \pm 0.016$ | $2.6 \pm 0.014$ | $2.5 \pm 0.014$ |
| 71 | $3.0 \pm 0.03$ | $2.9 \pm 0.026$ | $2.6 \pm 0.019$ | $2.6 \pm 0.017$ | $2.3 \pm 0.01$ | $2.3 \pm 0.008$ | $2.1 \pm 0.005$ |
| 72 | $2.0 \pm 0.05$ | $2.0 \pm 0.042$ | $1.8 \pm 0.026$ | $1.7 \pm 0.021$ | $1.5 \pm 0.01$ | $1.5 \pm 0.009$ | $1.4 \pm 0.011$ |
| 73 | $1.8 \pm 0.017$ | $1.8 \pm 0.015$ | $1.7 \pm 0.011$ | $1.7 \pm 0.01$ | $1.7 \pm 0.006$ | $1.7 \pm 0.005$ | $1.7 \pm 0.003$ |
| 74 | $0.9 \pm 0.024$ | $0.9 \pm 0.021$ | $0.7 \pm 0.015$ | $0.7 \pm 0.013$ | $0.6 \pm 0.007$ | $0.5 \pm 0.006$ | $0.5 \pm 0.004$ |
| 76 | $4.7 \pm 0.04$ | $4.5 \pm 0.036$ | $4.0 \pm 0.027$ | $3.9 \pm 0.024$ | $3.4 \pm 0.016$ | $3.2 \pm 0.013$ | $3.0 \pm 0.009$ |
| 77 | $10.5 \pm 0.695$ | $10.5 \pm 0.606$ | $10.6 \pm 0.415$ | $10.6 \pm 0.357$ | $10.7 \pm 0.201$ | $10.7 \pm 0.163$ | $10.7 \pm 0.133$ |
| 80 | $1.2 \pm 0.012$ | $1.2 \pm 0.011$ | $1.1 \pm 0.008$ | $1.0 \pm 0.007$ | $1.0 \pm 0.004$ | $0.9 \pm 0.003$ | $0.9 \pm 0.002$ |
| 81 | $2.6 \pm 0.042$ | $2.5 \pm 0.037$ | $2.3 \pm 0.027$ | $2.3 \pm 0.024$ | $2.1 \pm 0.016$ | $2.0 \pm 0.013$ | $2.0 \pm 0.009$ |
| 82 | $0.9 \pm 0.012$ | $0.8 \pm 0.011$ | $0.7 \pm 0.008$ | $0.7 \pm 0.007$ | $0.6 \pm 0.005$ | $0.5 \pm 0.004$ | $0.5 \pm 0.002$ |
| 83 | $2.4 \pm 0.057$ | $2.3 \pm 0.05$ | $2.1 \pm 0.036$ | $2.0 \pm 0.032$ | $1.8 \pm 0.019$ | $1.7 \pm 0.015$ | $1.6 \pm 0.01$ |
| 84 | $1.1 \pm 0.003$ | $1.1 \pm 0.003$ | $1.0 \pm 0.001$ | $1.0 \pm 0.001$ | $1.0 \pm 0.001$ | $0.9 \pm 0.001$ | $0.9 \pm 0.001$ |
| 85 | $2.5 \pm 0.109$ | $2.4 \pm 0.1$ | $2.2 \pm 0.078$ | $2.2 \pm 0.071$ | $2.0 \pm 0.05$ | $1.9 \pm 0.043$ | $1.8 \pm 0.032$ |
| 87 | $3.1 \pm 0.029$ | $3.0 \pm 0.026$ | $2.9 \pm 0.019$ | $2.8 \pm 0.016$ | $2.6 \pm 0.011$ | $2.5 \pm 0.009$ | $2.4 \pm 0.008$ |
| 88 | $1.9 \pm 0.019$ | $1.9 \pm 0.017$ | $1.8 \pm 0.012$ | $1.8 \pm 0.011$ | $1.7 \pm 0.006$ | $1.6 \pm 0.005$ | $1.6 \pm 0.003$ |
| 89 | $2.4 \pm 0.01$ | $2.3 \pm 0.009$ | $2.1 \pm 0.007$ | $2.1 \pm 0.006$ | $1.9 \pm 0.004$ | $1.8 \pm 0.003$ | $1.7 \pm 0.002$ |
| 90 | $1.2 \pm 0.031$ | $1.2 \pm 0.027$ | $1.1 \pm 0.02$ | $1.1 \pm 0.018$ | $1.0 \pm 0.011$ | $0.9 \pm 0.009$ | $0.9 \pm 0.006$ |


| ID |  | MACS 1206 ACS $r_{e}[\mathrm{kpc}]$ cont. |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{Z}_{850}$ |  |
| 91 | $2.1 \pm 0.003$ | $2.1 \pm 0.003$ | $1.9 \pm 0.002$ | $1.9 \pm 0.002$ | $1.7 \pm 0.001$ | $1.7 \pm 0.001$ | $1.6 \pm 0.001$ |  |
| 93 | $1.6 \pm 0.064$ | $1.6 \pm 0.056$ | $1.5 \pm 0.041$ | $1.5 \pm 0.036$ | $1.4 \pm 0.022$ | $1.3 \pm 0.018$ | $1.3 \pm 0.013$ |  |
| 94 | $0.6 \pm 0.003$ | $0.6 \pm 0.002$ | $0.5 \pm 0.002$ | $0.5 \pm 0.001$ | $0.5 \pm 0.001$ | $0.5 \pm 0.001$ | $0.4 \pm 0.0$ |  |
| 96 | $1.3 \pm 0.018$ | $1.3 \pm 0.016$ | $1.2 \pm 0.01$ | $1.2 \pm 0.009$ | $1.1 \pm 0.005$ | $1.1 \pm 0.004$ | $1.0 \pm 0.004$ |  |
| 97 | $2.5 \pm 0.03$ | $2.4 \pm 0.027$ | $2.2 \pm 0.02$ | $2.1 \pm 0.017$ | $1.9 \pm 0.01$ | $1.8 \pm 0.008$ | $1.7 \pm 0.005$ |  |
| 98 | $6.8 \pm 0.126$ | $6.7 \pm 0.108$ | $6.5 \pm 0.068$ | $6.5 \pm 0.057$ | $6.3 \pm 0.031$ | $6.2 \pm 0.028$ | $6.1 \pm 0.033$ |  |
| 99 | $3.1 \pm 0.051$ | $3.0 \pm 0.045$ | $2.6 \pm 0.032$ | $2.5 \pm 0.028$ | $2.2 \pm 0.018$ | $2.1 \pm 0.016$ | $2.0 \pm 0.014$ |  |
| 100 | $3.9 \pm 0.015$ | $3.8 \pm 0.013$ | $3.5 \pm 0.009$ | $3.4 \pm 0.008$ | $3.0 \pm 0.004$ | $2.9 \pm 0.003$ | $2.7 \pm 0.002$ |  |
| 101 | $1.8 \pm 0.006$ | $1.7 \pm 0.005$ | $1.5 \pm 0.004$ | $1.5 \pm 0.004$ | $1.3 \pm 0.002$ | $1.2 \pm 0.002$ | $1.1 \pm 0.002$ |  |
| 104 | $1.7 \pm 0.02$ | $1.7 \pm 0.018$ | $1.6 \pm 0.013$ | $1.5 \pm 0.011$ | $1.4 \pm 0.007$ | $1.4 \pm 0.005$ | $1.3 \pm 0.003$ |  |

## C. MegaMorph results

Table C.4.: $r_{e}$ assessments in the WFC3/IR bands for the fitted early-type galaxies in MACS 1206. To calculate the values in kpc the cosmological parameters $H_{0}=70 \mathrm{~km} \mathrm{~s}^{-1}$ $\mathrm{Mpc}^{-1}, \Omega_{\Lambda}=0.7$ and $\Omega_{M}=0.3$ have been adopted with a pixelscale of $0.065 \% / \mathrm{px}$ and a cluster redshift $z \sim 0.44$.

| ID | MACS $1206 \mathrm{WFC} 3 / \mathrm{IR} r_{e}[\mathrm{kpc}]$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $\mathrm{YJ}_{110}$ | $\mathrm{~J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| 1 | $3.6 \pm 0.005$ | $3.5 \pm 0.005$ | $3.5 \pm 0.005$ | $3.5 \pm 0.005$ | $3.6 \pm 0.007$ |
| 3 | $3.4 \pm 0.004$ | $3.2 \pm 0.003$ | $3.1 \pm 0.003$ | $3.0 \pm 0.003$ | $2.9 \pm 0.003$ |
| 4 | $3.8 \pm 0.036$ | $3.9 \pm 0.035$ | $3.9 \pm 0.03$ | $4.0 \pm 0.021$ | $4.1 \pm 0.032$ |
| 5 | $7.7 \pm 0.018$ | $7.6 \pm 0.017$ | $7.6 \pm 0.018$ | $7.6 \pm 0.021$ | $7.7 \pm 0.033$ |
| 6 | $4.6 \pm 0.006$ | $4.5 \pm 0.006$ | $4.5 \pm 0.007$ | $4.5 \pm 0.009$ | $4.7 \pm 0.013$ |
| 8 | $1.3 \pm 0.013$ | $1.3 \pm 0.013$ | $1.3 \pm 0.012$ | $1.3 \pm 0.011$ | $1.2 \pm 0.015$ |
| 10 | $3.1 \pm 0.002$ | $3.0 \pm 0.002$ | $2.9 \pm 0.001$ | $2.9 \pm 0.002$ | $3.0 \pm 0.002$ |
| 11 | $3.7 \pm 0.007$ | $3.6 \pm 0.007$ | $3.5 \pm 0.007$ | $3.5 \pm 0.007$ | $3.6 \pm 0.01$ |
| 12 | $6.0 \pm 0.0$ | $5.9 \pm 0.0$ | $5.8 \pm 0.0$ | $5.5 \pm 0.0$ | $5.1 \pm 0.0$ |
| 15 | $4.0 \pm 0.006$ | $3.8 \pm 0.006$ | $3.5 \pm 0.006$ | $3.2 \pm 0.005$ | $3.0 \pm 0.005$ |
| 17 | $4.0 \pm 0.002$ | $4.0 \pm 0.002$ | $4.0 \pm 0.002$ | $4.0 \pm 0.002$ | $4.1 \pm 0.003$ |
| 19 | $2.6 \pm 0.006$ | $2.5 \pm 0.006$ | $2.4 \pm 0.004$ | $2.3 \pm 0.003$ | $2.3 \pm 0.011$ |
| 20 | $3.1 \pm 0.034$ | $3.2 \pm 0.032$ | $3.3 \pm 0.032$ | $3.6 \pm 0.05$ | $4.0 \pm 0.096$ |
| 21 | $3.3 \pm 0.021$ | $3.2 \pm 0.022$ | $3.2 \pm 0.023$ | $3.1 \pm 0.028$ | $3.0 \pm 0.039$ |
| 22 | $1.7 \pm 0.004$ | $1.6 \pm 0.004$ | $1.6 \pm 0.004$ | $1.6 \pm 0.004$ | $1.6 \pm 0.007$ |
| 23 | $1.9 \pm 0.006$ | $1.8 \pm 0.005$ | $1.8 \pm 0.004$ | $1.7 \pm 0.004$ | $1.7 \pm 0.006$ |
| 24 | $0.9 \pm 0.014$ | $0.9 \pm 0.013$ | $0.8 \pm 0.012$ | $0.7 \pm 0.011$ | $0.7 \pm 0.012$ |
| 25 | $1.1 \pm 0.004$ | $1.1 \pm 0.003$ | $1.0 \pm 0.003$ | $1.0 \pm 0.003$ | $1.0 \pm 0.006$ |
| 26 | $1.2 \pm 0.002$ | $1.2 \pm 0.002$ | $1.2 \pm 0.002$ | $1.1 \pm 0.002$ | $1.1 \pm 0.003$ |
| 28 | $2.5 \pm 0.002$ | $2.4 \pm 0.002$ | $2.3 \pm 0.002$ | $2.3 \pm 0.001$ | $2.4 \pm 0.004$ |
| 29 | $2.5 \pm 0.001$ | $2.4 \pm 0.001$ | $2.4 \pm 0.001$ | $2.4 \pm 0.001$ | $2.6 \pm 0.002$ |
| 30 | $2.2 \pm 0.008$ | $2.2 \pm 0.007$ | $2.2 \pm 0.008$ | $2.3 \pm 0.011$ | $2.4 \pm 0.017$ |
| 31 | $4.1 \pm 0.019$ | $4.2 \pm 0.017$ | $4.3 \pm 0.016$ | $4.4 \pm 0.024$ | $4.5 \pm 0.048$ |
| 32 | $2.9 \pm 0.007$ | $2.9 \pm 0.006$ | $2.9 \pm 0.006$ | $2.9 \pm 0.008$ | $3.1 \pm 0.01$ |
| 33 | $1.4 \pm 0.02$ | $1.4 \pm 0.019$ | $1.3 \pm 0.02$ | $1.3 \pm 0.02$ | $1.3 \pm 0.028$ |
| 35 | $0.8 \pm 0.001$ | $0.7 \pm 0.001$ | $0.7 \pm 0.001$ | $0.7 \pm 0.001$ | $0.7 \pm 0.001$ |
| 37 | $3.8 \pm 0.043$ | $3.6 \pm 0.035$ | $3.4 \pm 0.031$ | $3.1 \pm 0.027$ | $2.6 \pm 0.029$ |
| 38 | $1.4 \pm 0.002$ | $1.3 \pm 0.002$ | $1.2 \pm 0.002$ | $1.1 \pm 0.002$ | $1.1 \pm 0.002$ |
| 39 | $1.3 \pm 0.005$ | $1.2 \pm 0.004$ | $1.1 \pm 0.003$ | $1.1 \pm 0.003$ | $1.0 \pm 0.004$ |
| 40 | $1.3 \pm 0.008$ | $1.3 \pm 0.007$ | $1.3 \pm 0.007$ | $1.2 \pm 0.008$ | $1.3 \pm 0.012$ |
| 41 | $1.1 \pm 0.001$ | $1.1 \pm 0.001$ | $1.0 \pm 0.001$ | $1.0 \pm 0.0$ | $1.0 \pm 0.001$ |
| 42 | $2.7 \pm 0.002$ | $2.6 \pm 0.002$ | $2.5 \pm 0.002$ | $2.4 \pm 0.002$ | $2.3 \pm 0.002$ |
| 43 | $5.5 \pm 0.001$ | $5.4 \pm 0.001$ | $5.4 \pm 0.001$ | $5.5 \pm 0.001$ | $5.8 \pm 0.001$ |
| 45 | $2.1 \pm 0.005$ | $2.0 \pm 0.005$ | $2.0 \pm 0.004$ | $2.0 \pm 0.004$ | $2.0 \pm 0.009$ |
| 46 | $1.6 \pm 0.002$ | $1.5 \pm 0.001$ | $1.5 \pm 0.001$ | $1.5 \pm 0.001$ | $1.5 \pm 0.002$ |
| 47 | $6.5 \pm 0.024$ | $6.3 \pm 0.022$ | $6.2 \pm 0.021$ | $6.2 \pm 0.023$ | $6.3 \pm 0.032$ |
|  |  |  |  |  |  |


| ID |  | MACS 1206 WFC3/IR $r_{e}[\mathrm{kpc}]$ cont. |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $\mathrm{YJ}_{110}$ | $\mathrm{~J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| 48 | $0.8 \pm 0.002$ | $0.8 \pm 0.001$ | $0.8 \pm 0.001$ | $0.7 \pm 0.001$ | $0.7 \pm 0.001$ |
| 50 | $1.9 \pm 0.001$ | $1.9 \pm 0.001$ | $1.8 \pm 0.001$ | $1.9 \pm 0.001$ | $1.9 \pm 0.003$ |
| 51 | $3.1 \pm 0.0$ | $3.0 \pm 0.001$ | $3.0 \pm 0.001$ | $3.0 \pm 0.001$ | $3.1 \pm 0.001$ |
| 52 | $6.3 \pm 0.003$ | $6.3 \pm 0.002$ | $6.3 \pm 0.002$ | $6.5 \pm 0.003$ | $6.9 \pm 0.007$ |
| 55 | $2.1 \pm 0.006$ | $2.1 \pm 0.007$ | $2.1 \pm 0.007$ | $2.2 \pm 0.009$ | $2.3 \pm 0.013$ |
| 56 | $0.9 \pm 0.001$ | $0.8 \pm 0.001$ | $0.8 \pm 0.001$ | $0.8 \pm 0.001$ | $0.8 \pm 0.001$ |
| 57 | $1.9 \pm 0.004$ | $1.9 \pm 0.004$ | $1.8 \pm 0.004$ | $1.8 \pm 0.004$ | $1.8 \pm 0.006$ |
| 58 | $1.1 \pm 0.019$ | $1.0 \pm 0.017$ | $0.9 \pm 0.015$ | $0.8 \pm 0.015$ | $0.8 \pm 0.018$ |
| 59 | $0.9 \pm 0.002$ | $0.9 \pm 0.002$ | $0.8 \pm 0.002$ | $0.8 \pm 0.002$ | $0.9 \pm 0.003$ |
| 61 | $1.6 \pm 0.005$ | $1.5 \pm 0.005$ | $1.5 \pm 0.005$ | $1.5 \pm 0.005$ | $1.5 \pm 0.007$ |
| 62 | $1.3 \pm 0.004$ | $1.3 \pm 0.003$ | $1.2 \pm 0.003$ | $1.2 \pm 0.004$ | $1.2 \pm 0.005$ |
| 63 | $1.6 \pm 0.018$ | $1.6 \pm 0.018$ | $1.5 \pm 0.017$ | $1.4 \pm 0.015$ | $1.2 \pm 0.016$ |
| 64 | $1.3 \pm 0.005$ | $1.3 \pm 0.004$ | $1.3 \pm 0.004$ | $1.2 \pm 0.006$ | $1.3 \pm 0.01$ |
| 65 | $2.5 \pm 0.031$ | $2.3 \pm 0.024$ | $2.2 \pm 0.02$ | $2.0 \pm 0.018$ | $1.9 \pm 0.026$ |
| 68 | $2.2 \pm 0.013$ | $2.2 \pm 0.012$ | $2.2 \pm 0.012$ | $2.2 \pm 0.017$ | $2.3 \pm 0.029$ |
| 69 | $2.8 \pm 0.014$ | $2.7 \pm 0.012$ | $2.7 \pm 0.009$ | $2.7 \pm 0.007$ | $3.0 \pm 0.021$ |
| 70 | $2.5 \pm 0.014$ | $2.6 \pm 0.013$ | $2.7 \pm 0.011$ | $2.9 \pm 0.012$ | $3.2 \pm 0.025$ |
| 71 | $2.0 \pm 0.005$ | $1.9 \pm 0.005$ | $1.9 \pm 0.005$ | $1.9 \pm 0.006$ | $1.9 \pm 0.008$ |
| 72 | $1.3 \pm 0.011$ | $1.2 \pm 0.01$ | $1.2 \pm 0.006$ | $1.1 \pm 0.007$ | $1.1 \pm 0.023$ |
| 73 | $1.7 \pm 0.003$ | $1.7 \pm 0.002$ | $1.7 \pm 0.002$ | $1.7 \pm 0.002$ | $1.8 \pm 0.003$ |
| 74 | $0.4 \pm 0.003$ | $0.4 \pm 0.003$ | $0.4 \pm 0.002$ | $0.4 \pm 0.003$ | $0.4 \pm 0.006$ |
| 76 | $2.8 \pm 0.006$ | $2.7 \pm 0.005$ | $2.6 \pm 0.005$ | $2.6 \pm 0.005$ | $2.7 \pm 0.007$ |
| 77 | $10.7 \pm 0.134$ | $10.7 \pm 0.135$ | $10.7 \pm 0.125$ | $10.6 \pm 0.107$ | $10.5 \pm 0.17$ |
| 80 | $0.9 \pm 0.002$ | $0.8 \pm 0.001$ | $0.8 \pm 0.001$ | $0.9 \pm 0.002$ | $0.9 \pm 0.003$ |
| 81 | $1.9 \pm 0.007$ | $1.9 \pm 0.006$ | $1.9 \pm 0.006$ | $1.9 \pm 0.009$ | $2.0 \pm 0.013$ |
| 82 | $0.4 \pm 0.002$ | $0.4 \pm 0.001$ | $0.4 \pm 0.001$ | $0.4 \pm 0.001$ | $0.4 \pm 0.002$ |
| 83 | $1.6 \pm 0.008$ | $1.6 \pm 0.007$ | $1.6 \pm 0.007$ | $1.6 \pm 0.008$ | $1.8 \pm 0.013$ |
| 84 | $0.9 \pm 0.001$ | $0.9 \pm 0.001$ | $0.9 \pm 0.001$ | $0.9 \pm 0.001$ | $0.9 \pm 0.002$ |
| 85 | $1.7 \pm 0.024$ | $1.7 \pm 0.02$ | $1.6 \pm 0.017$ | $1.6 \pm 0.017$ | $1.5 \pm 0.022$ |
| 87 | $2.3 \pm 0.008$ | $2.2 \pm 0.008$ | $2.0 \pm 0.008$ | $1.9 \pm 0.01$ | $1.7 \pm 0.014$ |
| 88 | $1.5 \pm 0.003$ | $1.5 \pm 0.003$ | $1.4 \pm 0.003$ | $1.4 \pm 0.003$ | $1.4 \pm 0.004$ |
| 89 | $1.6 \pm 0.001$ | $1.5 \pm 0.001$ | $1.5 \pm 0.001$ | $1.5 \pm 0.001$ | $1.5 \pm 0.001$ |
| 90 | $0.9 \pm 0.004$ | $0.8 \pm 0.004$ | $0.8 \pm 0.003$ | $0.8 \pm 0.003$ | $0.8 \pm 0.004$ |
| 91 | $1.5 \pm 0.0$ | $1.4 \pm 0.0$ | $1.4 \pm 0.0$ | $1.3 \pm 0.0$ | $1.3 \pm 0.001$ |
| 93 | $1.2 \pm 0.011$ | $1.1 \pm 0.01$ | $1.1 \pm 0.01$ | $1.0 \pm 0.008$ | $0.9 \pm 0.009$ |
| 94 | $0.4 \pm 0.0$ | $0.4 \pm 0.0$ | $0.4 \pm 0.0$ | $0.4 \pm 0.0$ | $0.4 \pm 0.0$ |
| 96 | $1.0 \pm 0.005$ | $1.0 \pm 0.005$ | $0.9 \pm 0.005$ | $0.9 \pm 0.005$ | $0.9 \pm 0.008$ |
| 97 | $1.6 \pm 0.004$ | $1.5 \pm 0.003$ | $1.5 \pm 0.003$ | $1.5 \pm 0.004$ | $1.5 \pm 0.007$ |
| 98 | $6.0 \pm 0.04$ | $5.9 \pm 0.042$ | $5.8 \pm 0.042$ | $5.6 \pm 0.043$ | $5.5 \pm 0.056$ |
| 99 | $1.9 \pm 0.013$ | $1.8 \pm 0.013$ | $1.8 \pm 0.012$ | $1.9 \pm 0.014$ | $2.0 \pm 0.023$ |
|  |  |  |  |  |  |


| ID | MACS 1206 WFC3/IR $r_{e}[\mathrm{kpc}]$ cont. |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $\mathrm{YJ}_{110}$ | $\mathrm{~J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| 100 | $2.6 \pm 0.003$ | $2.5 \pm 0.003$ | $2.4 \pm 0.003$ | $2.4 \pm 0.002$ | $2.4 \pm 0.003$ |
| 101 | $1.1 \pm 0.001$ | $1.0 \pm 0.001$ | $1.0 \pm 0.001$ | $0.9 \pm 0.001$ | $1.0 \pm 0.002$ |
| 104 | $1.3 \pm 0.003$ | $1.2 \pm 0.002$ | $1.2 \pm 0.002$ | $1.1 \pm 0.002$ | $1.1 \pm 0.003$ |


| Table C.5.: Measurements of $n$ in the ACS bands for the selected early-type galaxies in MACS 1206 |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ID |  | MACS 1206 ACS Sérsic $n$ |  |  |  |  |  |  | $\mathrm{I}_{8}$ | $\mathrm{I}_{814}$ | $\mathrm{z}_{850}$ |
|  | $\mathrm{~B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $3.2 \pm 0.004$ |  |  |  |  |  |
| 1 | $4.9 \pm 0.019$ | $4.7 \pm 0.017$ | $4.2 \pm 0.012$ | $4.1 \pm 0.011$ | $3.6 \pm 0.007$ | $3.5 \pm 0.005$ | $3.2 \pm 0.3$ |  |  |  |  |
| 3 | $4.8 \pm 0.013$ | $4.8 \pm 0.012$ | $4.6 \pm 0.009$ | $4.5 \pm 0.008$ | $4.3 \pm 0.005$ | $4.3 \pm 0.004$ | $4.2 \pm 0.003$ |  |  |  |  |
| 4 | $3.4 \pm 0.057$ | $3.4 \pm 0.047$ | $3.5 \pm 0.027$ | $3.5 \pm 0.022$ | $3.6 \pm 0.015$ | $3.6 \pm 0.017$ | $3.7 \pm 0.022$ |  |  |  |  |
| 5 | $3.0 \pm 0.031$ | $2.9 \pm 0.028$ | $2.8 \pm 0.02$ | $2.8 \pm 0.018$ | $2.7 \pm 0.011$ | $2.7 \pm 0.009$ | $2.6 \pm 0.006$ |  |  |  |  |
| 6 | $4.7 \pm 0.012$ | $4.6 \pm 0.011$ | $4.5 \pm 0.008$ | $4.4 \pm 0.008$ | $4.3 \pm 0.005$ | $4.2 \pm 0.004$ | $4.2 \pm 0.003$ |  |  |  |  |
| 8 | $3.0 \pm 0.144$ | $3.0 \pm 0.127$ | $2.8 \pm 0.089$ | $2.8 \pm 0.078$ | $2.6 \pm 0.047$ | $2.5 \pm 0.039$ | $2.4 \pm 0.032$ |  |  |  |  |
| 10 | $6.1 \pm 0.009$ | $6.0 \pm 0.008$ | $5.7 \pm 0.006$ | $5.6 \pm 0.005$ | $5.3 \pm 0.003$ | $5.2 \pm 0.003$ | $5.1 \pm 0.002$ |  |  |  |  |
| 11 | $4.2 \pm 0.03$ | $4.0 \pm 0.027$ | $3.8 \pm 0.018$ | $3.7 \pm 0.016$ | $3.4 \pm 0.009$ | $3.3 \pm 0.008$ | $3.1 \pm 0.006$ |  |  |  |  |
| 12 | $6.3 \pm 0.0$ | $6.4 \pm 0.0$ | $6.7 \pm 0.0$ | $6.8 \pm 0.0$ | $7.0 \pm 0.0$ | $7.1 \pm 0.0$ | $7.1 \pm 0.0$ |  |  |  |  |
| 15 | $5.2 \pm 0.024$ | $5.1 \pm 0.021$ | $5.0 \pm 0.016$ | $5.0 \pm 0.014$ | $4.9 \pm 0.009$ | $4.8 \pm 0.007$ | $4.7 \pm 0.005$ |  |  |  |  |
| 17 | $4.1 \pm 0.005$ | $4.1 \pm 0.004$ | $4.1 \pm 0.003$ | $4.0 \pm 0.003$ | $4.0 \pm 0.002$ | $4.0 \pm 0.002$ | $4.0 \pm 0.001$ |  |  |  |  |
| 19 | $5.3 \pm 0.033$ | $5.1 \pm 0.028$ | $4.7 \pm 0.018$ | $4.6 \pm 0.015$ | $4.1 \pm 0.008$ | $4.0 \pm 0.008$ | $3.7 \pm 0.008$ |  |  |  |  |
| 20 | $3.2 \pm 0.096$ | $3.1 \pm 0.082$ | $3.0 \pm 0.054$ | $3.0 \pm 0.046$ | $3.0 \pm 0.028$ | $2.9 \pm 0.025$ | $2.9 \pm 0.024$ |  |  |  |  |
| 21 | $4.3 \pm 0.071$ | $4.4 \pm 0.062$ | $4.7 \pm 0.044$ | $4.8 \pm 0.039$ | $5.1 \pm 0.025$ | $5.2 \pm 0.021$ | $5.3 \pm 0.018$ |  |  |  |  |
| 22 | $5.8 \pm 0.04$ | $5.7 \pm 0.035$ | $5.4 \pm 0.025$ | $5.3 \pm 0.022$ | $5.1 \pm 0.013$ | $5.0 \pm 0.011$ | $4.8 \pm 0.009$ |  |  |  |  |
| 23 | $5.2 \pm 0.034$ | $5.1 \pm 0.031$ | $4.7 \pm 0.024$ | $4.6 \pm 0.022$ | $4.3 \pm 0.015$ | $4.2 \pm 0.013$ | $4.0 \pm 0.01$ |  |  |  |  |
| 24 | $4.0 \pm 0.114$ | $4.0 \pm 0.102$ | $4.0 \pm 0.076$ | $4.0 \pm 0.067$ | $3.9 \pm 0.044$ | $3.8 \pm 0.039$ | $3.7 \pm 0.035$ |  |  |  |  |
| 25 | $8.0 \pm 0.06$ | $7.7 \pm 0.051$ | $7.0 \pm 0.034$ | $6.8 \pm 0.029$ | $6.2 \pm 0.019$ | $6.0 \pm 0.018$ | $5.8 \pm 0.018$ |  |  |  |  |
| 26 | $3.4 \pm 0.018$ | $3.3 \pm 0.015$ | $3.1 \pm 0.01$ | $3.0 \pm 0.008$ | $2.9 \pm 0.006$ | $2.8 \pm 0.006$ | $2.7 \pm 0.006$ |  |  |  |  |
| 28 | $6.8 \pm 0.019$ | $6.5 \pm 0.016$ | $5.8 \pm 0.01$ | $5.6 \pm 0.009$ | $4.9 \pm 0.005$ | $4.7 \pm 0.004$ | $4.3 \pm 0.004$ |  |  |  |  |
| 29 | $4.9 \pm 0.005$ | $4.9 \pm 0.004$ | $4.7 \pm 0.003$ | $4.7 \pm 0.003$ | $4.6 \pm 0.002$ | $4.6 \pm 0.002$ | $4.5 \pm 0.001$ |  |  |  |  |
| 30 | $4.6 \pm 0.045$ | $4.5 \pm 0.039$ | $4.2 \pm 0.028$ | $4.1 \pm 0.024$ | $3.9 \pm 0.015$ | $3.8 \pm 0.013$ | $3.8 \pm 0.011$ |  |  |  |  |
| 31 | $5.1 \pm 0.05$ | $5.3 \pm 0.042$ | $5.7 \pm 0.027$ | $5.9 \pm 0.023$ | $6.4 \pm 0.015$ | $6.5 \pm 0.015$ | $6.8 \pm 0.015$ |  |  |  |  |


| ID | MACS 1206 ACS Sérsic $n$ cont. |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{Z}_{850}$ |
| 32 | $5.6 \pm 0.021$ | $5.5 \pm 0.02$ | $5.3 \pm 0.016$ | $5.2 \pm 0.015$ | $5.0 \pm 0.011$ | $5.0 \pm 0.01$ | $5.0 \pm 0.008$ |
| 33 | $8 \pm 0.351$ | $7.7 \pm 0.312$ | $7.1 \pm 0.226$ | $6.8 \pm 0.2$ | $6.2 \pm 0.124$ | $6.0 \pm 0.102$ | $5.8 \pm 0.075$ |
| 35 | $4.2 \pm 0.024$ | $4.1 \pm 0.022$ | $3.9 \pm 0.017$ | $3.8 \pm 0.016$ | $3.6 \pm 0.011$ | $3.6 \pm 0.01$ | $3.5 \pm 0.007$ |
| 37 | $4.8 \pm 0.14$ | $5.0 \pm 0.126$ | $5.2 \pm 0.097$ | $5.3 \pm 0.087$ | $5.4 \pm 0.06$ | $5.5 \pm 0.052$ | $5.5 \pm 0.039$ |
| 38 | $5.7 \pm 0.029$ | $5.6 \pm 0.026$ | $5.5 \pm 0.019$ | $5.5 \pm 0.017$ | $5.3 \pm 0.012$ | $5.3 \pm 0.01$ | $5.1 \pm 0.008$ |
| 39 | $8.0 \pm 0.07$ | $7.7 \pm 0.063$ | $7.0 \pm 0.047$ | $6.8 \pm 0.042$ | $6.1 \pm 0.028$ | $5.8 \pm 0.024$ | $5.5 \pm 0.019$ |
| 40 | $4.9 \pm 0.151$ | $4.8 \pm 0.135$ | $4.6 \pm 0.099$ | $4.6 \pm 0.088$ | $4.5 \pm 0.058$ | $4.4 \pm 0.049$ | $4.4 \pm 0.037$ |
| 41 | $4.5 \pm 0.008$ | $4.4 \pm 0.007$ | $4.2 \pm 0.004$ | $4.2 \pm 0.003$ | $4.0 \pm 0.002$ | $4.0 \pm 0.002$ | $3.9 \pm 0.003$ |
| 42 | $8.0 \pm 0.013$ | $7.9 \pm 0.011$ | $7.7 \pm 0.007$ | $7.6 \pm 0.006$ | $7.4 \pm 0.003$ | $7.3 \pm 0.003$ | $7.1 \pm 0.003$ |
| 43 | $3.1 \pm 0.001$ | $3.1 \pm 0.001$ | $3.0 \pm 0.001$ | $2.9 \pm 0.001$ | $2.9 \pm 0.001$ | $2.8 \pm 0.0$ | $2.8 \pm 0.0$ |
| 45 | $6.0 \pm 0.038$ | $5.9 \pm 0.033$ | $5.6 \pm 0.023$ | $5.6 \pm 0.02$ | $5.3 \pm 0.013$ | $5.3 \pm 0.011$ | $5.2 \pm 0.01$ |
| 46 | $4.9 \pm 0.012$ | $4.8 \pm 0.011$ | $4.4 \pm 0.008$ | $4.3 \pm 0.007$ | $4.0 \pm 0.005$ | $3.9 \pm 0.004$ | $3.8 \pm 0.003$ |
| 47 | $6.8 \pm 0.035$ | $6.7 \pm 0.031$ | $6.6 \pm 0.023$ | $6.6 \pm 0.021$ | $6.5 \pm 0.014$ | $6.4 \pm 0.013$ | $6.3 \pm 0.01$ |
| 48 | $6.3 \pm 0.029$ | $6.2 \pm 0.026$ | $6.1 \pm 0.02$ | $6.0 \pm 0.019$ | $5.8 \pm 0.014$ | $5.7 \pm 0.012$ | $5.5 \pm 0.01$ |
| 50 | $2.8 \pm 0.008$ | $2.8 \pm 0.007$ | $2.8 \pm 0.004$ | $2.8 \pm 0.003$ | $2.8 \pm 0.002$ | $2.8 \pm 0.002$ | $2.8 \pm 0.003$ |
| 51 | $5.2 \pm 0.003$ | $5.2 \pm 0.003$ | $5.2 \pm 0.002$ | $5.2 \pm 0.002$ | $5.2 \pm 0.001$ | $5.2 \pm 0.001$ | $5.2 \pm 0.001$ |
| 52 | $6.0 \pm 0.003$ | $6.0 \pm 0.003$ | $5.9 \pm 0.001$ | $5.9 \pm 0.001$ | $6.0 \pm 0.001$ | $6.0 \pm 0.001$ | $6.0 \pm 0.001$ |
| 55 | $4.5 \pm 0.034$ | $4.5 \pm 0.029$ | $4.4 \pm 0.019$ | $4.3 \pm 0.017$ | $4.3 \pm 0.01$ | $4.3 \pm 0.009$ | $4.3 \pm 0.009$ |
| 56 | $4.5 \pm 0.026$ | $4.4 \pm 0.023$ | $4.0 \pm 0.015$ | $3.9 \pm 0.013$ | $3.6 \pm 0.007$ | $3.5 \pm 0.005$ | $3.3 \pm 0.005$ |
| 57 | $3.4 \pm 0.031$ | $3.4 \pm 0.027$ | $3.5 \pm 0.017$ | $3.5 \pm 0.015$ | $3.6 \pm 0.009$ | $3.6 \pm 0.009$ | $3.6 \pm 0.009$ |
| 58 | $5.0 \pm 0.17$ | $5.1 \pm 0.15$ | $5.5 \pm 0.112$ | $5.5 \pm 0.101$ | $5.7 \pm 0.075$ | $5.8 \pm 0.07$ | $5.8 \pm 0.066$ |
| 59 | $5.6 \pm 0.075$ | $5.3 \pm 0.066$ | $4.7 \pm 0.048$ | $4.5 \pm 0.042$ | $4.0 \pm 0.025$ | $3.8 \pm 0.021$ | $3.5 \pm 0.014$ |
| 61 | $3.2 \pm 0.06$ | $3.2 \pm 0.053$ | $3.3 \pm 0.037$ | $3.3 \pm 0.033$ | $3.3 \pm 0.02$ | $3.3 \pm 0.017$ | $3.3 \pm 0.014$ |
| 62 | $4.0 \pm 0.049$ | $4.0 \pm 0.044$ | $4.0 \pm 0.033$ | $4.0 \pm 0.03$ | $3.9 \pm 0.02$ | $3.9 \pm 0.018$ | $3.9 \pm 0.014$ |


| ID |  |  | MACS 1206 ACS Sérsic $n$ cont. |  | $\mathrm{I}_{814}$ | $\mathrm{Z}_{850}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{~B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $3.5 \pm 0.04$ |  |
| 63 | $2.1 \pm 0.156$ | $2.3 \pm 0.137$ | $2.7 \pm 0.097$ | $2.8 \pm 0.085$ | $3.2 \pm 0.054$ | $3.3 \pm 0.046$ | 3.5 |
| 64 | $6.3 \pm 0.097$ | $6.1 \pm 0.085$ | $5.6 \pm 0.059$ | $5.4 \pm 0.051$ | $4.9 \pm 0.03$ | $4.8 \pm 0.024$ | $4.5 \pm 0.018$ |
| 65 | $8.0 \pm 0.156$ | $7.8 \pm 0.143$ | $7.5 \pm 0.114$ | $7.4 \pm 0.104$ | $6.9 \pm 0.076$ | $6.8 \pm 0.067$ | $6.5 \pm 0.052$ |
| 68 | $4.7 \pm 0.072$ | $4.7 \pm 0.061$ | $4.8 \pm 0.038$ | $4.8 \pm 0.032$ | $4.8 \pm 0.023$ | $4.9 \pm 0.023$ | $4.9 \pm 0.026$ |
| 69 | $6.7 \pm 0.058$ | $6.6 \pm 0.046$ | $6.3 \pm 0.023$ | $6.2 \pm 0.018$ | $6.0 \pm 0.017$ | $6.0 \pm 0.02$ | $5.9 \pm 0.025$ |
| 70 | $5.6 \pm 0.07$ | $5.6 \pm 0.06$ | $5.5 \pm 0.038$ | $5.6 \pm 0.032$ | $5.6 \pm 0.019$ | $5.7 \pm 0.018$ | $5.8 \pm 0.018$ |
| 71 | $6.4 \pm 0.046$ | $6.3 \pm 0.041$ | $6.3 \pm 0.028$ | $6.2 \pm 0.024$ | $6.2 \pm 0.015$ | $6.2 \pm 0.013$ | $6.3 \pm 0.012$ |
| 72 | $7.9 \pm 0.258$ | $7.7 \pm 0.215$ | $7.3 \pm 0.128$ | $7.2 \pm 0.105$ | $6.9 \pm 0.069$ | $6.8 \pm 0.071$ | $6.6 \pm 0.08$ |
| 73 | $2.1 \pm 0.025$ | $2.2 \pm 0.022$ | $2.3 \pm 0.016$ | $2.4 \pm 0.015$ | $2.5 \pm 0.009$ | $2.6 \pm 0.008$ | $2.7 \pm 0.006$ |
| 74 | $8.0 \pm 0.484$ | $7.3 \pm 0.421$ | $5.9 \pm 0.285$ | $5.4 \pm 0.243$ | $4.2 \pm 0.132$ | $3.8 \pm 0.105$ | $3.3 \pm 0.083$ |
| 76 | $8.0 \pm 0.034$ | $7.7 \pm 0.031$ | $7.0 \pm 0.023$ | $6.8 \pm 0.021$ | $6.2 \pm 0.014$ | $5.9 \pm 0.012$ | $5.6 \pm 0.008$ |
| 77 | $4.3 \pm 0.113$ | $4.3 \pm 0.098$ | $4.2 \pm 0.067$ | $4.2 \pm 0.057$ | $4.1 \pm 0.032$ | $4.1 \pm 0.026$ | $4.0 \pm 0.022$ |
| 80 | $5.1 \pm 0.055$ | $5.0 \pm 0.048$ | $4.7 \pm 0.033$ | $4.6 \pm 0.028$ | $4.4 \pm 0.016$ | $4.3 \pm 0.014$ | $4.2 \pm 0.012$ |
| 81 | $3.1 \pm 0.036$ | $3.0 \pm 0.032$ | $2.8 \pm 0.023$ | $2.8 \pm 0.021$ | $2.6 \pm 0.013$ | $2.5 \pm 0.011$ | $2.4 \pm 0.009$ |
| 82 | $6.7 \pm 0.081$ | $6.4 \pm 0.073$ | $5.8 \pm 0.055$ | $5.6 \pm 0.049$ | $5.0 \pm 0.034$ | $4.8 \pm 0.03$ | $4.6 \pm 0.024$ |
| 83 | $7.4 \pm 0.126$ | $7.1 \pm 0.112$ | $6.5 \pm 0.08$ | $6.4 \pm 0.071$ | $5.9 \pm 0.043$ | $5.7 \pm 0.036$ | $5.5 \pm 0.026$ |
| 84 | $3.8 \pm 0.018$ | $3.8 \pm 0.015$ | $3.8 \pm 0.008$ | $3.8 \pm 0.006$ | $3.8 \pm 0.006$ | $3.8 \pm 0.007$ | $3.8 \pm 0.008$ |
| 85 | $3.7 \pm 0.132$ | $3.7 \pm 0.117$ | $3.7 \pm 0.086$ | $3.7 \pm 0.076$ | $3.7 \pm 0.052$ | $3.7 \pm 0.046$ | $3.7 \pm 0.039$ |
| 87 | $5.8 \pm 0.032$ | $5.9 \pm 0.027$ | $5.9 \pm 0.018$ | $5.9 \pm 0.016$ | $5.9 \pm 0.011$ | $5.8 \pm 0.01$ | $5.8 \pm 0.011$ |
| 88 | $2.8 \pm 0.043$ | $2.9 \pm 0.038$ | $3.1 \pm 0.025$ | $3.1 \pm 0.022$ | $3.3 \pm 0.013$ | $3.3 \pm 0.011$ | $3.3 \pm 0.011$ |
| 89 | $4.8 \pm 0.015$ | $4.6 \pm 0.013$ | $4.1 \pm 0.009$ | $3.9 \pm 0.008$ | $3.4 \pm 0.004$ | $3.3 \pm 0.003$ | $3.0 \pm 0.003$ |
| 90 | $4.4 \pm 0.178$ | $4.3 \pm 0.158$ | $4.0 \pm 0.114$ | $3.9 \pm 0.101$ | $3.7 \pm 0.062$ | $3.7 \pm 0.052$ | $3.5 \pm 0.039$ |
| 91 | $8.0 \pm 0.011$ | $7.6 \pm 0.009$ | $6.6 \pm 0.006$ | $6.3 \pm 0.006$ | $5.4 \pm 0.003$ | $5.1 \pm 0.003$ | $4.6 \pm 0.002$ |
| 93 | $5.3 \pm 0.196$ | $5.3 \pm 0.174$ | $5.3 \pm 0.126$ | $5.3 \pm 0.111$ | $5.3 \pm 0.071$ | $5.3 \pm 0.061$ | $5.2 \pm 0.049$ |


| ID | MACS 1206 ACS Sérsic $n$ cont. |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{z}_{850}$ |
| 94 | $8.0 \pm 0.046$ | $7.5 \pm 0.04$ | $6.4 \pm 0.028$ | $6.0 \pm 0.024$ | $5.0 \pm 0.014$ | $4.6 \pm 0.012$ | $4.1 \pm 0.009$ |
| 96 | $3.8 \pm 0.067$ | $3.7 \pm 0.057$ | $3.5 \pm 0.035$ | $3.4 \pm 0.029$ | $3.2 \pm 0.016$ | $3.2 \pm 0.015$ | $3.1 \pm 0.017$ |
| 97 | $6.3 \pm 0.064$ | $6.1 \pm 0.057$ | $5.8 \pm 0.041$ | $5.7 \pm 0.036$ | $5.3 \pm 0.022$ | $5.2 \pm 0.018$ | $5.0 \pm 0.013$ |
| 98 | $4.3 \pm 0.046$ | $4.3 \pm 0.039$ | $4.3 \pm 0.025$ | $4.4 \pm 0.02$ | $4.4 \pm 0.011$ | $4.4 \pm 0.011$ | $4.4 \pm 0.013$ |
| 99 | $3.9 \pm 0.047$ | $3.7 \pm 0.042$ | $3.4 \pm 0.03$ | $3.3 \pm 0.026$ | $3.0 \pm 0.018$ | $2.8 \pm 0.016$ | $2.6 \pm 0.014$ |
| 100 | $6.0 \pm 0.016$ | $5.9 \pm 0.014$ | $5.6 \pm 0.009$ | $5.5 \pm 0.008$ | $5.3 \pm 0.004$ | $5.2 \pm 0.003$ | $5.0 \pm 0.004$ |
| 101 | $4.8 \pm 0.016$ | $4.9 \pm 0.014$ | $5.0 \pm 0.01$ | $5.1 \pm 0.009$ | $5.2 \pm 0.007$ | $5.3 \pm 0.006$ | $5.4 \pm 0.006$ |
| 104 | $3.3 \pm 0.046$ | $3.3 \pm 0.04$ | $3.3 \pm 0.029$ | $3.2 \pm 0.026$ | $3.2 \pm 0.016$ | $3.2 \pm 0.014$ | $3.2 \pm 0.011$ |

Table C.6.: Measurements of $n$ in the WFC3/IR bands for the selected early-type galaxies in MACS 1206.

| ID | MACS 1206 WFC3/IR Sérsic $n$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $Y \mathrm{~J}_{110}$ | $\mathrm{J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| 1 | $3.1 \pm 0.003$ | $3.0 \pm 0.003$ | $3.0 \pm 0.003$ | $3.0 \pm 0.004$ | $3.2 \pm 0.005$ |
| 3 | $4.2 \pm 0.002$ | $4.1 \pm 0.002$ | $4.2 \pm 0.002$ | $4.2 \pm 0.003$ | $4.3 \pm 0.004$ |
| 4 | $3.7 \pm 0.024$ | $3.8 \pm 0.024$ | $3.8 \pm 0.021$ | $3.9 \pm 0.015$ | $4.0 \pm 0.023$ |
| 5 | $2.6 \pm 0.005$ | $2.6 \pm 0.005$ | $2.6 \pm 0.005$ | $2.6 \pm 0.006$ | $2.6 \pm 0.009$ |
| 6 | $4.1 \pm 0.002$ | $4.1 \pm 0.003$ | $4.2 \pm 0.003$ | $4.2 \pm 0.004$ | $4.4 \pm 0.006$ |
| 8 | $2.4 \pm 0.031$ | $2.3 \pm 0.032$ | $2.2 \pm 0.032$ | $2.2 \pm 0.032$ | $2.2 \pm 0.044$ |
| 10 | $5.1 \pm 0.002$ | $5.1 \pm 0.002$ | $5.1 \pm 0.002$ | $5.3 \pm 0.002$ | $5.5 \pm 0.003$ |
| 11 | $3.0 \pm 0.006$ | $2.9 \pm 0.006$ | $2.9 \pm 0.006$ | $2.8 \pm 0.006$ | $2.9 \pm 0.008$ |
| 12 | $7.2 \pm 0.0$ | $7.1 \pm 0.0$ | $7.1 \pm 0.0$ | $6.9 \pm 0.0$ | $6.7 \pm 0.0$ |
| 15 | $4.5 \pm 0.004$ | $4.4 \pm 0.004$ | $4.2 \pm 0.004$ | $4.0 \pm 0.004$ | $3.7 \pm 0.005$ |
| 17 | $4.0 \pm 0.001$ | $4.0 \pm 0.001$ | $4.0 \pm 0.001$ | $4.1 \pm 0.001$ | $4.2 \pm 0.002$ |
| 19 | $3.5 \pm 0.008$ | $3.3 \pm 0.007$ | $3.2 \pm 0.005$ | $3.1 \pm 0.005$ | $3.0 \pm 0.014$ |
| 20 | $3.0 \pm 0.025$ | $3.0 \pm 0.024$ | $3.1 \pm 0.024$ | $3.2 \pm 0.032$ | $3.4 \pm 0.058$ |
| 21 | $5.3 \pm 0.018$ | $5.3 \pm 0.019$ | $5.3 \pm 0.021$ | $5.2 \pm 0.026$ | $4.9 \pm 0.037$ |
| 22 | $4.7 \pm 0.008$ | $4.6 \pm 0.008$ | $4.5 \pm 0.008$ | $4.4 \pm 0.009$ | $4.4 \pm 0.014$ |
| 23 | $3.8 \pm 0.008$ | $3.7 \pm 0.007$ | $3.6 \pm 0.006$ | $3.5 \pm 0.008$ | $3.5 \pm 0.01$ |
| 24 | $3.5 \pm 0.039$ | $3.2 \pm 0.048$ | $3.0 \pm 0.058$ | $2.5 \pm 0.077$ | $2.0 \pm 0.1$ |
| 25 | $5.6 \pm 0.017$ | $5.5 \pm 0.016$ | $5.6 \pm 0.015$ | $5.8 \pm 0.021$ | $6.1 \pm 0.039$ |
| 26 | $2.7 \pm 0.007$ | $2.6 \pm 0.007$ | $2.6 \pm 0.008$ | $2.7 \pm 0.01$ | $2.8 \pm 0.016$ |
| 28 | $4.1 \pm 0.004$ | $3.9 \pm 0.004$ | $3.9 \pm 0.003$ | $3.9 \pm 0.004$ | $4.1 \pm 0.009$ |
| 29 | $4.6 \pm 0.001$ | $4.7 \pm 0.001$ | $4.7 \pm 0.001$ | $5.0 \pm 0.001$ | $5.2 \pm 0.002$ |
| 30 | $3.8 \pm 0.01$ | $3.9 \pm 0.01$ | $4.0 \pm 0.01$ | $4.3 \pm 0.013$ | $4.7 \pm 0.021$ |
| 31 | $7.1 \pm 0.015$ | $7.3 \pm 0.014$ | $7.5 \pm 0.013$ | $7.6 \pm 0.018$ | $7.7 \pm 0.035$ |
| 32 | $5.0 \pm 0.007$ | $5.0 \pm 0.006$ | $5.1 \pm 0.007$ | $5.3 \pm 0.008$ | $5.6 \pm 0.011$ |
| 33 | $5.6 \pm 0.065$ | $5.5 \pm 0.066$ | $5.5 \pm 0.07$ | $5.6 \pm 0.084$ | $5.9 \pm 0.119$ |
| 35 | $3.4 \pm 0.006$ | $3.4 \pm 0.005$ | $3.5 \pm 0.005$ | $3.6 \pm 0.006$ | $3.7 \pm 0.007$ |
| 37 | $5.5 \pm 0.031$ | $5.4 \pm 0.028$ | $5.2 \pm 0.026$ | $4.9 \pm 0.027$ | $4.4 \pm 0.034$ |
| 38 | $4.9 \pm 0.007$ | $4.8 \pm 0.006$ | $4.6 \pm 0.005$ | $4.3 \pm 0.005$ | $4.0 \pm 0.008$ |
| 39 | $5.1 \pm 0.016$ | $4.9 \pm 0.014$ | $4.8 \pm 0.012$ | $4.7 \pm 0.012$ | $4.7 \pm 0.02$ |
| 40 | $4.4 \pm 0.031$ | $4.4 \pm 0.03$ | $4.4 \pm 0.03$ | $4.5 \pm 0.037$ | $4.7 \pm 0.054$ |
| 41 | $3.9 \pm 0.003$ | $3.8 \pm 0.003$ | $3.9 \pm 0.002$ | $3.9 \pm 0.002$ | $4.0 \pm 0.004$ |
| 42 | $7.0 \pm 0.003$ | $6.8 \pm 0.003$ | $6.7 \pm 0.003$ | $6.6 \pm 0.003$ | $6.4 \pm 0.004$ |
| 43 | $2.8 \pm 0.0$ | $2.8 \pm 0.0$ | $2.8 \pm 0.0$ | $2.9 \pm 0.0$ | $3 \pm 0.0$ |
| 45 | $5.1 \pm 0.009$ | $5.1 \pm 0.008$ | $5.1 \pm 0.007$ | $5.2 \pm 0.009$ | $5.3 \pm 0.017$ |
| 46 | $3.7 \pm 0.003$ | $3.6 \pm 0.002$ | $3.6 \pm 0.002$ | $3.7 \pm 0.003$ | $3.9 \pm 0.005$ |
| 47 | $6.3 \pm 0.01$ | $6.2 \pm 0.009$ | $6.2 \pm 0.009$ | $6.2 \pm 0.011$ | $6.2 \pm 0.015$ |
| 48 | $5.3 \pm 0.009$ | $5.1 \pm 0.008$ | $4.9 \pm 0.008$ | $4.7 \pm 0.008$ | $4.4 \pm 0.01$ |
| 50 | $2.8 \pm 0.003$ | $2.8 \pm 0.002$ | $2.8 \pm 0.002$ | $2.9 \pm 0.003$ | $2.9 \pm 0.006$ |


| ID | MACS 1206 WFC3/IR Sérsic $n$ cont. |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | YJ 110 | $\mathrm{J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| 51 | $5.2 \pm 0.001$ | $5.3 \pm 0.001$ | $5.3 \pm 0.001$ | $5.4 \pm 0.001$ | $5.4 \pm 0.001$ |
| 52 | $6.1 \pm 0.001$ | $6.2 \pm 0.001$ | $6.3 \pm 0.001$ | $6.5 \pm 0.001$ | $6.7 \pm 0.003$ |
| 55 | $4.3 \pm 0.01$ | $4.4 \pm 0.01$ | $4.5 \pm 0.011$ | $4.7 \pm 0.014$ | $5.0 \pm 0.021$ |
| 56 | $3.1 \pm 0.006$ | $3.0 \pm 0.007$ | $3 \pm 0.007$ | $3.0 \pm 0.006$ | $3.0 \pm 0.008$ |
| 57 | $3.7 \pm 0.01$ | $3.7 \pm 0.01$ | $3.6 \pm 0.01$ | $3.6 \pm 0.011$ | $3.5 \pm 0.017$ |
| 58 | $5.7 \pm 0.066$ | $5.5 \pm 0.067$ | $5.3 \pm 0.071$ | $4.8 \pm 0.088$ | $4.1 \pm 0.123$ |
| 59 | $3.4 \pm 0.012$ | $3.3 \pm 0.011$ | $3.3 \pm 0.011$ | $3.5 \pm 0.013$ | $3.8 \pm 0.021$ |
| 61 | $3.3 \pm 0.014$ | $3.3 \pm 0.014$ | $3.2 \pm 0.014$ | $3.2 \pm 0.015$ | $3.1 \pm 0.022$ |
| 62 | $3.8 \pm 0.011$ | $3.8 \pm 0.011$ | $3.7 \pm 0.01$ | $3.7 \pm 0.012$ | $3.6 \pm 0.016$ |
| 63 | $3.6 \pm 0.039$ | $3.6 \pm 0.04$ | $3.5 \pm 0.04$ | $3.3 \pm 0.041$ | $3.0 \pm 0.053$ |
| 64 | $4.3 \pm 0.017$ | $4.2 \pm 0.015$ | $4.2 \pm 0.015$ | $4.2 \pm 0.02$ | $4.4 \pm 0.037$ |
| 65 | $6.2 \pm 0.042$ | $6.0 \pm 0.036$ | $5.8 \pm 0.034$ | $5.5 \pm 0.038$ | $5.3 \pm 0.05$ |
| 68 | $5.0 \pm 0.026$ | $5.1 \pm 0.024$ | $5.1 \pm 0.022$ | $5.2 \pm 0.03$ | $5.4 \pm 0.055$ |
| 69 | $5.8 \pm 0.026$ | $5.8 \pm 0.022$ | $5.8 \pm 0.015$ | $6.0 \pm 0.011$ | $6.2 \pm 0.038$ |
| 70 | $5.9 \pm 0.019$ | $6.1 \pm 0.018$ | $6.3 \pm 0.015$ | $6.6 \pm 0.015$ | $7.0 \pm 0.03$ |
| 71 | $6.3 \pm 0.012$ | $6.4 \pm 0.012$ | $6.5 \pm 0.012$ | $6.6 \pm 0.014$ | $6.8 \pm 0.021$ |
| 72 | $6.6 \pm 0.078$ | $6.6 \pm 0.065$ | $6.6 \pm 0.051$ | $6.8 \pm 0.084$ | $7.2 \pm 0.189$ |
| 73 | $2.8 \pm 0.005$ | $2.9 \pm 0.004$ | $2.9 \pm 0.004$ | $3.1 \pm 0.005$ | $3.2 \pm 0.007$ |
| 74 | $3.1 \pm 0.085$ | $3.2 \pm 0.09$ | $3.4 \pm 0.095$ | $4.2 \pm 0.125$ | $5.4 \pm 0.212$ |
| 76 | $5.4 \pm 0.007$ | $5.3 \pm 0.006$ | $5.3 \pm 0.006$ | $5.5 \pm 0.007$ | $5.8 \pm 0.01$ |
| 77 | $3.9 \pm 0.023$ | $3.9 \pm 0.023$ | $3.9 \pm 0.022$ | $3.8 \pm 0.018$ | $3.8 \pm 0.03$ |
| 80 | $4.2 \pm 0.011$ | $4.2 \pm 0.01$ | $4.3 \pm 0.009$ | $4.4 \pm 0.011$ | $4.7 \pm 0.022$ |
| 81 | $2.4 \pm 0.007$ | $2.3 \pm 0.007$ | $2.3 \pm 0.007$ | $2.4 \pm 0.009$ | $2.5 \pm 0.014$ |
| 82 | $4.4 \pm 0.02$ | $4.4 \pm 0.018$ | $4.4 \pm 0.017$ | $4.6 \pm 0.021$ | $4.9 \pm 0.031$ |
| 83 | $5.4 \pm 0.023$ | $5.4 \pm 0.022$ | $5.4 \pm 0.022$ | $5.7 \pm 0.025$ | $6.2 \pm 0.039$ |
| 84 | $3.7 \pm 0.008$ | $3.7 \pm 0.006$ | $3.7 \pm 0.005$ | $3.7 \pm 0.007$ | $3.7 \pm 0.017$ |
| 85 | $3.6 \pm 0.036$ | $3.6 \pm 0.034$ | $3.5 \pm 0.032$ | $3.3 \pm 0.035$ | $3.1 \pm 0.052$ |
| 87 | $5.7 \pm 0.011$ | $5.6 \pm 0.012$ | $5.4 \pm 0.013$ | $5.2 \pm 0.017$ | $4.9 \pm 0.026$ |
| 88 | $3.3 \pm 0.011$ | $3.3 \pm 0.011$ | $3.2 \pm 0.01$ | $3.1 \pm 0.01$ | $2.9 \pm 0.015$ |
| 89 | $2.8 \pm 0.003$ | $2.7 \pm 0.003$ | $2.7 \pm 0.002$ | $2.7 \pm 0.002$ | $2.8 \pm 0.004$ |
| 90 | $3.5 \pm 0.034$ | $3.4 \pm 0.033$ | $3.4 \pm 0.033$ | $3.5 \pm 0.037$ | $3.6 \pm 0.053$ |
| 91 | $4.3 \pm 0.002$ | $4.1 \pm 0.002$ | $4.0 \pm 0.002$ | $4.1 \pm 0.002$ | $4.4 \pm 0.003$ |
| 93 | $5.2 \pm 0.047$ | $5.1 \pm 0.047$ | $5.0 \pm 0.048$ | $4.8 \pm 0.052$ | $4.6 \pm 0.068$ |
| 94 | $3.7 \pm 0.008$ | $3.5 \pm 0.007$ | $3.4 \pm 0.006$ | $3.6 \pm 0.007$ | $4.0 \pm 0.013$ |
| 96 | $3.0 \pm 0.02$ | $3.0 \pm 0.022$ | $3.0 \pm 0.023$ | $3.1 \pm 0.029$ | $3.2 \pm 0.046$ |
| 97 | $4.8 \pm 0.011$ | $4.7 \pm 0.01$ | $4.7 \pm 0.011$ | $4.6 \pm 0.013$ | $4.6 \pm 0.021$ |
| 98 | $4.4 \pm 0.016$ | $4.3 \pm 0.017$ | $4.3 \pm 0.017$ | $4.3 \pm 0.019$ | $4.2 \pm 0.027$ |
| 99 | $2.5 \pm 0.015$ | $2.4 \pm 0.015$ | $2.3 \pm 0.016$ | $2.2 \pm 0.02$ | $2.2 \pm 0.028$ |


| ID | MACS 1206 WFC3/IR Sérsic $n$ cont. |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $\mathrm{YJ}_{110}$ | $\mathrm{~J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| 100 | $4.9 \pm 0.004$ | $4.8 \pm 0.004$ | $4.8 \pm 0.004$ | $4.8 \pm 0.004$ | $4.8 \pm 0.005$ |
| 101 | $5.4 \pm 0.006$ | $5.5 \pm 0.005$ | $5.5 \pm 0.005$ | $5.5 \pm 0.005$ | $5.5 \pm 0.009$ |
| 104 | $3.1 \pm 0.01$ | $3.1 \pm 0.009$ | $3.0 \pm 0.008$ | $3.0 \pm 0.008$ | $2.9 \pm 0.014$ |

Table C.7.: Magnitude assessments in the ACS bands for the fitted early-type galaxies in MACS 0416

| ID | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{z}_{850}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $23.5 \pm 0.0295$ | $22.8 \pm 0.0173$ | $21.8 \pm 0.011$ | $21.5 \pm 0.0104$ | $21.0 \pm 0.0067$ | $20.9 \pm 0.0055$ | $20.6 \pm 0.0042$ |
| 21 | 23.5 |  |  |  |  |  |  |
| 25 | $23.0 \pm 0.7552$ | $22.5 \pm 0.4667$ | $21.2 \pm 0.2749$ | $21.1 \pm 0.2653$ | $20.6 \pm 0.1841$ | $20.4 \pm 0.1563$ | $20.1 \pm 0.1731$ |
| 26 | $22.2 \pm 0.0022$ | $21.8 \pm 0.0018$ | $20.5 \pm 0.0014$ | $20.3 \pm 0.0014$ | $19.8 \pm 0.0011$ | $19.6 \pm 0.001$ | $19.3 \pm 0.0009$ |
| 30 | $24.4 \pm 0.0097$ | $24.5 \pm 0.008$ | $23.6 \pm 0.0057$ | $23.4 \pm 0.0045$ | $22.9 \pm 0.003$ | $22.8 \pm 0.0022$ | $22.5 \pm 0.0019$ |
| 31 | $23.7 \pm 0.0109$ | $23.1 \pm 0.0068$ | $21.8 \pm 0.0046$ | $21.5 \pm 0.0038$ | $21.0 \pm 0.0026$ | $20.8 \pm 0.0022$ | $20.5 \pm 0.0012$ |
| 33 | $23.6 \pm 0.0008$ | $23.1 \pm 0.001$ | $21.8 \pm 0.0008$ | $21.7 \pm 0.0005$ | $21.1 \pm 0.0004$ | $20.9 \pm 0.0003$ | $20.6 \pm 0.0002$ |
| 38 | $24.3 \pm 0.0029$ | $23.9 \pm 0.0025$ | $22.6 \pm 0.0021$ | $22.5 \pm 0.002$ | $22.0 \pm 0.0015$ | $21.8 \pm 0.0011$ | $21.4 \pm 0.0008$ |
| 43 | $24.0 \pm 0.0205$ | $23.9 \pm 0.0161$ | $23.2 \pm 0.0101$ | $23.1 \pm 0.0098$ | $22.6 \pm 0.0065$ | $22.4 \pm 0.0052$ | $22.1 \pm 0.0045$ |
| 45 | $22.5 \pm 0.0004$ | $22.1 \pm 0.0002$ | $20.8 \pm 0.0001$ | $20.6 \pm 0.0001$ | $20.1 \pm 0.0001$ | $19.9 \pm 0.0001$ | $19.5 \pm 0.0001$ |
| 50 | $22.6 \pm 0.0022$ | $22.2 \pm 0.0023$ | $21.2 \pm 0.0021$ | $21.1 \pm 0.0021$ | $20.6 \pm 0.0018$ | $20.4 \pm 0.0013$ | $20.1 \pm 0.0008$ |
| 53 | $23.3 \pm 0.0015$ | $22.8 \pm 0.001$ | $21.4 \pm 0.001$ | $21.3 \pm 0.0007$ | $20.8 \pm 0.0005$ | $20.5 \pm 0.0004$ | $20.2 \pm 0.0004$ |
| 54 | $24.1 \pm 0.0017$ | $23.6 \pm 0.0015$ | $22.4 \pm 0.0014$ | $22.3 \pm 0.0013$ | $21.7 \pm 0.0011$ | $21.5 \pm 0.0009$ | $21.2 \pm 0.0008$ |
| 55 | $24.7 \pm 0.0291$ | $23.9 \pm 0.0227$ | $22.9 \pm 0.02$ | $22.7 \pm 0.0185$ | $22.3 \pm 0.0184$ | $22.1 \pm 0.017$ | $21.8 \pm 0.0179$ |
| 57 | $23.1 \pm 0.0017$ | $22.8 \pm 0.002$ | $21.6 \pm 0.0016$ | $21.5 \pm 0.0014$ | $20.9 \pm 0.0009$ | $20.7 \pm 0.0008$ | $20.4 \pm 0.0002$ |
| 63 | $22.8 \pm 0.0027$ | $22.5 \pm 0.0025$ | $21.3 \pm 0.0018$ | $21.2 \pm 0.0022$ | $20.6 \pm 0.0011$ | $20.4 \pm 0.0007$ | $20.1 \pm 0.0001$ |
| 65 | $22.2 \pm 0.0004$ | $21.8 \pm 0.0005$ | $20.6 \pm 0.0004$ | $20.4 \pm 0.0002$ | $19.9 \pm 0.0004$ | $19.7 \pm 0.0004$ | $19.3 \pm 0.0003$ |
| 66 | $21.3 \pm 0.0018$ | $21.0 \pm 0.0008$ | $19.6 \pm 0.0006$ | $19.4 \pm 0.0007$ | $18.8 \pm 0.0002$ | $18.5 \pm 0.0002$ | $18.1 \pm 0.0005$ |
| 68 | $24.0 \pm 0.0034$ | $23.5 \pm 0.0018$ | $22.3 \pm 0.0015$ | $22.1 \pm 0.0012$ | $21.6 \pm 0.0008$ | $21.4 \pm 0.0005$ | $21.0 \pm 0.0004$ |
| 70 | $22.9 \pm 0.0004$ | $22.5 \pm 0.0003$ | $21.2 \pm 0.0002$ | $21.1 \pm 0.0002$ | $20.6 \pm 0.0001$ | $20.4 \pm 0.0001$ | $20.0 \pm 0.0001$ |
| 71 | $23.6 \pm 0.008$ | $23.3 \pm 0.0054$ | $22.1 \pm 0.0031$ | $22.0 \pm 0.0029$ | $21.5 \pm 0.002$ | $21.3 \pm 0.0017$ | $21.0 \pm 0.0013$ |
| 85 | $21.9 \pm 0.0178$ | $21.5 \pm 0.0127$ | $20.2 \pm 0.0088$ | $20.0 \pm 0.0082$ | $19.5 \pm 0.0056$ | $19.3 \pm 0.0048$ | $18.9 \pm 0.0036$ |
| 87 | $22.0 \pm 0.0005$ | $21.7 \pm 0.0003$ | $20.5 \pm 0.0002$ | $20.3 \pm 0.0002$ | $19.8 \pm 0.0003$ | $19.6 \pm 0.0002$ | $19.2 \pm 0.0003$ |
| 88 | $22.6 \pm 0.0023$ | $22.1 \pm 0.0018$ | $20.6 \pm 0.0015$ | $20.5 \pm 0.0014$ | $20.0 \pm 0.0013$ | $19.7 \pm 0.001$ | $19.4 \pm 0.0009$ |


| ID |  | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | MACS 0416 ACS AB Magnitudes cont. |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ |  |  |  |
| 93 | $23.6 \pm 0.0003$ | $23.5 \pm 0.0002$ | $22.4 \pm 0.0001$ | $22.3 \pm 0.0003$ | $21.8 \pm 0.0001$ | $21.6 \pm 0.0001$ | $21.3 \pm 0.0001$ |
| 94 | $22.6 \pm 0.0018$ | $22.3 \pm 0.0013$ | $21.1 \pm 0.0009$ | $20.9 \pm 0.0008$ | $20.4 \pm 0.0006$ | $20.2 \pm 0.0005$ | $19.8 \pm 0.0004$ |
| 95 | $23.8 \pm 0.0024$ | $23.7 \pm 0.0021$ | $22.6 \pm 0.0019$ | $22.5 \pm 0.0018$ | $22.0 \pm 0.0011$ | $21.8 \pm 0.0009$ | $21.5 \pm 0.0002$ |
| 96 | $22.4 \pm 0.0252$ | $21.9 \pm 0.0156$ | $20.7 \pm 0.0105$ | $20.5 \pm 0.0103$ | $20.1 \pm 0.0081$ | $19.8 \pm 0.0063$ | $19.6 \pm 0.005$ |
| 98 | $24.9 \pm 0.0043$ | $24.1 \pm 0.0025$ | $22.8 \pm 0.0013$ | $22.7 \pm 0.0009$ | $22.1 \pm 0.0012$ | $21.9 \pm 0.0008$ | $21.6 \pm 0.0008$ |
| 102 | $23.7 \pm 0.0764$ | $23.4 \pm 0.0434$ | $22.2 \pm 0.0208$ | $22.1 \pm 0.0198$ | $21.6 \pm 0.0125$ | $21.3 \pm 0.0099$ | $21.0 \pm 0.0086$ |
| 103 | $22.5 \pm 0.0014$ | $22.4 \pm 0.0009$ | $21.4 \pm 0.0005$ | $21.2 \pm 0.0005$ | $20.7 \pm 0.0003$ | $20.5 \pm 0.0003$ | $20.2 \pm 0.0002$ |
| 104 | $22.1 \pm 0.0005$ | $21.7 \pm 0.0005$ | $20.4 \pm 0.0004$ | $20.2 \pm 0.0004$ | $19.7 \pm 0.0002$ | $19.5 \pm 0.0001$ | $19.1 \pm 0.0$ |
| 105 | $23.3 \pm 0.0276$ | $22.8 \pm 0.0152$ | $21.5 \pm 0.0086$ | $21.3 \pm 0.0084$ | $20.8 \pm 0.0069$ | $20.5 \pm 0.0062$ | $20.2 \pm 0.0061$ |
| 107 | $22.4 \pm 0.0006$ | $22.0 \pm 0.0004$ | $20.6 \pm 0.0003$ | $20.5 \pm 0.0003$ | $19.9 \pm 0.0001$ | $19.7 \pm 0.0001$ | $19.3 \pm 0.0001$ |
| 109 | $22.6 \pm 0.0012$ | $22.1 \pm 0.0007$ | $20.7 \pm 0.0005$ | $20.6 \pm 0.0005$ | $20.1 \pm 0.0009$ | $19.9 \pm 0.001$ | $19.5 \pm 0.0009$ |
| 113 | $23.3 \pm 0.0198$ | $22.0 \pm 0.0105$ | $20.7 \pm 0.0073$ | $20.6 \pm 0.0066$ | $20.0 \pm 0.0046$ | $19.8 \pm 0.0038$ | $19.5 \pm 0.0028$ |
| 115 | $23.5 \pm 0.0814$ | $22.6 \pm 0.0342$ | $21.3 \pm 0.0188$ | $21.0 \pm 0.0181$ | $20.7 \pm 0.0128$ | $20.3 \pm 0.0097$ | $20.1 \pm 0.0102$ |
| 133 | $23.4 \pm 0.0803$ | $23.0 \pm 0.054$ | $21.6 \pm 0.0313$ | $21.5 \pm 0.0279$ | $20.9 \pm 0.0171$ | $20.7 \pm 0.0142$ | $20.4 \pm 0.0166$ |
| 154 | $24.0 \pm 0.0012$ | $23.9 \pm 0.0009$ | $23.1 \pm 0.0015$ | $23.0 \pm 0.0009$ | $22.5 \pm 0.0007$ | $22.3 \pm 0.0007$ | $22.0 \pm 0.0005$ |

## C. MegaMorph results

Table C.8.: Magnitude assessments in the WFC3/IR bands for the fitted early-type galaxies in MACS 0416

| ID | MACS 0416 WFC3/IR AB Magnitudes |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $Y J_{110}$ | $\mathrm{J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| 21 | $20.5 \pm 0.0039$ | $20.3 \pm 0.0039$ | $20.3 \pm 0.0035$ | $20.1 \pm 0.0028$ | $20.0 \pm 0.0044$ |
| 25 | $20.0 \pm 0.1749$ | $19.9 \pm 0.7321$ | $19.9 \pm 0.1774$ | $19.8 \pm 0.1822$ | $19.6 \pm 0.2743$ |
| 26 | $19.2 \pm 0.0007$ | $19.2 \pm 0.0006$ | $19.2 \pm 0.0005$ | $18.9 \pm 0.0006$ | $18.7 \pm 8.0$ |
| 30 | $22.4 \pm 0.0018$ | $22.4 \pm 0.0012$ | $22.3 \pm 0.0011$ | $22.2 \pm 0.0011$ | $22.1 \pm 0.0$ |
| 31 | $20.4 \pm 0.0008$ | $20.3 \pm 0.0008$ | $20.1 \pm 0.0128$ | $19.9 \pm 0.0892$ | $20.0 \pm 0.0011$ |
| 33 | $20.5 \pm 0.0001$ | $20.4 \pm 0.0001$ | $20.3 \pm 0.0001$ | $20.2 \pm 0.0001$ | $20.0 \pm 1.848$ |
| 38 | $21.4 \pm 0.0005$ | $21.3 \pm 0.0003$ | $21.2 \pm 0.0002$ | $21.1 \pm 0.0002$ | $21.0 \pm 3.35$ |
| 43 | $22.0 \pm 0.0028$ | $21.9 \pm 0.0022$ | $21.9 \pm 0.0022$ | $21.7 \pm 0.0019$ | $21.6 \pm 0.0026$ |
| 45 | $19.5 \pm 0.0$ | $19.3 \pm 0.0$ | $19.2 \pm 0.0$ | $19.1 \pm 0.0$ | $18.9 \pm 3.354$ |
| 50 | $20.0 \pm 0.0005$ | $19.9 \pm 0.0003$ | $19.8 \pm 0.0003$ | $19.6 \pm 0.0006$ | $19.4 \pm 0.0015$ |
| 53 | $20.1 \pm 0.0001$ | $20.0 \pm 0.0001$ | $19.9 \pm 0.0001$ | $19.7 \pm 0.0001$ | $19.6 \pm 1.3645$ |
| 54 | $21.2 \pm 0.0006$ | $21.1 \pm 0.0005$ | $21.0 \pm 0.0004$ | $20.9 \pm 0.0005$ | $20.8 \pm 8.3619$ |
| 55 | $21.8 \pm 0.0179$ | $21.7 \pm 0.0181$ | $21.7 \pm 0.0185$ | $21.7 \pm 0.0168$ | $21.7 \pm 0.0137$ |
| 57 | $20.3 \pm 0.0002$ | $20.2 \pm 0.0004$ | $20.1 \pm 0.0004$ | $19.9 \pm 0.0003$ | $19.7 \pm 3.06$ |
| 63 | $20.0 \pm 0.0005$ | $19.9 \pm 0.0006$ | $19.8 \pm 0.0006$ | $19.6 \pm 0.0004$ | $19.5 \pm 3.81$ |
| 65 | $19.2 \pm 0.0002$ | $19.1 \pm 0.0$ | $19.0 \pm 0.0$ | $18.8 \pm 0.0002$ | $18.7 \pm 3.8695$ |
| 66 | $18.0 \pm 0.0001$ | $17.9 \pm 0.0008$ | $17.7 \pm 0.0011$ | $17.6 \pm 0.0002$ | $17.5 \pm 8.0$ |
| 68 | $21.0 \pm 0.0002$ | $20.9 \pm 0.0002$ | $20.8 \pm 0.0002$ | $20.7 \pm 0.0002$ | $20.5 \pm 4.0$ |
| 70 | $19.9 \pm 0.0001$ | $19.8 \pm 0.0001$ | $19.7 \pm 0.0$ | $19.6 \pm 0.0$ | $19.4 \pm 5.0$ |
| 71 | $20.9 \pm 0.0012$ | $20.8 \pm 0.001$ | $20.8 \pm 0.001$ | $20.6 \pm 0.0011$ | $20.4 \pm 0.0$ |
| 85 | $18.8 \pm 0.003$ | $18.7 \pm 0.0026$ | $18.6 \pm 0.0024$ | $18.4 \pm 0.0026$ | $18.2 \pm 0.0$ |
| 87 | $19.2 \pm 0.0002$ | $19.1 \pm 0.0003$ | $19.0 \pm 0.0003$ | $18.8 \pm 0.0004$ | $18.7 \pm 4.05$ |
| 88 | $19.3 \pm 0.0008$ | $19.2 \pm 0.001$ | $19.1 \pm 0.0013$ | $18.9 \pm 0.0021$ | $18.8 \pm 0.0$ |
| 93 | $21.3 \pm 0.0$ | $21.1 \pm 0.0$ | $21 \pm 0.0$ | $20.9 \pm 0.0$ | $20.7 \pm 6.2$ |
| 94 | $19.8 \pm 0.0003$ | $19.7 \pm 0.0003$ | $19.6 \pm 0.0003$ | $19.4 \pm 0.0003$ | $19.3 \pm 3.7$ |
| 95 | $21.5 \pm 0.0002$ | $21.4 \pm 0.0004$ | $21.3 \pm 0.0007$ | $21.2 \pm 0.0008$ | $21.0 \pm 2.0$ |
| 96 | $19.5 \pm 0.0035$ | $19.4 \pm 0.0044$ | $19.3 \pm 0.1181$ | $19.3 \pm 1.128$ | $19.0 \pm 0.0072$ |
| 98 | $21.5 \pm 0.0008$ | $21.4 \pm 0.0004$ | $21.3 \pm 0.0002$ | $21.2 \pm 0.0007$ | $21.1 \pm 7.6235$ |
| 102 | $20.9 \pm 0.0053$ | $20.9 \pm 0.0039$ | $20.7 \pm 0.0042$ | $20.6 \pm 0.0036$ | $20.5 \pm 0.0052$ |
| 103 | $20.2 \pm 0.0002$ | $20.1 \pm 0.0001$ | $20.0 \pm 0.0001$ | $19.9 \pm 0.0001$ | $19.8 \pm 1.0$ |
| 104 | $19.1 \pm 0.0001$ | $19.0 \pm 0.0001$ | $18.9 \pm 0.0$ | $18.8 \pm 0.0002$ | $18.7 \pm 5.37$ |
| 105 | $20.0 \pm 0.0057$ | $19.9 \pm 0.0043$ | $20.0 \pm 0.0299$ | $20.2 \pm 0.2371$ | $19.5 \pm 0.0102$ |
| 107 | $19.2 \pm 0.0001$ | $19.1 \pm 0.0001$ | $19.0 \pm 0.0001$ | $18.8 \pm 0.0001$ | $18.7 \pm 4.0$ |
| 109 | $19.5 \pm 0.0006$ | $19.4 \pm 0.0004$ | $19.3 \pm 0.0004$ | $19.1 \pm 0.0009$ | $19.0 \pm 0.0018$ |
| 113 | $19.5 \pm 0.0023$ | $19.4 \pm 0.0021$ | $19.2 \pm 0.0071$ | $18.1 \pm 0.0291$ | $18.8 \pm 0.0$ |
| 115 | $20.0 \pm 0.007$ | $19.0 \pm 0.049$ | $19.7 \pm 0.0055$ | $19.5 \pm 0.0046$ | $19.4 \pm 0.0079$ |
| 133 | $20.3 \pm 0.0175$ | $20.1 \pm 0.0674$ | $20.0 \pm 0.018$ | $19.8 \pm 0.0217$ | $19.6 \pm 0.0323$ |
| 154 | $21.9 \pm 0.0003$ | $21.8 \pm 0.0002$ | $21.7 \pm 0.0001$ | $21.6 \pm 0.0002$ | $21.4 \pm 0.0007$ |

Table C.9.: $r_{e}$ assessments in the ACS bands for the fitted early-type galaxies in MACS 0416. To calculate the values in kpe the cosmological parameters $H_{0}=70 \mathrm{~km} \mathrm{~s}^{-1} \mathrm{Mpc}^{-1}, \Omega_{\Lambda}=0.7$ and $\Omega_{M}=0.3$ have been adopted with a pixelscale of $0.065 " / \mathrm{px}$ and a cluster redshift $z \sim 0.397$

| ID | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{Z}_{850}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $2.6 \pm 0.0898$ | $2.6 \pm 0.0801$ | $2.4 \pm 0.0585$ | $2.4 \pm 0.0517$ | $2.2 \pm 0.0318$ | $2.2 \pm 0.0258$ | $2.1 \pm 0.0173$ |
| 21 | 2.6 |  |  |  |  |  |  |
| 25 | $2.8 \pm 0.8873$ | $2.7 \pm 0.7813$ | $2.4 \pm 0.5513$ | $2.3 \pm 0.4807$ | $2.0 \pm 0.2862$ | $1.9 \pm 0.2352$ | $1.8 \pm 0.1801$ |
| 26 | $7.9 \pm 0.0348$ | $7.5 \pm 0.0313$ | $6.5 \pm 0.0234$ | $6.2 \pm 0.0208$ | $5.2 \pm 0.0131$ | $4.9 \pm 0.0106$ | $4.3 \pm 0.0065$ |
| 30 | $1.8 \pm 0.0205$ | $1.8 \pm 0.0181$ | $1.7 \pm 0.013$ | $1.7 \pm 0.0114$ | $1.7 \pm 0.0068$ | $1.6 \pm 0.0055$ | $1.6 \pm 0.0037$ |
| 31 | $3.0 \pm 0.024$ | $2.8 \pm 0.0214$ | $2.5 \pm 0.0155$ | $2.4 \pm 0.0136$ | $2.1 \pm 0.0079$ | $2.0 \pm 0.0061$ | $1.8 \pm 0.0032$ |
| 33 | $1.9 \pm 0.0038$ | $1.9 \pm 0.0033$ | $1.8 \pm 0.0024$ | $1.8 \pm 0.0021$ | $1.7 \pm 0.0012$ | $1.7 \pm 0.001$ | $1.6 \pm 0.0006$ |
| 38 | $1.4 \pm 0.0062$ | $1.4 \pm 0.0056$ | $1.3 \pm 0.0042$ | $1.3 \pm 0.0038$ | $1.2 \pm 0.0025$ | $1.2 \pm 0.0021$ | $1.2 \pm 0.0014$ |
| 43 | $1.8 \pm 0.0272$ | $1.7 \pm 0.0243$ | $1.6 \pm 0.0179$ | $1.6 \pm 0.0158$ | $1.5 \pm 0.0097$ | $1.4 \pm 0.0079$ | $1.4 \pm 0.005$ |
| 45 | $3.1 \pm 0.0056$ | $3.0 \pm 0.005$ | $2.8 \pm 0.0036$ | $2.8 \pm 0.0031$ | $2.6 \pm 0.0019$ | $2.5 \pm 0.0015$ | $2.4 \pm 0.0009$ |
| 50 | $2.9 \pm 0.037$ | $2.8 \pm 0.0326$ | $2.4 \pm 0.023$ | $2.3 \pm 0.02$ | $1.9 \pm 0.0117$ | $1.8 \pm 0.0093$ | $1.6 \pm 0.0064$ |
| 53 | $2.5 \pm 0.0036$ | $2.4 \pm 0.0032$ | $2.2 \pm 0.0025$ | $2.1 \pm 0.0022$ | $1.9 \pm 0.0014$ | $1.8 \pm 0.0012$ | $1.7 \pm 0.0007$ |
| 54 | $2.2 \pm 0.0151$ | $2.1 \pm 0.0139$ | $2.0 \pm 0.0112$ | $1.9 \pm 0.0103$ | $1.8 \pm 0.0074$ | $1.7 \pm 0.0064$ | $1.6 \pm 0.0047$ |
| 55 | $6.6 \pm 0.9888$ | $6.6 \pm 0.8939$ | $6.7 \pm 0.6812$ | $6.7 \pm 0.6127$ | $6.6 \pm 0.405$ | $6.5 \pm 0.3378$ | $6.2 \pm 0.2312$ |
| 57 | $2.4 \pm 0.0112$ | $2.3 \pm 0.0099$ | $2.0 \pm 0.007$ | $2.0 \pm 0.0061$ | $1.8 \pm 0.0034$ | $1.7 \pm 0.0027$ | $1.6 \pm 0.0017$ |
| 63 | $3.4 \pm 0.0184$ | $3.2 \pm 0.0164$ | $2.9 \pm 0.012$ | $2.8 \pm 0.0106$ | $2.6 \pm 0.0065$ | $2.5 \pm 0.0051$ | $2.3 \pm 0.0032$ |
| 65 | $3.4 \pm 0.0054$ | $3.3 \pm 0.0049$ | $3.1 \pm 0.0036$ | $3.0 \pm 0.0032$ | $2.8 \pm 0.002$ | $2.7 \pm 0.0016$ | $2.6 \pm 0.0011$ |
| 66 | $4.0 \pm 0.0174$ | $4.8 \pm 0.0152$ | $6.7 \pm 0.012$ | $7.2 \pm 0.0115$ | $9.0 \pm 0.0116$ | $9.6 \pm 0.012$ | $10.6 \pm 0.0125$ |
| 68 | $2.4 \pm 0.0121$ | $2.3 \pm 0.0109$ | $2.2 \pm 0.0082$ | $2.1 \pm 0.0073$ | $2.0 \pm 0.0048$ | $2.0 \pm 0.0039$ | $1.9 \pm 0.0027$ |
| 70 | $2.2 \pm 0.0027$ | $2.1 \pm 0.0024$ | $2.0 \pm 0.0018$ | $1.9 \pm 0.0016$ | $1.7 \pm 0.001$ | $1.7 \pm 0.0008$ | $1.6 \pm 0.0005$ |
| 71 | $1.9 \pm 0.0533$ | $1.8 \pm 0.0473$ | $1.6 \pm 0.0341$ | $1.6 \pm 0.0299$ | $1.4 \pm 0.0179$ | $1.4 \pm 0.0143$ | $1.3 \pm 0.0092$ |
| 85 | $6.7 \pm 0.0458$ | $6.4 \pm 0.0406$ | $5.7 \pm 0.0291$ | $5.4 \pm 0.0256$ | $4.9 \pm 0.0154$ | $4.7 \pm 0.0125$ | $4.5 \pm 0.0087$ |


| ID |  | MACS 0416 ACS $r_{e}[\mathrm{kpc}]$ cont. |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{Z}_{850}$ |  |
| 87 | $3.9 \pm 0.0229$ | $3.7 \pm 0.0205$ | $3.4 \pm 0.0151$ | $3.3 \pm 0.0134$ | $3.0 \pm 0.0082$ | $2.9 \pm 0.0066$ | $2.8 \pm 0.0041$ |  |
| 88 | $4.4 \pm 0.044$ | $4.3 \pm 0.0396$ | $4.0 \pm 0.0298$ | $3.9 \pm 0.0267$ | $3.6 \pm 0.0172$ | $3.5 \pm 0.0142$ | $3.3 \pm 0.0095$ |  |
| 93 | $1.0 \pm 0.0003$ | $1.0 \pm 0.0003$ | $0.9 \pm 0.0002$ | $0.9 \pm 0.0002$ | $0.8 \pm 0.0001$ | $0.7 \pm 0.0001$ | $0.7 \pm 0.0001$ |  |
| 94 | $4.8 \pm 0.0173$ | $4.6 \pm 0.0154$ | $4.3 \pm 0.0114$ | $4.1 \pm 0.0101$ | $3.7 \pm 0.0061$ | $3.6 \pm 0.0048$ | $3.4 \pm 0.0028$ |  |
| 95 | $2.8 \pm 0.0117$ | $2.7 \pm 0.0104$ | $2.4 \pm 0.0074$ | $2.4 \pm 0.0064$ | $2.2 \pm 0.0036$ | $2.1 \pm 0.0027$ | $2.0 \pm 0.0014$ |  |
| 96 | $5.2 \pm 0.1297$ | $5.0 \pm 0.1157$ | $4.4 \pm 0.0844$ | $4.2 \pm 0.0744$ | $3.6 \pm 0.0445$ | $3.4 \pm 0.0351$ | $3.1 \pm 0.0211$ |  |
| 98 | $1.6 \pm 0.0263$ | $1.7 \pm 0.0218$ | $1.9 \pm 0.0127$ | $2.0 \pm 0.0104$ | $2.2 \pm 0.0073$ | $2.3 \pm 0.0079$ | $2.4 \pm 0.0093$ |  |
| 102 | $3.5 \pm 0.0899$ | $3.5 \pm 0.0804$ | $3.4 \pm 0.0592$ | $3.3 \pm 0.0525$ | $3.2 \pm 0.0325$ | $3.2 \pm 0.0263$ | $3.2 \pm 0.017$ |  |
| 103 | $4.8 \pm 0.0435$ | $4.6 \pm 0.039$ | $4.1 \pm 0.0289$ | $3.9 \pm 0.0256$ | $3.5 \pm 0.0158$ | $3.3 \pm 0.0126$ | $3.1 \pm 0.0076$ |  |
| 104 | $6.7 \pm 0.0095$ | $6.5 \pm 0.0082$ | $5.9 \pm 0.0053$ | $5.7 \pm 0.0045$ | $5.2 \pm 0.0022$ | $5.0 \pm 0.0017$ | $4.6 \pm 0.0014$ |  |
| 105 | $2.2 \pm 0.0453$ | $2.1 \pm 0.0401$ | $2.1 \pm 0.0287$ | $2.1 \pm 0.0251$ | $2.1 \pm 0.0148$ | $2.1 \pm 0.0118$ | $2.0 \pm 0.0081$ |  |
| 107 | $3.8 \pm 0.0048$ | $3.7 \pm 0.0043$ | $3.5 \pm 0.0033$ | $3.4 \pm 0.0029$ | $3.2 \pm 0.002$ | $3.1 \pm 0.0016$ | $3.0 \pm 0.0011$ |  |
| 109 | $5.0 \pm 0.0343$ | $4.9 \pm 0.0306$ | $4.4 \pm 0.0225$ | $4.3 \pm 0.0198$ | $3.9 \pm 0.0119$ | $3.7 \pm 0.0094$ | $3.5 \pm 0.0055$ |  |
| 113 | $3.8 \pm 0.0862$ | $3.7 \pm 0.0764$ | $3.3 \pm 0.0548$ | $3.2 \pm 0.048$ | $2.9 \pm 0.0278$ | $2.8 \pm 0.0217$ | $2.7 \pm 0.0132$ |  |
| 115 | $2.6 \pm 0.1185$ | $2.5 \pm 0.1052$ | $2.3 \pm 0.076$ | $2.3 \pm 0.0668$ | $2.1 \pm 0.0408$ | $2.1 \pm 0.0333$ | $2.0 \pm 0.0232$ |  |
| 133 | $1.9 \pm 0.1281$ | $1.9 \pm 0.1088$ | $1.7 \pm 0.0683$ | $1.7 \pm 0.0565$ | $1.6 \pm 0.0291$ | $1.6 \pm 0.026$ | $1.5 \pm 0.0284$ |  |
| 154 | $0.9 \pm 0.0028$ | $0.9 \pm 0.0025$ | $0.8 \pm 0.0018$ | $0.7 \pm 0.0016$ | $0.6 \pm 0.001$ | $0.6 \pm 0.0007$ | $0.5 \pm 0.0004$ |  |

Table C.10.: $r_{e}$ assessments in the WFC3/IR bands for the fitted early-type galaxies in MACS 0416. To calculate the values in kpc the cosmological parameters $H_{0}=70 \mathrm{~km} \mathrm{~s}^{-1}$ $\mathrm{Mpc}^{-1}, \Omega_{\Lambda}=0.7$ and $\Omega_{M}=0.3$ have been adopted with a pixelscale of $0.065 " / \mathrm{px}$ and a cluster redshift $z \sim 0.397$.

| ID | MACS $0416 \mathrm{WFC} 3 / \mathrm{IR} r_{e}[\mathrm{kpc}]$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $\mathrm{YJ}_{110}$ | $\mathrm{~J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| 21 | $2.0 \pm 0.013$ | $2.0 \pm 0.0116$ | $2.0 \pm 0.0108$ | $2.0 \pm 0.0098$ | $2.0 \pm 0.014$ |
| 25 | $1.6 \pm 0.1687$ | $1.5 \pm 0.1691$ | $1.5 \pm 0.1675$ | $1.5 \pm 0.1784$ | $1.5 \pm 0.2641$ |
| 26 | $3.9 \pm 0.0037$ | $3.7 \pm 0.0024$ | $3.6 \pm 0.0022$ | $3.6 \pm 0.0023$ | $3.8 \pm 0.003$ |
| 30 | $1.6 \pm 0.0029$ | $1.6 \pm 0.0027$ | $1.6 \pm 0.0025$ | $1.6 \pm 0.0026$ | $1.6 \pm 0.0045$ |
| 31 | $1.6 \pm 0.0017$ | $1.5 \pm 0.0018$ | $1.5 \pm 0.002$ | $1.4 \pm 0.0018$ | $1.4 \pm 0.0016$ |
| 33 | $1.6 \pm 0.0005$ | $1.6 \pm 0.0004$ | $1.5 \pm 0.0004$ | $1.5 \pm 0.0004$ | $1.5 \pm 0.0003$ |
| 38 | $1.2 \pm 0.0009$ | $1.1 \pm 0.0007$ | $1.1 \pm 0.0006$ | $1.1 \pm 0.0007$ | $1.1 \pm 0.001$ |
| 43 | $1.3 \pm 0.0035$ | $1.3 \pm 0.0029$ | $1.3 \pm 0.0027$ | $1.3 \pm 0.0026$ | $1.3 \pm 0.0038$ |
| 45 | $2.3 \pm 0.0006$ | $2.2 \pm 0.0006$ | $2.2 \pm 0.0005$ | $2.2 \pm 0.0004$ | $2.3 \pm 0.0007$ |
| 50 | $1.5 \pm 0.0053$ | $1.5 \pm 0.0051$ | $1.4 \pm 0.0051$ | $1.5 \pm 0.0067$ | $1.6 \pm 0.012$ |
| 53 | $1.6 \pm 0.0004$ | $1.5 \pm 0.0003$ | $1.4 \pm 0.0002$ | $1.4 \pm 0.0003$ | $1.4 \pm 0.0003$ |
| 54 | $1.5 \pm 0.0032$ | $1.5 \pm 0.0022$ | $1.4 \pm 0.0015$ | $1.4 \pm 0.0011$ | $1.4 \pm 0.0014$ |
| 55 | $5.9 \pm 0.1586$ | $5.6 \pm 0.1206$ | $5.2 \pm 0.0997$ | $4.5 \pm 0.0741$ | $3.6 \pm 0.0577$ |
| 57 | $1.6 \pm 0.0014$ | $1.5 \pm 0.0015$ | $1.6 \pm 0.0015$ | $1.6 \pm 0.0018$ | $1.8 \pm 0.0029$ |
| 63 | $2.2 \pm 0.002$ | $2.1 \pm 0.0016$ | $2.1 \pm 0.0015$ | $2.1 \pm 0.0013$ | $2.2 \pm 0.0022$ |
| 65 | $2.5 \pm 0.0007$ | $2.4 \pm 0.0006$ | $2.3 \pm 0.0005$ | $2.3 \pm 0.0005$ | $2.3 \pm 0.0003$ |
| 66 | $11.3 \pm 0.0125$ | $11.7 \pm 0.012$ | $11.9 \pm 0.0114$ | $11.8 \pm 0.012$ | $11.4 \pm 0.0164$ |
| 68 | $1.8 \pm 0.0018$ | $1.8 \pm 0.0014$ | $1.8 \pm 0.0012$ | $1.7 \pm 0.0011$ | $1.7 \pm 0.0017$ |
| 70 | $1.5 \pm 0.0003$ | $1.5 \pm 0.0002$ | $1.4 \pm 0.0002$ | $1.4 \pm 0.0002$ | $1.4 \pm 2.489$ |
| 71 | $1.2 \pm 0.0069$ | $1.2 \pm 0.0064$ | $1.1 \pm 0.0063$ | $1.1 \pm 0.0073$ | $1.2 \pm 0.0122$ |
| 85 | $4.4 \pm 0.007$ | $4.5 \pm 0.0066$ | $4.7 \pm 0.0066$ | $5.2 \pm 0.0085$ | $6.0 \pm 0.0144$ |
| 87 | $2.6 \pm 0.0027$ | $2.5 \pm 0.0022$ | $2.5 \pm 0.0021$ | $2.4 \pm 0.0019$ | $2.4 \pm 0.0026$ |
| 88 | $3.2 \pm 0.0065$ | $3.1 \pm 0.0051$ | $3.0 \pm 0.0046$ | $3.0 \pm 0.0048$ | $3.0 \pm 0.0068$ |
| 93 | $0.6 \pm 0.0$ | $0.6 \pm 0.0$ | $0.6 \pm 0.0$ | $0.6 \pm 0.0$ | $0.6 \pm 0.4806$ |
| 94 | $3.2 \pm 0.0019$ | $3.1 \pm 0.0021$ | $3.0 \pm 0.0026$ | $3.0 \pm 0.0032$ | $3.1 \pm 0.0037$ |
| 95 | $1.9 \pm 0.0011$ | $1.8 \pm 0.0014$ | $1.8 \pm 0.0016$ | $1.8 \pm 0.0016$ | $1.9 \pm 0.0014$ |
| 96 | $2.9 \pm 0.0141$ | $2.7 \pm 0.0134$ | $2.7 \pm 0.014$ | $2.6 \pm 0.0135$ | $2.8 \pm 0.0146$ |
| 98 | $2.5 \pm 0.0095$ | $2.6 \pm 0.0082$ | $2.6 \pm 0.0059$ | $2.6 \pm 0.0052$ | $2.6 \pm 0.0148$ |
| 102 | $3.1 \pm 0.0117$ | $3.1 \pm 0.0099$ | $3.1 \pm 0.0093$ | $3.1 \pm 0.009$ | $3.2 \pm 0.013$ |
| 103 | $2.9 \pm 0.0044$ | $2.8 \pm 0.0032$ | $2.7 \pm 0.003$ | $2.7 \pm 0.0027$ | $2.8 \pm 0.0037$ |
| 104 | $4.3 \pm 0.0014$ | $4.1 \pm 0.0013$ | $4.0 \pm 0.0009$ | $3.8 \pm 0.0013$ | $3.8 \pm 0.0036$ |
| 105 | $2.0 \pm 0.0072$ | $2.0 \pm 0.0075$ | $2.0 \pm 0.0077$ | $2.0 \pm 0.008$ | $2.0 \pm 0.0103$ |
| 107 | $2.9 \pm 0.0007$ | $2.8 \pm 0.0006$ | $2.8 \pm 0.0005$ | $2.7 \pm 0.0004$ | $2.7 \pm 0.0006$ |
| 109 | $3.3 \pm 0.0035$ | $3.2 \pm 0.0033$ | $3.1 \pm 0.0036$ | $3.0 \pm 0.0037$ | $3.0 \pm 0.0036$ |
|  |  |  |  |  |  |


| ID |  | MACS 0416 WFC3/IR $r_{e}[\mathrm{kpc}]$ cont. |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $\mathrm{YJ}_{110}$ | $\mathrm{~J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |  |
| 113 | $2.6 \pm 0.0103$ | $2.5 \pm 0.0108$ | $2.5 \pm 0.0115$ | $2.5 \pm 0.0112$ | $2.6 \pm 0.0132$ |  |
| 115 | $2.0 \pm 0.0186$ | $1.9 \pm 0.0166$ | $1.9 \pm 0.0148$ | $1.9 \pm 0.0134$ | $2.0 \pm 0.0229$ |  |
| 133 | $1.5 \pm 0.0319$ | $1.5 \pm 0.0325$ | $1.6 \pm 0.0323$ | $1.6 \pm 0.0415$ | $1.8 \pm 0.0726$ |  |
| 154 | $0.5 \pm 0.0003$ | $0.4 \pm 0.0002$ | $0.4 \pm 0.0002$ | $0.4 \pm 0.0002$ | $0.5 \pm 3.3667$ |  |

Table C.11.: Measurements of $n$ in the ACS bands for the selected early-type galaxies in MACS 0416.

| ID |  | MACS 0416 ACS Sérsic $n$ |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{~B}_{435}$ | $\mathrm{~g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{Z}_{850}$ |
| 21 | $5.5 \pm 0.1398$ | $5.3 \pm 0.1244$ | $4.7 \pm 0.0905$ | $4.5 \pm 0.08$ | $3.9 \pm 0.0496$ | $3.7 \pm 0.0408$ | $3.4 \pm 0.029$ |
| 25 | $6.4 \pm 1.7148$ | $6.0 \pm 1.4384$ | $5.3 \pm 0.87$ | $5.0 \pm 0.7131$ | $4.3 \pm 0.4355$ | $4.0 \pm 0.4454$ | $3.6 \pm 0.5226$ |
| 26 | $8.0 \pm 0.0153$ | $7.6 \pm 0.0139$ | $6.6 \pm 0.0106$ | $6.3 \pm 0.0095$ | $5.3 \pm 0.0063$ | $5.0 \pm 0.0053$ | $4.5 \pm 0.0036$ |
| 30 | $3.2 \pm 0.0335$ | $3.1 \pm 0.0291$ | $2.9 \pm 0.0198$ | $2.8 \pm 0.017$ | $2.6 \pm 0.0096$ | $2.6 \pm 0.008$ | $2.5 \pm 0.0067$ |
| 31 | $8.0 \pm 0.0447$ | $7.6 \pm 0.0397$ | $6.5 \pm 0.0287$ | $6.2 \pm 0.0252$ | $5.2 \pm 0.0149$ | $4.9 \pm 0.0117$ | $4.3 \pm 0.0072$ |
| 33 | $6.2 \pm 0.0074$ | $5.9 \pm 0.0065$ | $5.3 \pm 0.0047$ | $5.1 \pm 0.0041$ | $4.5 \pm 0.0025$ | $4.3 \pm 0.002$ | $4.1 \pm 0.0015$ |
| 38 | $4.6 \pm 0.0141$ | $4.4 \pm 0.0127$ | $3.9 \pm 0.0097$ | $3.8 \pm 0.0087$ | $3.4 \pm 0.0058$ | $3.3 \pm 0.0048$ | $3.1 \pm 0.0034$ |
| 43 | $3.4 \pm 0.06$ | $3.3 \pm 0.0523$ | $3.1 \pm 0.0359$ | $3.0 \pm 0.031$ | $2.7 \pm 0.018$ | $2.7 \pm 0.015$ | $2.6 \pm 0.0125$ |
| 45 | $4.7 \pm 0.006$ | $4.6 \pm 0.0054$ | $4.2 \pm 0.0039$ | $4.0 \pm 0.0034$ | $3.6 \pm 0.002$ | $3.5 \pm 0.0016$ | $3.3 \pm 0.0011$ |
| 50 | $8.0 \pm 0.064$ | $7.4 \pm 0.056$ | $6.2 \pm 0.0391$ | $5.9 \pm 0.034$ | $4.8 \pm 0.0203$ | $4.4 \pm 0.0168$ | $3.9 \pm 0.0131$ |
| 53 | $5.2 \pm 0.0044$ | $4.9 \pm 0.0041$ | $4.2 \pm 0.0032$ | $4.0 \pm 0.0029$ | $3.4 \pm 0.002$ | $3.2 \pm 0.0017$ | $2.9 \pm 0.0012$ |
| 54 | $7.8 \pm 0.0576$ | $7.4 \pm 0.0494$ | $6.5 \pm 0.0318$ | $6.2 \pm 0.0265$ | $5.3 \pm 0.0128$ | $5.1 \pm 0.0103$ | $4.6 \pm 0.0105$ |
| 55 | $4.3 \pm 0.362$ | $4.5 \pm 0.3289$ | $4.9 \pm 0.255$ | $5.0 \pm 0.2313$ | $5.3 \pm 0.1603$ | $5.4 \pm 0.1375$ | $5.4 \pm 0.1017$ |
| 57 | $7.4 \pm 0.0228$ | $7.1 \pm 0.02$ | $6.5 \pm 0.0139$ | $6.3 \pm 0.0121$ | $5.7 \pm 0.0068$ | $5.6 \pm 0.0054$ | $5.4 \pm 0.004$ |
| 63 | $5.6 \pm 0.0182$ | $5.3 \pm 0.0162$ | $4.6 \pm 0.0118$ | $4.4 \pm 0.0105$ | $3.8 \pm 0.0065$ | $3.6 \pm 0.0053$ | $3.3 \pm 0.0036$ |
| 65 | $4.3 \pm 0.0053$ | $4.2 \pm 0.0048$ | $4.0 \pm 0.0036$ | $3.9 \pm 0.0032$ | $3.7 \pm 0.0021$ | $3.6 \pm 0.0018$ | $3.5 \pm 0.0013$ |
| 66 | $3.2 \pm 0.0062$ | $3.7 \pm 0.0055$ | $4.8 \pm 0.0043$ | $5.1 \pm 0.004$ | $6.2 \pm 0.0036$ | $6.6 \pm 0.0035$ | $7.2 \pm 0.0035$ |
| 68 | $3.6 \pm 0.0141$ | $3.5 \pm 0.0125$ | $3.2 \pm 0.009$ | $3.2 \pm 0.0079$ | $2.9 \pm 0.0048$ | $2.8 \pm 0.004$ | $2.7 \pm 0.003$ |
| 70 | $4.7 \pm 0.0038$ | $4.4 \pm 0.0034$ | $3.9 \pm 0.0026$ | $3.8 \pm 0.0023$ | $3.3 \pm 0.0016$ | $3.1 \pm 0.0013$ | $2.8 \pm 0.0009$ |
| 71 | $4.9 \pm 0.1268$ | $4.6 \pm 0.112$ | $4.0 \pm 0.0799$ | $3.8 \pm 0.07$ | $3.2 \pm 0.0423$ | $3.0 \pm 0.0345$ | $2.7 \pm 0.0248$ |
| 85 | $7.2 \pm 0.0265$ | $7.0 \pm 0.0231$ | $6.6 \pm 0.016$ | $6.4 \pm 0.0139$ | $6.2 \pm 0.0084$ | $6.1 \pm 0.0071$ | $6.1 \pm 0.0059$ |
| 87 | $3.8 \pm 0.0202$ | $3.7 \pm 0.0181$ | $3.4 \pm 0.0134$ | $3.4 \pm 0.012$ | $3.1 \pm 0.0077$ | $3.0 \pm 0.0064$ | $2.9 \pm 0.0045$ |
| 88 | $7.0 \pm 0.0368$ | $6.8 \pm 0.0332$ | $6.4 \pm 0.0251$ | $6.2 \pm 0.0225$ | $5.8 \pm 0.0149$ | $5.6 \pm 0.0125$ | $5.4 \pm 0.0088$ |


| ID |  | $\mathrm{g}_{475}$ | $\mathrm{~V}_{606}$ | $\mathrm{r}_{625}$ | $\mathrm{i}_{775}$ | $\mathrm{I}_{814}$ | $\mathrm{z}_{850}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{~B}_{435}$ | $\mathrm{~g}_{475}$ |  |  |  |  |  |
| 93 | $6.6 \pm 0.0012$ | $6.1 \pm 0.0011$ | $5.2 \pm 0.0009$ | $4.8 \pm 0.0008$ | $3.9 \pm 0.0006$ | $3.6 \pm 0.0005$ | $3.2 \pm 0.0004$ |
| 94 | $7.7 \pm 0.0142$ | $7.5 \pm 0.0126$ | $7.0 \pm 0.009$ | $6.8 \pm 0.0078$ | $6.4 \pm 0.0045$ | $6.2 \pm 0.0035$ | $6.0 \pm 0.0024$ |
| 95 | $4.1 \pm 0.0133$ | $3.9 \pm 0.0117$ | $3.4 \pm 0.0083$ | $3.3 \pm 0.0072$ | $2.8 \pm 0.004$ | $2.6 \pm 0.003$ | $2.4 \pm 0.0018$ |
| 96 | $8.0 \pm 0.1075$ | $7.6 \pm 0.0961$ | $6.8 \pm 0.0707$ | $6.6 \pm 0.0627$ | $5.8 \pm 0.0392$ | $5.6 \pm 0.0321$ | $5.2 \pm 0.0223$ |
| 98 | $2.3 \pm 0.0379$ | $2.4 \pm 0.0315$ | $2.5 \pm 0.0188$ | $2.6 \pm 0.0155$ | $2.8 \pm 0.0108$ | $2.8 \pm 0.0113$ | $2.9 \pm 0.0127$ |
| 102 | $3.0 \pm 0.0606$ | $2.9 \pm 0.0542$ | $2.8 \pm 0.0399$ | $2.7 \pm 0.0354$ | $2.6 \pm 0.0222$ | $2.5 \pm 0.0181$ | $2.5 \pm 0.0123$ |
| 103 | $5.2 \pm 0.0341$ | $5.0 \pm 0.0306$ | $4.3 \pm 0.0229$ | $4.1 \pm 0.0204$ | $3.5 \pm 0.013$ | $3.3 \pm 0.0106$ | $2.9 \pm 0.007$ |
| 104 | $6.3 \pm 0.0062$ | $6.1 \pm 0.0053$ | $5.7 \pm 0.0033$ | $5.6 \pm 0.0028$ | $5.2 \pm 0.0016$ | $5.0 \pm 0.0015$ | $4.8 \pm 0.0016$ |
| 105 | $4.5 \pm 0.0808$ | $4.4 \pm 0.0712$ | $4.2 \pm 0.0504$ | $4.2 \pm 0.044$ | $4.1 \pm 0.0261$ | $4.1 \pm 0.0213$ | $4.1 \pm 0.0163$ |
| 107 | $3.5 \pm 0.0038$ | $3.4 \pm 0.0034$ | $3.4 \pm 0.0025$ | $3.4 \pm 0.0022$ | $3.3 \pm 0.0014$ | $3.3 \pm 0.0012$ | $3.3 \pm 0.0009$ |
| 109 | $6.3 \pm 0.0272$ | $6.1 \pm 0.0243$ | $5.8 \pm 0.0178$ | $5.6 \pm 0.0157$ | $5.2 \pm 0.0096$ | $5.1 \pm 0.0077$ | $4.9 \pm 0.005$ |
| 113 | $8.0 \pm 0.0978$ | $7.5 \pm 0.0866$ | $6.4 \pm 0.062$ | $6.0 \pm 0.0542$ | $5.0 \pm 0.0318$ | $4.7 \pm 0.0252$ | $4.3 \pm 0.0169$ |
| 115 | $6.8 \pm 0.2237$ | $6.5 \pm 0.1975$ | $5.7 \pm 0.1407$ | $5.4 \pm 0.1232$ | $4.7 \pm 0.0747$ | $4.5 \pm 0.0616$ | $4.2 \pm 0.0461$ |
| 133 | $6.2 \pm 0.3148$ | $5.9 \pm 0.2663$ | $5.2 \pm 0.1644$ | $5.0 \pm 0.1347$ | $4.3 \pm 0.0691$ | $4.1 \pm 0.064$ | $3.9 \pm 0.0744$ |
| 154 | $6.0 \pm 0.0261$ | $5.7 \pm 0.0224$ | $5.2 \pm 0.0148$ | $5.0 \pm 0.0126$ | $4.4 \pm 0.0077$ | $4.3 \pm 0.007$ | $4.0 \pm 0.0068$ |

Table C.12.: Measurements of $n$ in the WFC3/IR bands for the selected early-type galaxies in MACS 0416.

| ID | MACS 0416 WFC3/IR Sérsic $n$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{Y}_{105}$ | $Y \mathrm{~J}_{110}$ | $\mathrm{J}_{125}$ | $\mathrm{JH}_{140}$ | $\mathrm{H}_{160}$ |
| 21 | $3.2 \pm 0.0237$ | $3.1 \pm 0.0221$ | $3.0 \pm 0.021$ | $3.0 \pm 0.0202$ | $3.1 \pm 0.028$ |
| 25 | $3.3 \pm 0.5588$ | $3.2 \pm 0.5262$ | $3.1 \pm 0.4549$ | $3.0 \pm 0.479$ | $3.2 \pm 0.9391$ |
| 26 | $4.1 \pm 0.0025$ | $3.9 \pm 0.002$ | $3.8 \pm 0.0018$ | $3.8 \pm 0.0019$ | $4.0 \pm 0.0024$ |
| 30 | $2.4 \pm 0.0067$ | $2.4 \pm 0.0065$ | $2.5 \pm 0.0057$ | $2.5 \pm 0.0053$ | $2.7 \pm 0.01$ |
| 31 | $3.9 \pm 0.0054$ | $3.7 \pm 0.0054$ | $3.6 \pm 0.0057$ | $3.6 \pm 0.0056$ | $3.8 \pm 0.0067$ |
| 33 | $3.9 \pm 0.0012$ | $3.8 \pm 0.0011$ | $3.9 \pm 0.001$ | $4.0 \pm 0.0011$ | $4.4 \pm 0.0$ |
| 38 | $3.0 \pm 0.0025$ | $3.0 \pm 0.0021$ | $3.1 \pm 0.0021$ | $3.3 \pm 0.0025$ | $3.7 \pm 0.0036$ |
| 43 | $2.5 \pm 0.0122$ | $2.5 \pm 0.0116$ | $2.5 \pm 0.0101$ | $2.5 \pm 0.0086$ | $2.7 \pm 0.0164$ |
| 45 | $3.2 \pm 0.0009$ | $3.1 \pm 0.0008$ | $3.0 \pm 0.0008$ | $3.1 \pm 0.0007$ | $3.2 \pm 0.001$ |
| 50 | $3.7 \pm 0.0122$ | $3.6 \pm 0.0124$ | $3.7 \pm 0.0136$ | $4.2 \pm 0.0197$ | $5.0 \pm 0.0327$ |
| 53 | $2.6 \pm 0.0007$ | $2.5 \pm 0.0005$ | $2.5 \pm 0.0005$ | $2.6 \pm 0.0008$ | $2.8 \pm 0.001$ |
| 54 | $4.2 \pm 0.0123$ | $4.0 \pm 0.0126$ | $3.9 \pm 0.0111$ | $3.8 \pm 0.0068$ | $4.0 \pm 0.0124$ |
| 55 | $5.4 \pm 0.077$ | $5.2 \pm 0.0628$ | $4.9 \pm 0.054$ | $4.3 \pm 0.0464$ | $3.5 \pm 0.0507$ |
| 57 | $5.3 \pm 0.0038$ | $5.3 \pm 0.004$ | $5.5 \pm 0.0042$ | $5.9 \pm 0.0051$ | $6.5 \pm 0.0078$ |
| 63 | $3.0 \pm 0.0027$ | $2.9 \pm 0.0023$ | $2.8 \pm 0.0021$ | $2.8 \pm 0.0019$ | $3.0 \pm 0.0$ |
| 65 | $3.5 \pm 0.001$ | $3.5 \pm 0.0009$ | $3.5 \pm 0.0008$ | $3.5 \pm 0.0009$ | $3.6 \pm 0.0013$ |
| 66 | $7.6 \pm 0.0035$ | $7.9 \pm 0.0033$ | $8.0 \pm 0.0032$ | $8.0 \pm 0.0034$ | $7.7 \pm 0.004$ |
| 68 | $2.6 \pm 0.0025$ | $2.6 \pm 0.0023$ | $2.5 \pm 0.002$ | $2.5 \pm 0.0019$ | $2.6 \pm 0.0033$ |
| 70 | $2.6 \pm 0.0006$ | $2.5 \pm 0.0005$ | $2.5 \pm 0.0004$ | $2.5 \pm 0.0005$ | $2.6 \pm 6.5$ |
| 71 | $2.6 \pm 0.0213$ | $2.5 \pm 0.0209$ | $2.5 \pm 0.0217$ | $2.7 \pm 0.0277$ | $3.1 \pm 0.0446$ |
| 85 | $6.2 \pm 0.0056$ | $6.4 \pm 0.0054$ | $6.7 \pm 0.0051$ | $7.2 \pm 0.0058$ | $8.0 \pm 0.01$ |
| 87 | $2.8 \pm 0.0035$ | $2.7 \pm 0.0031$ | $2.7 \pm 0.0029$ | $2.8 \pm 0.0032$ | $2.9 \pm 0.004$ |
| 88 | $5.2 \pm 0.0066$ | $5.1 \pm 0.0056$ | $5.1 \pm 0.0054$ | $5.1 \pm 0.0061$ | $5.3 \pm 0.0086$ |
| 93 | $2.9 \pm 0.0003$ | $2.7 \pm 0.0002$ | $2.7 \pm 0.0002$ | $2.9 \pm 0.0002$ | $3.4 \pm 2.0$ |
| 94 | $5.9 \pm 0.0023$ | $5.8 \pm 0.0027$ | $5.8 \pm 0.003$ | $5.9 \pm 0.0034$ | $6.2 \pm 0.0041$ |
| 95 | $2.2 \pm 0.0017$ | $2.0 \pm 0.002$ | $2.0 \pm 0.0023$ | $2.0 \pm 0.0023$ | $2.1 \pm 0.0024$ |
| 96 | $4.9 \pm 0.0179$ | $4.8 \pm 0.0172$ | $4.7 \pm 0.0176$ | $4.8 \pm 0.0188$ | $5.1 \pm 0.0236$ |
| 98 | $3.1 \pm 0.0126$ | $3.2 \pm 0.0108$ | $3.2 \pm 0.0084$ | $3.4 \pm 0.0118$ | $3.5 \pm 0.0266$ |
| 102 | $2.4 \pm 0.0091$ | $2.4 \pm 0.0079$ | $2.4 \pm 0.0075$ | $2.5 \pm 0.0076$ | $2.6 \pm 0.0112$ |
| 103 | $2.6 \pm 0.0048$ | $2.5 \pm 0.0039$ | $2.3 \pm 0.0035$ | $2.3 \pm 0.0035$ | $2.3 \pm 0.0048$ |
| 104 | $4.6 \pm 0.0016$ | $4.5 \pm 0.0013$ | $4.4 \pm 0.001$ | $4.3 \pm 0.0014$ | $4.3 \pm 0.0$ |
| 105 | $4.3 \pm 0.0159$ | $4.4 \pm 0.0172$ | $4.6 \pm 0.0187$ | $5.0 \pm 0.0225$ | $5.5 \pm 0.0311$ |
| 107 | $3.3 \pm 0.0007$ | $3.4 \pm 0.0007$ | $3.4 \pm 0.0006$ | $3.5 \pm 0.0006$ | $3.6 \pm 9.3$ |
| 109 | $4.7 \pm 0.0039$ | $4.5 \pm 0.004$ | $4.4 \pm 0.0043$ | $4.3 \pm 0.0047$ | $4.3 \pm 0.0054$ |
| 113 | $4.1 \pm 0.0147$ | $4.1 \pm 0.0155$ | $4.2 \pm 0.0166$ | $4.6 \pm 0.0184$ | $5.4 \pm 0.0245$ |
| 115 | $4.0 \pm 0.0407$ | $4.0 \pm 0.0387$ | $4.1 \pm 0.0364$ | $4.5 \pm 0.0372$ | $5.1 \pm 0.0598$ |
| 133 | $3.8 \pm 0.0849$ | $3.8 \pm 0.0871$ | $3.9 \pm 0.0876$ | $4.3 \pm 0.1122$ | $4.9 \pm 0.1917$ |
| 154 | $3.8 \pm 0.0067$ | $3.6 \pm 0.006$ | $3.6 \pm 0.0046$ | $3.6 \pm 0.0045$ | $3.7 \pm 0.0108$ |

## D. Gradient results

## D. Gradient results

Table C.13.: Color gradients of the early-type galaxy sample of MACS 1206.

| ID | $\begin{aligned} & \nabla\left(g_{475}-I_{814}\right) \\ & {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{aligned}$ |  | $\begin{gathered} \nabla\left(r_{625}-Y_{105}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  | $\begin{gathered} \nabla\left(I_{814}-H_{160}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  | $\begin{gathered} \nabla\left(Y_{105}-H_{160}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | fitted | calculated | fitted | calculated | fitted | calculated | fitted | calculated |
| 1 | -0.20 | -0.07 | -0.22 | -0.14 | -0.11 | -0.10 | -0.03 | -0.02 |
| 3 | -0.18 | -0.13 | -0.21 | -0.18 | -0.29 | -0.28 | -0.18 | -0.18 |
| 4 | -0.08 | -0.09 | -0.06 | -0.07 | 0.00 | -0.02 | 0.02 | 0.01 |
| 5 | -0.13 | -0.08 | -0.13 | -0.10 | -0.08 | -0.08 | -0.03 | -0.03 |
| 6 | -0.18 | -0.14 | -0.17 | -0.15 | -0.07 | -0.07 | 0.00 | 0.00 |
| 8 | 0.01 | 0.05 | 0.00 | 0.04 | -0.05 | -0.01 | -0.04 | -0.03 |
| 10 | -0.18 | -0.13 | -0.20 | -0.17 | -0.16 | -0.16 | -0.08 | -0.07 |
| 11 | -0.24 | -0.14 | -0.25 | -0.18 | -0.11 | -0.10 | 0.01 | 0.00 |
| 12 | 0.02 | 0.01 | 0.00 | -0.01 | -0.09 | -0.08 | -0.08 | -0.07 |
| 15 | -0.18 | -0.15 | -0.22 | -0.18 | -0.38 | -0.34 | -0.23 | -0.24 |
| 17 | -0.12 | -0.11 | -0.11 | -0.11 | -0.07 | -0.08 | -0.02 | -0.03 |
| 19 | -0.19 | -0.08 | -0.22 | -0.14 | -0.18 | -0.16 | -0.06 | -0.08 |
| 20 | -0.19 | -0.15 | -0.10 | -0.09 | 0.16 | 0.13 | 0.18 | 0.15 |
| 21 | -0.20 | -0.23 | -0.17 | -0.18 | -0.12 | -0.11 | -0.05 | -0.04 |
| 22 | -0.14 | -0.09 | -0.15 | -0.11 | -0.08 | -0.06 | -0.01 | -0.01 |
| 23 | -0.15 | -0.08 | -0.18 | -0.12 | -0.17 | -0.15 | -0.08 | -0.08 |
| 24 | -0.29 | -0.25 | -0.33 | -0.28 | -0.64 | -0.61 | -0.29 | -0.47 |
| 25 | -0.12 | -0.06 | -0.16 | -0.12 | -0.20 | -0.20 | -0.13 | -0.12 |
| 26 | -0.12 | -0.05 | -0.16 | -0.12 | -0.30 | -0.28 | -0.21 | -0.21 |
| 28 | -0.19 | -0.07 | -0.23 | -0.15 | -0.13 | -0.13 | -0.03 | -0.03 |
| 29 | -0.20 | -0.17 | -0.20 | -0.19 | -0.13 | -0.14 | -0.05 | -0.05 |
| 30 | -0.26 | -0.18 | -0.26 | -0.22 | -0.12 | -0.13 | -0.02 | -0.02 |
| 31 | -0.10 | -0.14 | -0.07 | -0.10 | -0.01 | -0.04 | 0.01 | 0.00 |
| 32 | -0.17 | -0.13 | -0.16 | -0.14 | -0.09 | -0.10 | -0.02 | -0.03 |
| 33 | -0.13 | -0.06 | -0.16 | -0.12 | -0.18 | -0.17 | -0.10 | -0.10 |
| 35 | -0.18 | -0.12 | -0.21 | -0.17 | -0.25 | -0.25 | -0.15 | -0.15 |
| 37 | -0.10 | -0.12 | -0.13 | -0.13 | -0.36 | -0.31 | -0.27 | -0.24 |
| 38 | -0.28 | -0.23 | -0.31 | -0.28 | -0.31 | -0.29 | -0.11 | -0.15 |
| 39 | -0.16 | -0.07 | -0.22 | -0.16 | -0.28 | -0.27 | -0.16 | -0.17 |
| 40 | -0.24 | -0.20 | -0.26 | -0.23 | -0.19 | -0.19 | -0.08 | -0.08 |
| 41 | -0.15 | -0.11 | -0.18 | -0.15 | -0.21 | -0.20 | -0.12 | -0.12 |
| 42 | -0.09 | -0.07 | -0.10 | -0.08 | -0.13 | -0.11 | -0.07 | -0.07 |
| 43 | -0.23 | -0.17 | -0.21 | -0.18 | -0.08 | -0.08 | 0.00 | 0.00 |
| 45 | -0.13 | -0.09 | -0.14 | -0.11 | -0.11 | -0.10 | -0.05 | -0.05 |
| 46 | -0.22 | -0.13 | -0.25 | -0.19 | -0.20 | -0.19 | -0.09 | -0.09 |


| ID | $\begin{gathered} \nabla\left(g_{475}-I_{814}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  | $\begin{gathered} \nabla\left(r_{625}-Y_{105}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  | $\begin{gathered} \nabla\left(I_{814}-H_{160}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  | $\begin{gathered} \nabla\left(Y_{105}-H_{160}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | fitted | calculated | fitted | calculated | fitted | calculated | fitted | calculated |
| 47 | -0.14 | -0.12 | -0.13 | -0.12 | -0.08 | -0.07 | -0.02 | -0.02 |
| 48 | -0.26 | -0.21 | -0.29 | -0.25 | -0.23 | -0.22 | -0.06 | -0.09 |
| 50 | -0.16 | -0.14 | -0.14 | -0.13 | -0.06 | -0.06 | 0.00 | -0.01 |
| 51 | -0.18 | -0.17 | -0.18 | -0.17 | -0.10 | -0.11 | -0.03 | -0.03 |
| 52 | -0.15 | -0.14 | -0.13 | -0.13 | -0.03 | -0.04 | 0.02 | 0.01 |
| 55 | -0.20 | -0.17 | -0.18 | -0.17 | -0.06 | -0.08 | 0.00 | 0.00 |
| 56 | -0.23 | -0.12 | -0.27 | -0.19 | -0.18 | -0.17 | -0.06 | -0.06 |
| 57 | -0.13 | -0.13 | -0.11 | -0.12 | -0.08 | -0.07 | -0.03 | -0.02 |
| 58 | -0.29 | -0.30 | -0.28 | -0.27 | -0.34 | -0.31 | -0.16 | -0.19 |
| 59 | -0.17 | -0.04 | -0.20 | -0.12 | -0.18 | -0.17 | -0.10 | -0.09 |
| 61 | -0.22 | -0.21 | -0.19 | -0.17 | -0.05 | -0.03 | 0.03 | 0.04 |
| 62 | -0.29 | -0.25 | -0.30 | -0.27 | -0.25 | -0.23 | -0.09 | -0.10 |
| 63 | -0.12 | -0.16 | -0.10 | -0.15 | -0.30 | -0.23 | -0.25 | -0.17 |
| 64 | -0.11 | -0.04 | -0.13 | -0.08 | -0.11 | -0.10 | -0.05 | -0.05 |
| 65 | -0.20 | -0.15 | -0.26 | -0.22 | -0.34 | -0.33 | -0.17 | -0.20 |
| 68 | -0.27 | -0.26 | -0.25 | -0.25 | -0.11 | -0.11 | -0.01 | -0.01 |
| 69 | -0.25 | -0.20 | -0.25 | -0.22 | -0.07 | -0.07 | 0.03 | 0.03 |
| 70 | -0.16 | -0.15 | -0.10 | -0.11 | 0.05 | 0.02 | 0.08 | 0.06 |
| 71 | -0.17 | -0.16 | -0.18 | -0.18 | -0.16 | -0.17 | -0.08 | -0.08 |
| 72 | -0.14 | -0.10 | -0.17 | -0.15 | -0.22 | -0.22 | -0.13 | -0.13 |
| 73 | -0.18 | -0.21 | -0.16 | -0.19 | -0.13 | -0.17 | -0.06 | -0.09 |
| 74 | -0.21 | 0.05 | -0.34 | -0.18 | -0.36 | -0.35 | -0.32 | -0.18 |
| 76 | -0.14 | -0.06 | -0.17 | -0.12 | -0.13 | -0.13 | -0.06 | -0.06 |
| 77 | 0.05 | 0.06 | 0.05 | 0.06 | 0.02 | 0.03 | 0.00 | 0.01 |
| 80 | -0.16 | -0.11 | -0.15 | -0.13 | -0.05 | -0.06 | 0.00 | 0.00 |
| 81 | -0.19 | -0.10 | -0.17 | -0.12 | -0.03 | -0.02 | 0.04 | 0.04 |
| 82 | -0.26 | -0.15 | -0.33 | -0.26 | -0.23 | -0.23 | -0.09 | -0.07 |
| 83 | -0.12 | -0.05 | -0.11 | -0.08 | -0.03 | -0.03 | 0.01 | 0.01 |
| 84 | -0.12 | -0.11 | -0.12 | -0.11 | -0.10 | -0.09 | -0.04 | -0.04 |
| 85 | -0.25 | -0.23 | -0.25 | -0.22 | -0.21 | -0.18 | -0.07 | -0.07 |
| 87 | -0.14 | -0.13 | -0.16 | -0.14 | -0.27 | -0.24 | -0.17 | -0.16 |
| 88 | -0.30 | -0.31 | -0.25 | -0.25 | -0.14 | -0.11 | -0.01 | -0.01 |
| 89 | -0.12 | 0.01 | -0.14 | -0.05 | -0.14 | -0.12 | -0.07 | -0.08 |
| 90 | -0.15 | -0.08 | -0.18 | -0.13 | -0.22 | -0.21 | -0.13 | -0.13 |
| 91 | -0.01 | 0.09 | -0.04 | 0.04 | -0.17 | -0.15 | -0.15 | -0.15 |
| 93 | -0.16 | -0.15 | -0.18 | -0.17 | -0.26 | -0.23 | -0.15 | -0.15 |
| 94 | 0.04 | 0.16 | 0.01 | 0.11 | -0.22 | -0.19 | -0.23 | -0.22 |
| 96 | -0.12 | -0.06 | -0.15 | -0.11 | -0.22 | -0.21 | -0.15 | -0.14 |


| ID | $\begin{aligned} & \nabla\left(g_{475}-I_{814}\right) \\ & {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{aligned}$ |  | $\begin{gathered} \nabla\left(r_{625}-Y_{105}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  | $\begin{gathered} \nabla\left(I_{814}-H_{160}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  | $\begin{gathered} \nabla\left(Y_{105}-H_{160}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | fitted | calculated | fitted | calculated | fitted | calculated | fitted | calculated |
| 97 | -0.17 | -0.11 | -0.19 | -0.15 | -0.11 | -0.10 | -0.02 | -0.02 |
| 98 | -0.08 | -0.07 | -0.08 | -0.08 | -0.12 | -0.11 | -0.07 | -0.07 |
| 99 | -0.29 | -0.14 | -0.28 | -0.18 | 0.08 | 0.08 | 0.19 | 0.17 |
| 100 | -0.16 | -0.11 | -0.18 | -0.14 | -0.15 | -0.15 | -0.07 | -0.07 |
| 101 | -0.32 | -0.32 | -0.32 | -0.31 | -0.24 | -0.24 | -0.09 | -0.09 |
| 104 | -0.25 | -0.22 | -0.25 | -0.22 | -0.20 | -0.18 | -0.08 | -0.08 |

Table C.14.: Color gradients of the early-type galaxy sample of MACS 0416.

| ID | $\begin{gathered} \nabla\left(g_{475}-I_{814}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  | $\begin{gathered} \nabla\left(r_{625}-Y_{105}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  | $\begin{gathered} \nabla\left(I_{814}-H_{160}\right) \\ {\left[\mathrm{mag} \mathrm{dex}^{-1}\right]} \end{gathered}$ |  | $\begin{gathered} \nabla\left(Y_{105}-H_{160}\right) \\ {\left[\operatorname{mag~dex}^{-1}\right]} \end{gathered}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | fitted | calculated | fitted | calculated | fitted | calculated | fitted | calculated |
| 21 | 0.00 | 0.11 | 0.01 | 0.09 | 0.03 | 0.06 | 0.02 | 0.02 |
| 25 | -0.15 | 0.00 | -0.19 | -0.08 | -0.18 | -0.17 | -0.10 | -0.10 |
| 26 | -0.20 | -0.05 | -0.27 | -0.17 | -0.15 | -0.15 | -0.02 | -0.03 |
| 30 | -0.02 | 0.04 | 0.02 | 0.06 | -0.03 | -0.04 | -0.07 | -0.08 |
| 31 | -0.12 | 0.02 | -0.19 | -0.08 | -0.26 | -0.25 | -0.17 | -0.18 |
| 33 | 0.07 | 0.14 | 0.04 | 0.10 | -0.10 | -0.10 | -0.12 | -0.13 |
| 38 | 0.05 | 0.13 | 0.01 | 0.06 | -0.15 | -0.16 | -0.15 | -0.17 |
| 43 | -0.16 | -0.06 | -0.13 | -0.08 | -0.12 | -0.11 | -0.06 | -0.06 |
| 45 | -0.08 | 0.01 | -0.09 | -0.03 | -0.05 | -0.04 | -0.02 | -0.02 |
| 50 | -0.17 | 0.02 | -0.23 | -0.11 | -0.20 | -0.20 | -0.14 | -0.09 |
| 53 | -0.11 | 0.06 | -0.16 | -0.04 | -0.30 | -0.28 | -0.23 | -0.23 |
| 54 | -0.03 | 0.07 | -0.05 | 0.03 | -0.11 | -0.08 | -0.08 | -0.08 |
| 55 | -0.12 | -0.15 | -0.14 | -0.15 | -0.50 | -0.39 | -0.38 | -0.32 |
| 57 | -0.12 | -0.05 | -0.12 | -0.08 | -0.04 | -0.05 | -0.01 | -0.01 |
| 63 | -0.10 | 0.04 | -0.12 | -0.02 | -0.01 | 0.00 | 0.02 | 0.02 |
| 65 | -0.15 | -0.09 | -0.18 | -0.14 | -0.22 | -0.21 | -0.12 | -0.12 |
| 66 | 0.13 | 0.23 | 0.08 | 0.00 | 0.03 | -0.02 | 0.00 | 0.00 |
| 68 | -0.08 | 0.00 | -0.09 | -0.03 | -0.12 | -0.10 | -0.06 | -0.06 |
| 70 | -0.07 | 0.06 | -0.08 | 0.02 | -0.15 | -0.12 | -0.11 | -0.11 |
| 71 | -0.13 | 0.04 | -0.21 | -0.10 | -0.17 | -0.17 | -0.10 | -0.07 |
| 85 | -0.17 | -0.12 | -0.13 | -0.12 | 0.02 | 0.00 | 0.06 | 0.05 |
| 87 | -0.20 | -0.10 | -0.23 | -0.16 | -0.26 | -0.25 | -0.15 | -0.14 |
| 88 | -0.07 | -0.02 | -0.08 | -0.05 | -0.11 | -0.10 | -0.07 | -0.07 |
| 93 | -0.05 | 0.12 | -0.16 | -0.02 | -0.23 | -0.22 | -0.19 | -0.16 |
| 94 | -0.11 | -0.06 | -0.14 | -0.11 | -0.12 | -0.11 | -0.06 | -0.06 |
| 95 | -0.11 | 0.06 | -0.10 | 0.01 | -0.05 | -0.03 | 0.01 | 0.00 |
| 96 | -0.17 | -0.07 | -0.21 | -0.14 | -0.15 | -0.15 | -0.07 | -0.06 |
| 98 | 0.26 | 0.24 | 0.14 | 0.09 | 0.00 | -0.08 | -0.02 | -0.07 |
| 102 | -0.02 | 0.03 | -0.01 | 0.02 | -0.05 | -0.05 | -0.05 | -0.06 |
| 103 | -0.15 | 0.02 | -0.18 | -0.05 | -0.03 | -0.03 | 0.04 | 0.01 |
| 104 | -0.14 | -0.08 | -0.18 | -0.13 | -0.20 | -0.18 | -0.09 | -0.10 |
| 105 | -0.01 | 0.01 | -0.07 | -0.07 | -0.19 | -0.24 | -0.15 | -0.18 |
| 107 | -0.20 | -0.17 | -0.20 | -0.17 | -0.23 | -0.23 | -0.13 | -0.14 |
| 109 | -0.15 | -0.09 | -0.16 | -0.12 | -0.12 | -0.10 | -0.04 | -0.04 |
| 113 | -0.01 | 0.12 | -0.05 | 0.03 | -0.16 | -0.17 | -0.15 | -0.14 |
| 115 | 0.02 | 0.11 | -0.01 | 0.06 | -0.10 | -0.11 | -0.11 | -0.12 |
| 133 | 0.00 | 0.11 | 0.01 | 0.07 | 0.01 | -0.02 | 0.00 | -0.02 |
| 154 | -0.30 | -0.16 | -0.32 | -0.24 | -0.23 | -0.22 | -0.07 | -0.08 |

## E. SSP results

Table C.15.: SSP results for the MACS 1206 sample in the $r_{625}-Y_{105}$ color. The first column contains the IDs of the galaxies. The ages for the populations located at the respective radii are derived by assuming a fixed solar metallicity. The metallicities are determined by fixing the age to 12 Gyr at $z=0$.

| ID | Age $[\mathrm{Gyr}]$ |  |  |  | $\mathrm{Z}\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 1 | 10.2 | 13.0 | 12.0 | 7.1 | 0.0 | 0.2 | 0.0 | -0.4 |
| 3 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.3 | 0.3 | 0.2 |
| 4 | 8.7 | 7.5 | 7.5 | 7.5 | -0.1 | -0.4 | -0.4 | -0.4 |
| 5 | 12.0 | 12.0 | 9.8 | 7.5 | 0.0 | 0.0 | -0.1 | -0.4 |
| 6 | 13.0 | 13.0 | 12.0 | 8.7 | 0.2 | 0.2 | 0.0 | -0.1 |
| 8 | 9.8 | 13.0 | 13.0 | 12.0 | -0.1 | 0.2 | 0.2 | 0 |
| 10 | 13.0 | 13.0 | 12.0 | 8.7 | 0.2 | 0.2 | 0.0 | -0.4 |
| 11 | 10.2 | 12.0 | 8.7 | 6.6 | -0.1 | 0.0 | -0.4 | -0.4 |
| 12 | 9.8 | 8.7 | 8.7 | 9.8 | -0.1 | -0.1 | -0.1 | -0.1 |
| 15 | 13.0 | 12.0 | 8.7 | 7.5 | 0.2 | 0.0 | -0.1 | -0.4 |
| 17 | 12.0 | 10.2 | 8.7 | 7.5 | 0.0 | -0.1 | -0.1 | -0.4 |
| 19 | 7.5 | 9.8 | 7.5 | 6.6 | -0.4 | -0.1 | -0.4 | -0.4 |
| 20 | 6.6 | 6.3 | 6.3 | 6.3 | -0.4 | -0.4 | -0.4 | -0.4 |
| 21 | 10.2 | 7.1 | 7.1 | 6.6 | -0.1 | -0.4 | -0.4 | -0.4 |
| 22 | 13.0 | 13.0 | 12.0 | 8.7 | 0.2 | 0.0 | 0.0 | -0.4 |
| 23 | 10.2 | 12.0 | 8.7 | 7.1 | -0.1 | 0.0 | -0.1 | -0.4 |
| 24 | 12.0 | 9.8 | 7.5 | 6.3 | 0.0 | -0.1 | -0.4 | -0.4 |
| 25 | 12.0 | 12.0 | 10.2 | 7.5 | 0.0 | 0.0 | -0.1 | -0.4 |
| 26 | 13.0 | 13.0 | 12.0 | 8.7 | 0.0 | 0.2 | 0.0 | -0.4 |
| 28 | 13.0 | 13.0 | 13.0 | 7.5 | 0.2 | 0.2 | 0.0 | -0.4 |
| 29 | 13.0 | 13.0 | 10.2 | 8.7 | 0.2 | 0.0 | -0.1 | -0.4 |
| 30 | 13.0 | 12.0 | 8.7 | 7.1 | 0.2 | 0.0 | -0.1 | -0.4 |
| 31 | 12.0 | 8.7 | 7.5 | 8.7 | 0 | -0.4 | -0.4 | -0.4 |
| 32 | 13.0 | 13.0 | 12.0 | 8.7 | 0.2 | 0.0 | 0.0 | -0.1 |
| 33 | 12.0 | 12.0 | 10.2 | 7.5 | 0.0 | 0.0 | -0.1 | -0.4 |
| 35 | 13.0 | 13.0 | 13.0 | 9.8 | 0.3 | 0.2 | 0.2 | -0.1 |
| 37 | 12.0 | 8.7 | 7.5 | 7.5 | 0.0 | -0.4 | -0.4 | -0.4 |
| 38 | 12.0 | 8.7 | 7.1 | 6.3 | 0.0 | -0.4 | -0.4 | -0.4 |
| 39 | 12.0 | 12.0 | 8.7 | 7.1 | 0.0 | 0.0 | -0.1 | -0.4 |
| 40 | 13.0 | 13.0 | 10.2 | 7.5 | 0.3 | 0.0 | -0.1 | -0.4 |
| 41 | 13.0 | 13.0 | 12.0 | 8.7 | 0.2 | 0.2 | 0.0 | -0.4 |
| 42 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.2 | 0.2 | 0.0 |
| 43 | 13.0 | 13.0 | 10.2 | 7.5 | 0.2 | 0.0 | -0.1 | -0.4 |
| 45 | 12.0 | 12.0 | 9.8 | 7.5 | 0.0 | 0.0 | -0.1 | -0.4 |
|  |  |  |  |  |  |  |  |  |


| ID | Age $[\mathrm{Gyr}]$ |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 46 | 13.0 | 13.0 | 12.0 | 7.5 | 0.2 | 0.2 | 0.0 | -0.4 |
| 47 | 13.0 | 13.0 | 12.0 | 10.2 | 0.2 | 0.2 | 0.0 | -0.1 |
| 48 | 13.0 | 13.0 | 12.0 | 7.5 | 0.3 | 0.2 | 0.0 | -0.4 |
| 50 | 13.0 | 12.0 | 10.2 | 8.7 | 0.2 | 0.0 | -0.1 | -0.4 |
| 51 | 12.0 | 8.7 | 7.5 | 7.1 | 0.0 | -0.4 | -0.4 | -0.4 |
| 52 | 13.0 | 10.2 | 8.7 | 7.5 | 0.0 | -0.1 | -0.1 | -0.4 |
| 55 | 13.0 | 12.0 | 9.8 | 7.5 | 0.2 | 0.0 | -0.1 | -0.4 |
| 56 | 13.0 | 13.0 | 12.0 | 7.5 | 0.2 | 0.2 | 0.0 | -0.4 |
| 57 | 8.7 | 7.5 | 7.1 | 7.1 | -0.1 | -0.4 | -0.4 | -0.4 |
| 58 | 13.0 | 10.2 | 7.5 | 7.1 | 0.3 | -0.1 | -0.4 | -0.4 |
| 59 | 13.0 | 13.0 | 13.0 | 9.8 | 0.2 | 0.3 | 0.2 | -0.1 |
| 61 | 13.0 | 12.0 | 8.7 | 7.5 | 0.2 | 0.0 | -0.1 | -0.4 |
| 62 | 13.0 | 12.0 | 9.8 | 7.1 | 0.3 | 0.0 | -0.1 | -0.4 |
| 63 | 13.0 | 8.7 | 8.7 | 9.8 | 0.2 | -0.1 | -0.4 | -0.1 |
| 64 | 6.6 | 7.1 | 6.6 | 6.3 | -0.4 | -0.4 | -0.4 | -0.4 |
| 65 | 13.0 | 10.2 | 8.7 | 7.1 | 0.2 | 0.0 | -0.4 | -0.4 |
| 68 | 13.0 | 10.2 | 8.7 | 7.5 | 0.3 | -0.1 | -0.4 | -0.4 |
| 69 | 12.0 | 8.7 | 7.5 | 6.6 | 0.0 | -0.4 | -0.4 | -0.4 |
| 70 | 12.0 | 8.7 | 7.5 | 7.5 | 0.0 | -0.4 | -0.4 | -0.4 |
| 71 | 13.0 | 9.8 | 8.7 | 7.5 | 0.2 | -0.1 | -0.4 | -0.4 |
| 72 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.3 | 0.2 | 0.0 |
| 73 | 13.0 | 7.5 | 7.1 | 7.5 | 0.2 | -0.4 | -0.4 | -0.4 |
| 74 | 13.0 | 13.0 | 13.0 | 10.2 | 0.3 | 0.3 | 0.3 | -0.1 |
| 76 | 13.0 | 13.0 | 12.0 | 8.7 | 0.2 | 0.2 | 0.0 | -0.4 |
| 77 | 7.5 | 8.7 | 9.8 | 9.8 | -0.4 | -0.1 | -0.1 | -0.1 |
| 80 | 10.2 | 9.8 | 7.5 | 7.1 | -0.1 | -0.1 | -0.4 | -0.4 |
| 81 | 9.8 | 12.0 | 9.8 | 7.1 | -0.1 | 0.0 | -0.1 | -0.4 |
| 82 | 13.0 | 13.0 | 10.2 | 7.1 | 0.3 | 0.2 | 0.0 | -0.4 |
| 83 | 13.0 | 13.0 | 13.0 | 12.0 | 0.2 | 0.2 | 0.2 | 0.0 |
| 84 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.3 | 0.2 | 0.2 |
| 85 | 13.0 | 8.7 | 7.5 | 6.6 | 0.0 | -0.1 | -0.4 | -0.4 |
| 87 | 12.0 | 8.7 | 7.5 | 7.1 | 0.0 | -0.4 | -0.4 | -0.4 |
| 88 | 13.0 | 12.0 | 9.8 | 7.5 | 0.3 | 0.0 | -0.1 | -0.4 |
| 89 | 7.5 | 13.0 | 12.0 | 7.5 | -0.4 | 0.0 | 0.0 | -0.4 |
| 90 | 13.0 | 13.0 | 12.0 | 8.7 | 0.2 | 0.2 | 0.0 | -0.4 |
| 91 | 7.5 | 12.0 | 12.0 | 8.7 | -0.4 | 0.0 | 0.0 | -0.4 |
| 93 | 13.0 | 12.0 | 9.8 | 7.5 | 0.2 | 0.0 | -0.1 | -0.4 |
| 94 | 8.7 | 13.0 | 13.0 | 13.0 | -0.4 | 0.3 | 0.3 | 0.0 |
| 96 | 13.0 | 13.0 | 13.0 | 12.0 | 0.3 | 0.3 | 0.2 | 0.0 |
|  |  |  |  |  |  |  |  |  |


| ID | Age $[\mathrm{Gyr}]$ |  |  |  | $\mathrm{Z}\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 97 | 13.0 | 13.0 | 13.0 | 10.2 | 0.3 | 0.2 | 0.2 | -0.1 |
| 98 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.2 | 0.2 | 0.2 |
| 99 | 10.2 | 13.0 | 10.2 | 6.6 | -0.1 | 0.0 | -0.1 | -0.4 |
| 100 | 13.0 | 13.0 | 12.0 | 9.8 | 0.2 | 0.2 | 0.0 | -0.1 |
| 101 | 13.0 | 9.8 | 7.5 | 7.1 | 0.3 | -0.1 | -0.4 | -0.4 |
| 104 | 13.0 | 13.0 | 10.2 | 7.5 | 0.3 | 0.2 | -0.1 | -0.4 |

Table C.16.: SSP results for the MACS 1206 sample in the $Y_{105}-H_{160}$ color. The first column contains the IDs of the galaxies. The ages for the populations located at the respective radii are derived by assuming a fixed solar metallicity. The metallicities are determined by fixing the age to 12 Gyr at $z=0$.

| ID | Age $[\mathrm{Gyr}]$ |  |  |  | $\mathrm{Z}[\log (\mathrm{Z} / \mathrm{Z} \odot)]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 1 | 12.0 | 8.7 | 8.7 | 8.7 | 0.2 | 0.0 | 0.0 | 0.0 |
| 3 | 13.0 | 13.0 | 13.0 | 10.2 | 0.3 | 0.3 | 0.2 | 0.0 |
| 4 | 6.6 | 6.6 | 6.6 | 6.6 | -0.1 | -0.4 | -0.1 | -0.1 |
| 5 | 13.0 | 10.2 | 10.2 | 10.2 | 0.2 | 0.2 | 0.0 | 0.0 |
| 6 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.2 | 0.2 | 0.3 |
| 8 | 7.1 | 7.5 | 7.5 | 6.6 | -0.1 | 0.0 | 0.0 | -0.1 |
| 10 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.3 | 0.3 | 0.2 |
| 11 | 8.7 | 12.0 | 12.0 | 10.2 | 0.0 | 0.2 | 0.2 | 0.0 |
| 12 | 8.7 | 7.5 | 7.5 | 7.1 | 0.0 | 0.0 | 0.0 | -0.1 |
| 15 | 12.0 | 12.0 | 7.5 | 6.3 | 0.2 | 0.2 | 0.0 | -0.4 |
| 17 | 13.0 | 10.2 | 10.2 | 10.2 | 0.2 | 0.0 | 0.0 | 0.0 |
| 19 | 7.1 | 9.8 | 8.7 | 7.1 | 0.0 | 0.0 | 0.0 | -0.1 |
| 20 | 6.3 | 6.3 | 6.6 | 8.7 | -0.4 | -0.4 | -0.1 | 0.0 |
| 21 | 7.5 | 8.7 | 7.5 | 7.1 | 0.0 | 0.0 | 0.0 | -0.1 |
| 22 | 13.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.2 | 0.2 | 0.2 |
| 23 | 10.2 | 12.0 | 9.8 | 7.1 | 0.0 | 0.2 | 0.0 | 0.0 |
| 24 | 6.6 | 13.0 | 13.0 | 5.7 | -0.1 | 0.3 | 0.3 | -0.4 |
| 25 | 13.0 | 13.0 | 10.2 | 8.7 | 0.3 | 0.2 | 0.0 | 0.0 |
| 26 | 13.0 | 13.0 | 12.0 | 7.5 | 0.3 | 0.3 | 0.2 | 0.0 |
| 28 | 13.0 | 13.0 | 13.0 | 12.0 | 0.2 | 0.2 | 0.2 | 0.2 |
| 29 | 13.0 | 13.0 | 12.0 | 13.0 | 0.3 | 0.2 | 0.2 | 0.2 |
| 30 | 13.0 | 10.2 | 10.2 | 13.0 | 0.3 | 0.2 | 0.0 | 0.2 |
| 31 | 13.0 | 12.0 | 12.0 | 13.0 | 0.2 | 0.2 | 0.2 | 0.2 |
| 32 | 13.0 | 12.0 | 12.0 | 13.0 | 0.3 | 0.2 | 0.2 | 0.2 |
| 33 | 13.0 | 7.5 | 7.1 | 7.1 | 0.2 | 0.0 | 0.0 | -0.1 |
| 35 | 13.0 | 12.0 | 8.7 | 7.5 | 0.3 | 0.2 | 0.0 | 0.0 |
| 37 | 10.2 | 8.7 | 6.6 | 5.7 | 0.2 | 0.0 | -0.1 | -0.4 |
| 38 | 7.5 | 12.0 | 8.7 | 6.6 | 0.0 | 0.2 | 0.0 | -0.1 |
| 39 | 13.0 | 13.0 | 12.0 | 7.5 | 0.3 | 0.3 | 0.2 | 0.0 |
| 40 | 13.0 | 10.2 | 9.8 | 8.7 | 0.3 | 0.2 | 0.0 | 0.0 |
| 41 | 13.0 | 13.0 | 10.2 | 8.7 | 0.3 | 0.2 | 0.0 | 0.0 |
| 42 | 13.0 | 13.0 | 12.0 | 10.2 | 0.2 | 0.2 | 0.2 | 0.0 |
| 43 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.2 | 0.2 | 0.3 |
| 45 | 12.0 | 9.8 | 8.7 | 8.7 | 0.2 | 0.0 | 0.0 | 0.0 |
| 46 | 13.0 | 10.2 | 8.7 | 7.5 | 0.3 | 0.0 | 0.0 | 0.0 |
| 47 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.3 | 0.3 | 0.3 |
|  |  |  |  |  |  |  |  |  |


| ID | Age $[\mathrm{Gyr}]$ |  |  |  | $\mathrm{Z}[\log (\mathrm{Z} / \mathrm{Z} \odot)]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 48 | 12.0 | 13.0 | 13.0 | 10.2 | 0.2 | 0.3 | 0.2 | 0.2 |
| 50 | 12.0 | 10.2 | 10.2 | 10.2 | 0.2 | 0.0 | 0.0 | 0.2 |
| 51 | 13.0 | 10.2 | 10.2 | 10.2 | 0.2 | 0.2 | 0.0 | 0.0 |
| 52 | 13.0 | 12.0 | 13.0 | 13.0 | 0.2 | 0.2 | 0.2 | 0.2 |
| 55 | 13.0 | 10.2 | 10.2 | 13.0 | 0.3 | 0.2 | 0.2 | 0.2 |
| 56 | 12.0 | 12.0 | 10.2 | 8.7 | 0.2 | 0.2 | 0.0 | 0.0 |
| 57 | 9.8 | 10.2 | 10.2 | 8.7 | 0.0 | 0.2 | 0.0 | 0.0 |
| 58 | 12.0 | 13.0 | 13.0 | 7.1 | 0.2 | 0.3 | 0.2 | -0.1 |
| 59 | 13.0 | 9.8 | 7.5 | 7.5 | 0.3 | 0.0 | 0.0 | 0.0 |
| 61 | 7.5 | 12.0 | 12.0 | 12.0 | 0.0 | 0.2 | 0.2 | 0.2 |
| 62 | 13.0 | 12.0 | 10.2 | 7.5 | 0.2 | 0.2 | 0.0 | 0.0 |
| 63 | 7.5 | 7.5 | 6.6 | 5.4 | 0.0 | 0.0 | -0.1 | -0.4 |
| 64 | 7.1 | 6.6 | 6.6 | 6.6 | -0.1 | -0.1 | -0.1 | -0.1 |
| 65 | 9.8 | 8.7 | 7.1 | 6.3 | 0.0 | 0.0 | -0.1 | -0.4 |
| 68 | 13.0 | 10.2 | 10.2 | 12.0 | 0.2 | 0.2 | 0.2 | 0.2 |
| 69 | 8.7 | 7.5 | 8.7 | 10.2 | 0.0 | 0.0 | 0.0 | 0.0 |
| 70 | 7.5 | 7.1 | 7.5 | 12.0 | 0.0 | -0.1 | 0.0 | 0.2 |
| 71 | 13.0 | 8.7 | 7.5 | 7.1 | 0.2 | 0.0 | 0.0 | 0.0 |
| 72 | 13.0 | 8.7 | 7.5 | 7.1 | 0.3 | 0.0 | 0.0 | -0.1 |
| 73 | 13.0 | 7.1 | 7.1 | 7.5 | 0.2 | 0.0 | -0.1 | 0.0 |
| 74 | 13.0 | 6.3 | 5.7 | 5.7 | 0.3 | -0.4 | -0.4 | -0.4 |
| 76 | 13.0 | 12.0 | 10.2 | 9.8 | 0.3 | 0.2 | 0.0 | 0.0 |
| 77 | 7.1 | 7.5 | 7.5 | 7.5 | -0.1 | 0.0 | 0.0 | 0.0 |
| 80 | 7.5 | 7.1 | 7.1 | 7.1 | 0.0 | -0.1 | -0.1 | -0.1 |
| 81 | 7.5 | 7.5 | 7.5 | 9.8 | 0.0 | 0.0 | 0.0 | 0.0 |
| 82 | 13.0 | 7.5 | 7.1 | 7.1 | 0.2 | 0.0 | 0.0 | -0.1 |
| 83 | 13.0 | 12.0 | 12.0 | 13.0 | 0.3 | 0.2 | 0.2 | 0.2 |
| 84 | 13.0 | 13.0 | 12.0 | 12.0 | 0.2 | 0.2 | 0.2 | 0.2 |
| 85 | 7.1 | 10.2 | 8.7 | 7.1 | -0.1 | 0.0 | 0.0 | -0.1 |
| 87 | 13.0 | 12.0 | 8.7 | 6.6 | 0.2 | 0.2 | 0.0 | -0.1 |
| 88 | 8.7 | 13.0 | 13.0 | 10.2 | 0.0 | 0.2 | 0.2 | 0.0 |
| 89 | 12.0 | 10.2 | 8.7 | 7.5 | 0.2 | 0.0 | 0.0 | 0.0 |
| 90 | 13.0 | 10.2 | 7.5 | 7.1 | 0.2 | 0.0 | 0.0 | -0.1 |
| 91 | 13.0 | 8.7 | 7.1 | 6.6 | 0.3 | 0.0 | 0.0 | -0.1 |
| 93 | 13.0 | 13.0 | 10.2 | 7.1 | 0.3 | 0.2 | 0.2 | -0.1 |
| 94 | 13.0 | 13.0 | 12.0 | 8.7 | 0.3 | 0.3 | 0.2 | 0.0 |
| 96 | 13.0 | 13.0 | 13.0 | 12.0 | 0.3 | 0.3 | 0.2 | 0.2 |
| 97 | 13.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.2 | 0.2 | 0.2 |
| 98 | 13.0 | 13.0 | 12.0 | 9.8 | 0.2 | 0.2 | 0.2 | 0.0 |


| ID | Age [Gyr] |  |  |  | $\mathrm{Z}[\log (\mathrm{Z} / \mathrm{Z} \odot)]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 99 | 6.6 | 12.0 | 13.0 | 13.0 | -0.4 | 0.2 | 0.2 | 0.3 |
| 100 | 13.0 | 12.0 | 10.2 | 8.7 | 0.2 | 0.2 | 0.0 | 0.0 |
| 101 | 12.0 | 7.5 | 7.5 | 7.1 | 0.2 | 0.0 | 0.0 | -0.1 |
| 104 | 9.8 | 10.2 | 8.7 | 7.1 | 0.0 | 0.2 | 0.0 | -0.1 |

Table C.17.: SSP results for the MACS 0416 sample in the $r_{625}-Y_{105}$ color. The first column contains the IDs of the galaxies. The ages for the populations located at the respective radii are derived by assuming a fixed solar metallicity. The metallicities are determined by fixing the age to 12 Gyr at $z=0$.

| ID | Age $[\mathrm{Gyr}]$ |  |  |  | $\mathrm{Z}\left[\log \left(\mathrm{Z} / \mathrm{Z}{ }_{\odot}\right)\right]$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |  |
| 21 | 6.3 | 8.7 | 8.7 | 7.1 | -0.4 | -0.4 | -0.4 | -0.4 |  |
| 25 | 8.7 | 13.0 | 12.0 | 7.1 | -0.1 | 0.2 | 0.0 | -0.4 |  |
| 26 | 12.0 | 13.0 | 12.0 | 7.1 | 0.0 | 0.2 | 0.0 | -0.4 |  |
| 30 | 6.3 | 7.1 | 7.1 | 6.6 | -0.4 | -0.4 | -0.4 | -0.4 |  |
| 31 | 10.2 | 13.0 | 13.0 | 7.5 | -0.1 | 0.2 | 0.0 | -0.4 |  |
| 33 | 7.5 | 13.0 | 13.0 | 12.0 | -0.4 | 0.0 | 0.2 | 0.0 |  |
| 38 | 7.1 | 9.8 | 9.8 | 7.5 | -0.4 | -0.1 | -0.1 | -0.4 |  |
| 43 | 7.5 | 10.2 | 8.7 | 7.1 | -0.4 | -0.1 | -0.4 | -0.4 |  |
| 45 | 7.5 | 10.2 | 8.7 | 7.1 | -0.4 | -0.1 | -0.1 | -0.4 |  |
| 50 | 10.2 | 13.0 | 13.0 | 7.5 | -0.1 | 0.2 | 0.0 | -0.4 |  |
| 53 | 9.8 | 13.0 | 13.0 | 8.7 | -0.1 | 0.3 | 0.2 | -0.4 |  |
| 54 | 7.5 | 13.0 | 12.0 | 8.7 | -0.4 | 0.0 | 0.0 | -0.4 |  |
| 55 | 6.6 | 6.3 | 5.7 | 5.7 | -0.4 | -0.4 | -0.4 | -0.4 |  |
| 57 | 12.0 | 13.0 | 12.0 | 8.7 | 0.0 | 0.0 | 0.0 | -0.1 |  |
| 63 | 10.2 | 13.0 | 13.0 | 9.8 | -0.1 | 0.3 | 0.2 | -0.1 |  |
| 65 | 13.0 | 12.0 | 10.2 | 7.5 | 0.0 | 0.0 | -0.1 | -0.4 |  |
| 66 | 10.2 | 7.5 | 7.5 | 10.2 | -0.1 | -0.4 | -0.4 | -0.1 |  |
| 68 | 7.5 | 10.2 | 8.7 | 7.1 | -0.4 | -0.1 | -0.1 | -0.4 |  |
| 70 | 8.7 | 13.0 | 13.0 | 9.8 | -0.4 | 0.3 | 0.2 | -0.1 |  |
| 71 | 8.7 | 13.0 | 12.0 | 7.1 | -0.4 | 0.2 | 0.0 | -0.4 |  |
| 85 | 13.0 | 10.2 | 8.7 | 7.5 | 0.0 | -0.1 | -0.1 | -0.4 |  |
| 87 | 10.2 | 12.0 | 8.7 | 6.6 | -0.1 | 0.0 | -0.4 | -0.4 |  |
| 88 | 12.0 | 13.0 | 12.0 | 9.8 | 0.0 | 0.0 | 0.0 | -0.1 |  |
| 93 | 7.1 | 12.0 | 10.2 | 6.6 | -0.4 | 0.0 | -0.1 | -0.4 |  |
| 94 | 9.8 | 8.7 | 7.5 | 7.1 | -0.1 | -0.1 | -0.4 | -0.4 |  |
| 95 | 6.3 | 10.2 | 8.7 | 6.6 | -0.4 | -0.1 | -0.1 | -0.4 |  |
| 96 | 8.7 | 8.7 | 7.5 | 6.6 | -0.4 | -0.1 | -0.4 | -0.4 |  |
| 98 | 7.1 | 6.6 | 7.1 | 8.7 | -0.4 | -0.4 | -0.4 | -0.4 |  |
| 102 | 8.7 | 10.2 | 10.2 | 8.7 | -0.4 | -0.1 | -0.1 | -0.4 |  |
| 103 | 7.1 | 12.0 | 9.8 | 6.6 | -0.4 | 0.0 | -0.1 | -0.4 |  |
| 104 | 10.2 | 10.2 | 8.7 | 7.1 | -0.1 | -0.1 | -0.4 | -0.4 |  |
| 105 | 13.0 | 12.0 | 12.0 | 10.2 | 0.2 | 0.0 | 0.0 | -0.1 |  |
| 107 | 13.0 | 13.0 | 13.0 | 9.8 | 0.3 | 0.2 | 0.0 | -0.1 |  |
| 109 | 9.8 | 9.8 | 8.7 | 7.1 | -0.1 | -0.1 | -0.4 | -0.4 |  |
| 113 | 7.5 | 13.0 | 12.0 | 8.7 | -0.4 | 0.0 | 0.0 | -0.4 |  |
|  |  |  |  |  |  |  |  |  |  |


| ID | Age [Gyr] |  |  |  | $\mathrm{Z}[\log (\mathrm{Z} / \mathrm{Z} \odot)]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 115 | 6.6 | 7.5 | 7.5 | 7.1 | -0.4 | -0.4 | -0.4 | -0.4 |
| 133 | 8.7 | 13.0 | 13.0 | 10.2 | -0.4 | 0.2 | 0.2 | 0.0 |
| 154 | 13.0 | 13.0 | 9.8 | 6.6 | 0.0 | 0.0 | -0.1 | -0.4 |

Table C.18.: SSP results for the MACS 0416 sample in the $Y_{105}-H_{160}$ color. The first column contains the IDs of the galaxies. The ages for the populations located at the respective radii are derived by assuming a fixed solar metallicity. The metallicities are determined by fixing the age to 12 Gyr at $z=0$.

| ID | Age $[\mathrm{Gyr}]$ |  |  |  | $\mathrm{Z}\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 21 | 7.5 | 9.8 | 10.2 | 10.2 | 0.0 | 0.0 | 0.0 | 0.0 |
| 25 | 7.5 | 7.5 | 7.1 | 6.6 | 0.0 | 0.0 | -0.1 | -0.1 |
| 26 | 10.2 | 10.2 | 10.2 | 8.7 | 0.2 | 0.2 | 0.0 | 0.0 |
| 30 | 6.6 | 5.7 | 5.7 | 5.7 | -0.1 | -0.4 | -0.4 | -0.4 |
| 31 | 12.0 | 7.5 | 7.1 | 6.3 | 0.2 | 0.0 | -0.1 | -0.4 |
| 33 | 13.0 | 7.5 | 7.1 | 7.1 | 0.3 | 0.0 | -0.1 | -0.1 |
| 38 | 13.0 | 6.3 | 5.7 | 6.3 | 0.2 | -0.4 | -0.4 | -0.4 |
| 43 | 7.5 | 6.6 | 6.6 | 6.6 | 0.0 | -0.1 | -0.1 | -0.1 |
| 45 | 13.0 | 13.0 | 13.0 | 13.0 .0 | 0.3 | 0.3 | 0.3 | 0.2 |
| 50 | 13.0 | 7.5 | 7.1 | 7.5 | 0.3 | 0.0 | -0.1 | 0.0 |
| 53 | 13.0 | 12.0 | 7.5 | 7.1 | 0.3 | 0.2 | 0.0 | -0.1 |
| 54 | 7.5 | 7.5 | 7.1 | 6.6 | 0.0 | 0.0 | -0.1 | -0.1 |
| 55 | 6.6 | 7.1 | 5.7 | 5.0 | -0.1 | -0.1 | -0.4 | -0.4 |
| 57 | 13.0 | 8.7 | 8.7 | 10.2 | 0.2 | 0.0 | 0.0 | 0.2 |
| 63 | 8.7 | 8.7 | 8.7 | 9.8 | 0.0 | 0.0 | 0.0 | 0.0 |
| 65 | 13.0 | 12.0 | 8.7 | 7.5 | 0.3 | 0.2 | 0.0 | 0.0 |
| 66 | 8.7 | 8.7 | 8.7 | 8.7 | 0.0 | 0.0 | 0.0 | 0.0 |
| 68 | 12.0 | 10.2 | 9.8 | 7.5 | 0.2 | 0.2 | 0.0 | 0.0 |
| 70 | 13.0 | 12.0 | 9.8 | 7.5 | 0.2 | 0.2 | 0.0 | 0.0 |
| 71 | 13.0 | 7.1 | 7.1 | 7.1 | 0.3 | -0.1 | -0.1 | -0.1 |
| 85 | 8.7 | 7.1 | 7.1 | 10.2 | 0.0 | -0.1 | 0.0 | 0.0 |
| 87 | 13.0 | 12.0 | 8.7 | 7.1 | 0.3 | 0.2 | 0.0 | 0.0 |
| 88 | 13.0 | 10.2 | 8.7 | 7.5 | 0.2 | 0.0 | 0.0 | 0.0 |
| 93 | 13.0 | 13.0 | 10.2 | 9.8 | 0.3 | 0.2 | 0.2 | 0.0 |
| 94 | 12.0 | 8.7 | 7.5 | 7.5 | 0.2 | 0.0 | 0.0 | 0.0 |
| 95 | 8.7 | 10.2 | 10.2 | 9.8 | 0.0 | 0.2 | 0.2 | 0.0 |
| 96 | 13.0 | 8.7 | 8.7 | 7.5 | 0.2 | 0.0 | 0.0 | 0.0 |
| 98 | 7.1 | 6.3 | 6.3 | 6.6 | -0.1 | -0.4 | -0.4 | -0.1 |
| 102 | 7.5 | 6.6 | 6.6 | 6.6 | 0.0 | -0.1 | -0.4 | -0.1 |
| 103 | 6.3 | 7.5 | 7.5 | 7.1 | -0.4 | 0.0 | 0.0 | -0.1 |
| 104 | 8.7 | 7.5 | 7.1 | 6.6 | 0.0 | 0.0 | -0.1 | -0.1 |
| 105 | 13.0 | 7.1 | 6.6 | 6.6 | 0.3 | -0.1 | -0.1 | -0.1 |
| 107 | 13.0 | 8.7 | 7.5 | 7.1 | 0.3 | 0.0 | 0.0 | -0.1 |
| 109 | 12.0 | 13.0 | 12.0 | 10.2 | 0.2 | 0.2 | 0.2 | 0.0 |
| 113 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.2 | 0.2 | 0.2 |
|  |  |  |  |  |  |  |  |  |


| ID | Age [Gyr] |  |  |  | $\mathrm{Z}\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 115 | 13.0 | 8.7 | 7.5 | 8.7 | 0.3 | 0.0 | 0.0 | 0.0 |
| 133 | 13.0 | 12.0 | 12.0 | 13.0 | 0.3 | 0.2 | 0.2 | 0.3 |
| 154 | 13.0 | 12.0 | 10.2 | 8.7 | 0.2 | 0.2 | 0.0 | 0.0 |

Table C.19.: SSP results for the MACS 1206 sample using a color-color grid with $r_{625}-Y_{105}$ and $Y_{105}-H_{160}$. The first column contains the IDs of the galaxies. The ages and metallicities for the stellar populations at the respective radii are derived by finding the nearest grid points with their colors corresponding to modelled ages and metallicities.

| ID | Age $[\mathrm{Gyr}]$ |  |  |  |  |  |  |  |  | $\mathrm{Z}[\log (\mathrm{Z} / \mathrm{Z} \odot)]$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |  |  |  |  |  |  |
| 1 | 12.0 | 12.0 | 13.0 | 13.0 | 0.0 | 0.0 | -0.1 | -0.1 |  |  |  |  |  |  |
| 3 | 13.0 | 13.0 | 13.0 | 13.0 | 0.4 | 0.3 | 0.2 | 0.2 |  |  |  |  |  |  |
| 4 | 13.0 | 10.2 | 10.2 | 10.2 | -0.4 | -0.4 | -0.4 | -0.4 |  |  |  |  |  |  |
| 5 | 8.7 | 12.0 | 9.8 | 9.8 | 0.2 | 0.0 | 0.0 | 0.0 |  |  |  |  |  |  |
| 6 | 13.0 | 10.2 | 9.8 | 9.8 | 0.2 | 0.2 | 0.2 | 0.2 |  |  |  |  |  |  |
| 8 | 13.0 | 13.0 | 13.0 | 13.0 | -0.4 | -0.1 | -0.1 | -0.1 |  |  |  |  |  |  |
| 10 | 10.2 | 12.0 | 9.8 | 9.8 | 0.3 | 0.2 | 0.2 | 0.2 |  |  |  |  |  |  |
| 11 | 13.0 | 12.0 | 7.5 | 7.5 | -0.1 | 0.0 | 0.2 | 0.2 |  |  |  |  |  |  |
| 12 | 9.8 | 8.7 | 10.2 | 10.2 | 0.0 | 0.0 | -0.1 | -0.1 |  |  |  |  |  |  |
| 15 | 13.0 | 12.0 | 12.0 | 12.0 | 0.0 | 0.0 | -0.1 | -0.1 |  |  |  |  |  |  |
| 17 | 12.0 | 10.2 | 8.7 | 8.7 | 0.0 | 0.0 | 0.0 | 0.0 |  |  |  |  |  |  |
| 19 | 9.8 | 9.8 | 7.5 | 7.5 | -0.1 | 0.0 | 0.0 | 0.0 |  |  |  |  |  |  |
| 20 | 7.5 | 7.1 | 6.3 | 6.3 | -0.4 | -0.4 | 0.0 | 0.0 |  |  |  |  |  |  |
| 21 | 12.0 | 7.1 | 6.6 | 6.6 | -0.1 | 0.2 | 0.2 | 0.2 |  |  |  |  |  |  |
| 22 | 13.0 | 10.2 | 8.7 | 8.7 | 0.0 | 0.2 | 0.2 | 0.2 |  |  |  |  |  |  |
| 23 | 10.2 | 12.0 | 8.7 | 8.7 | 0.0 | 0.0 | 0.0 | 0.0 |  |  |  |  |  |  |
| 24 | 13.0 | 7.1 | 6.3 | 6.3 | -0.1 | 0.4 | 0.4 | 0.4 |  |  |  |  |  |  |
| 25 | 7.5 | 12.0 | 10.2 | 10.2 | 0.4 | 0.0 | 0.0 | 0.0 |  |  |  |  |  |  |
| 26 | 7.5 | 10.2 | 12.0 | 12.0 | 0.4 | 0.2 | 0.0 | 0.0 |  |  |  |  |  |  |
| 28 | 12.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.2 | 0.0 | 0.0 |  |  |  |  |  |  |
| 29 | 9.8 | 13.0 | 10.2 | 10.2 | 0.3 | 0.0 | 0.0 | 0.0 |  |  |  |  |  |  |
| 30 | 8.7 | 12.0 | 7.5 | 7.5 | 0.3 | 0.0 | 0.2 | 0.2 |  |  |  |  |  |  |
| 31 | 9.8 | 7.5 | 7.5 | 7.5 | 0.2 | 0.2 | 0.2 | 0.2 |  |  |  |  |  |  |
| 32 | 13.0 | 13.0 | 8.7 | 8.7 | 0.2 | 0.0 | 0.2 | 0.2 |  |  |  |  |  |  |
| 33 | 9.8 | 13.0 | 12.0 | 12.0 | 0.2 | -0.1 | -0.1 | -0.1 |  |  |  |  |  |  |
| 35 | 12.0 | 12.0 | 12.0 | 12.0 | 0.3 | 0.2 | 0.0 | 0.0 |  |  |  |  |  |  |
| 37 | 12.0 | 8.7 | 12.0 | 12.0 | 0.0 | 0.0 | -0.4 | -0.4 |  |  |  |  |  |  |
| 38 | 13.0 | 7.5 | 6.3 | 6.3 | -0.1 | 0.2 | 0.3 | 0.3 |  |  |  |  |  |  |
| 39 | 7.1 | 7.5 | 7.5 | 7.5 | 0.4 | 0.3 | 0.2 | 0.2 |  |  |  |  |  |  |
| 40 | 13.0 | 12.0 | 10.2 | 10.2 | 0.2 | 0.0 | 0.0 | 0.0 |  |  |  |  |  |  |
| 41 | 9.8 | 13.0 | 12.0 | 12.0 | 0.3 | 0.0 | 0.0 | 0.0 |  |  |  |  |  |  |
| 42 | 13.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.2 | 0.0 | 0.0 |  |  |  |  |  |  |
| 43 | 13.0 | 13.0 | 8.7 | 8.7 | 0.2 | 0.0 | 0.2 | 0.2 |  |  |  |  |  |  |
| 45 | 12.0 | 13.0 | 9.8 | 9.8 | 0.0 | -0.1 | 0.0 | 0.0 |  |  |  |  |  |  |
| 46 | 12.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.0 | -0.1 | -0.1 |  |  |  |  |  |  |
| 47 | 13.0 | 10.2 | 9.8 | 9.8 | 0.2 | 0.2 | 0.2 | 0.2 |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |


| ID | Age $[\mathrm{Gyr}]$ |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 48 | 13.0 | 12.0 | 8.7 | 8.7 | 0.2 | 0.2 | 0.2 | 0.2 |
| 50 | 13.0 | 12.0 | 10.2 | 10.2 | 0.0 | 0.0 | 0.0 | 0.0 |
| 51 | 8.7 | 7.5 | 7.1 | 7.1 | 0.2 | 0.2 | 0.2 | 0.2 |
| 52 | 13.0 | 10.2 | 7.5 | 7.5 | 0.0 | 0.0 | 0.2 | 0.2 |
| 55 | 12.0 | 12.0 | 10.2 | 10.2 | 0.2 | 0.0 | 0.0 | 0.0 |
| 56 | 13.0 | 13.0 | 12.0 | 12.0 | 0.0 | 0.0 | 0.0 | 0.0 |
| 57 | 9.8 | 7.1 | 6.6 | 6.6 | 0.0 | 0.2 | 0.3 | 0.3 |
| 58 | 13.0 | 7.5 | 7.1 | 7.1 | 0.2 | 0.3 | 0.3 | 0.3 |
| 59 | 13.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.2 | 0.0 | 0.0 |
| 61 | 13.0 | 12.0 | 7.5 | 7.5 | 0.0 | 0.0 | 0.2 | 0.2 |
| 62 | 13.0 | 12.0 | 9.8 | 9.8 | 0.2 | 0.0 | 0.0 | 0.0 |
| 63 | 13.0 | 12.0 | 13.0 | 13.0 | 0.0 | -0.1 | -0.4 | -0.4 |
| 64 | 6.6 | 7.5 | 7.1 | 7.1 | 0.2 | -0.1 | -0.1 | -0.1 |
| 65 | 13.0 | 13.0 | 10.2 | 10.2 | 0.0 | -0.1 | -0.1 | -0.1 |
| 68 | 13.0 | 10.2 | 7.5 | 7.5 | 0.2 | 0.0 | 0.2 | 0.2 |
| 69 | 13.0 | 8.7 | 7.1 | 7.1 | -0.1 | 0.0 | 0.2 | 0.2 |
| 70 | 13.0 | 10.2 | 7.5 | 7.5 | -0.1 | -0.1 | 0.0 | 0.0 |
| 71 | 13.0 | 12.0 | 10.2 | 10.2 | 0.0 | -0.1 | -0.1 | -0.1 |
| 72 | 13.0 | 13.0 | 13.0 | 13.0 | 0.3 | 0.2 | 0.0 | 0.0 |
| 73 | 13.0 | 7.5 | 8.7 | 8.7 | 0.0 | 0.0 | -0.1 | -0.1 |
| 74 | 9.8 | 13.0 | 13.0 | 13.0 | 0.4 | 0.2 | -0.1 | -0.1 |
| 76 | 10.2 | 13.0 | 12.0 | 12.0 | 0.2 | 0.0 | 0.0 | 0.0 |
| 77 | 9.8 | 12.0 | 12.0 | 12.0 | -0.1 | -0.1 | -0.1 | -0.1 |
| 80 | 12.0 | 13.0 | 9.8 | 9.8 | -0.1 | -0.4 | -0.1 | -0.1 |
| 81 | 12.0 | 13.0 | 12.0 | 12.0 | -0.1 | -0.1 | -0.1 | -0.1 |
| 82 | 13.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.0 | -0.1 | -0.1 |
| 83 | 13.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.2 | 0.0 | 0.0 |
| 84 | 13.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.2 | 0.0 | 0.0 |
| 85 | 13.0 | 9.8 | 7.1 | 7.1 | -0.1 | 0.0 | 0.2 | 0.2 |
| 87 | 8.7 | 7.5 | 7.5 | 7.5 | 0.2 | 0.2 | 0.0 | 0.0 |
| 88 | 13.0 | 9.8 | 7.5 | 7.5 | 0.2 | 0.2 | 0.3 | 0.3 |
| 89 | 7.5 | 12.0 | 13.0 | 13.0 | 0.2 | 0.0 | -0.1 | -0.1 |
| 90 | 12.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.0 | -0.1 | -0.1 |
| 91 | 6.3 | 13.0 | 13.0 | 13.0 | 0.4 | -0.1 | -0.1 | -0.1 |
| 93 | 12.0 | 8.7 | 9.8 | 9.8 | 0.2 | 0.2 | 0.0 | 0.0 |
| 94 | 7.1 | 10.2 | 13.0 | 13.0 | 0.4 | 0.3 | 0.2 | 0.2 |
| 96 | 8.7 | 13.0 | 13.0 | 13.0 | 0.4 | 0.2 | 0.2 | 0.2 |
|  | 13.0 | 13.0 | 13.0 | 13.0 | 0.2 | 0.2 | 0.0 | 0.0 |
|  | 13.0 | 13.0 | 0.2 | 0.2 | 0.0 | 0.0 |  |  |


| ID | Age [Gyr] |  |  |  | $\mathrm{Z}[\log (\mathrm{Z} / \mathrm{Z} \odot)]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 99 | 13.0 | 13.0 | 7.5 | 7.5 | -0.4 | 0.0 | 0.3 | 0.3 |
| 100 | 12.0 | 13.0 | 12.0 | 12.0 | 0.2 | 0.0 | 0.0 | 0.0 |
| 101 | 13.0 | 12.0 | 7.5 | 7.5 | 0.2 | -0.1 | 0.0 | 0.0 |
| 104 | 13.0 | 13.0 | 13.0 | 13.0 | 0.0 | 0.0 | -0.1 | -0.1 |

Table C.20.: SSP results for the MACS 0416 sample using a color-color grid with $r_{625}-Y_{105}$ and $Y_{105}-H_{160}$. The first column contains the IDs of the galaxies. The ages and metallicities for the stellar populations at the respective radii are derived by finding the nearest grid points with their colors corresponding to modelled ages and metallicities.

| ID | Age $[\mathrm{Gyr}]$ |  |  |  | $\mathrm{Z}\left[\mathrm{Log}\left(\mathrm{Z} / \mathrm{Z}{ }_{\odot}\right)\right]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 21 | 5.7 | 8.7 | 7.5 | 6.3 | 0.4 | 0.0 | 0.2 | 0.3 |
| 25 | 8.7 | 13.0 | 13.0 | 7.5 | 0.0 | -0.1 | -0.1 | -0.1 |
| 26 | 12.0 | 13.0 | 12.0 | 6.3 | 0.0 | 0.0 | 0.0 | 0.3 |
| 30 | 6.6 | 8.7 | 8.7 | 7.5 | -0.1 | -0.4 | -0.4 | -0.4 |
| 31 | 10.2 | 13.0 | 13.0 | 10.2 | 0.0 | 0.0 | -0.1 | -0.4 |
| 33 | 6.3 | 13.0 | 13.0 | 13.0 | 0.4 | -0.1 | -0.1 | -0.1 |
| 38 | 6.3 | 13.0 | 13.0 | 9.8 | 0.3 | -0.4 | -0.4 | -0.4 |
| 43 | 7.5 | 13.0 | 13.0 | 7.5 | 0.0 | -0.4 | -0.4 | -0.1 |
| 45 | 6.3 | 7.5 | 7.1 | 6.3 | 0.4 | 0.3 | 0.3 | 0.4 |
| 50 | 7.1 | 13.0 | 13.0 | 7.5 | 0.4 | 0.0 | -0.1 | 0.0 |
| 53 | 7.1 | 13.0 | 13.0 | 13.0 | 0.4 | 0.2 | 0.0 | -0.4 |
| 54 | 7.5 | 13.0 | 13.0 | 13.0 | 0.0 | -0.1 | -0.1 | -0.4 |
| 55 | 7.1 | 5.7 | 6.3 | 5.7 | -0.1 | 0.3 | -0.4 | -0.1 |
| 57 | 9.8 | 13.0 | 13.0 | 7.5 | 0.2 | -0.1 | -0.1 | 0.2 |
| 63 | 12.0 | 13.0 | 13.0 | 9.8 | -0.1 | 0.0 | 0.0 | 0.0 |
| 65 | 10.2 | 12.0 | 13.0 | 7.5 | 0.2 | 0.0 | -0.1 | 0.0 |
| 66 | 9.8 | 7.1 | 7.1 | 10.2 | 0.0 | 0.2 | 0.2 | 0.0 |
| 68 | 6.6 | 10.2 | 8.7 | 6.3 | 0.3 | 0.0 | 0.0 | 0.3 |
| 70 | 7.1 | 13.0 | 13.0 | 12.0 | 0.3 | 0.2 | 0.0 | -0.1 |
| 71 | 6.6 | 13.0 | 13.0 | 6.6 | 0.4 | -0.1 | -0.1 | 0.2 |
| 85 | 13.0 | 12.0 | 10.2 | 7.1 | -0.1 | -0.1 | -0.1 | 0.2 |
| 87 | 7.1 | 12.0 | 8.7 | 6.6 | 0.4 | 0.0 | 0.0 | 0.2 |
| 88 | 8.7 | 12.0 | 13.0 | 12.0 | 0.2 | 0.0 | -0.1 | -0.1 |
| 93 | 6.3 | 9.8 | 10.2 | 6.3 | 0.4 | 0.2 | 0.0 | 0.3 |
| 94 | 10.2 | 8.7 | 7.5 | 7.1 | 0.0 | 0.0 | 0.0 | 0.0 |
| 95 | 5.7 | 10.2 | 7.5 | 5.7 | 0.4 | 0.0 | 0.2 | 0.4 |
| 96 | 7.1 | 8.7 | 7.5 | 6.3 | 0.3 | 0.0 | 0.0 | 0.2 |
| 98 | 7.1 | 7.5 | 8.7 | 13.0 | 0.0 | -0.4 | -0.4 | -0.4 |
| 102 | 10.2 | 13.0 | 13.0 | 13.0 | -0.1 | -0.4 | -0.4 | -0.4 |
| 103 | 8.7 | 13.0 | 12.0 | 6.3 | -0.4 | -0.1 | -0.1 | 0.2 |
| 104 | 13.0 | 12.0 | 10.2 | 7.5 | -0.1 | -0.1 | -0.1 | -0.1 |
| 105 | 8.7 | 13.0 | 13.0 | 13.0 | 0.3 | -0.1 | -0.4 | -0.4 |
| 107 | 10.2 | 13.0 | 13.0 | 12.0 | 0.3 | 0.0 | -0.1 | -0.1 |
| 109 | 10.2 | 8.7 | 7.5 | 6.3 | 0.0 | 0.2 | 0.2 | 0.3 |
| 113 | 7.1 | 10.2 | 12.0 | 7.1 | 0.4 | 0.2 | 0.0 | 0.3 |
|  |  |  |  |  |  |  |  |  |


| ID | Age [Gyr] |  |  |  | $\mathrm{Z}\left[\log \left(\mathrm{Z} / \mathrm{Z}_{\odot}\right)\right]$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ | $0.1 r_{e}$ | $0.5 r_{e}$ | $1 r_{e}$ | $2 r_{e}$ |
| 115 | 6.3 | 8.7 | 7.5 | 6.3 | 0.4 | 0.0 | 0.0 | 0.3 |
| 133 | 6.6 | 13.0 | 13.0 | 7.5 | 0.4 | 0.0 | 0.0 | 0.3 |
| 154 | 13.0 | 13.0 | 9.8 | 6.6 | 0.0 | 0.0 | 0.0 | 0.2 |


[^0]:    ${ }^{1} \mathrm{https}: / /$ archive.stsci.edu/prepds/clash/

[^1]:    ${ }^{1}$ http://www.iaa.es/sites/default/files/MOLINO\%20BENITO.pdf

[^2]:    ${ }^{1}$ Except in regard of computational time, as described in chapter 4.

[^3]:    ${ }^{2}$ The formation redshifts were $z_{f}=[0.5,0.55,0.6,0.7,0.75,0.85,0.95,1.3,1.75,2,4,10]$

